

Effects of topological defects on the CMB

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PFNG - HRI - 2010

[arxiv:1010.5662](#), [1005.2663](#), [0911.1241](#), [0908.0432](#), [0812.1929](#), [0803.2059](#), [0711.1842](#),

[0704.3800](#), [astro-ph/0702223](#)

Inflationary cosmology & topological defects

- Simplest model of the early Universe: Inflation^a
- Some models of inflation (“hybrid”) end by producing topological defects^b
- Defects may also be formed in subsequent thermal transitions^c
- String/M-theory: defects from $(D\bar{D})$ -brane collisions^d
- Defects have gravitational fields & contribute to perturbations^e

^aStarobinsky (1980); Sato (1981); Guth (1981); Mukhanov & Chibisov (1981); Linde (1982); Hawking & Moss (1982); Albrecht & Steinhardt (1982); Guth & Pi (1982); Hawking (1982); Hawking & Moss (1983); Bardeen, Steinhardt, Turner, (1983)

^bYokoyama (1989); Copeland et al (1994); Kofman, Linde, Starobinski (1996); Garcia-Bellido et al (2010)

^cKibble (1976); Zurek (1996); Rajantie (2002)

^dJones, Stoica, Tye (2002); Dvali & Vilenkin (2003); Copeland, Myers, Polchinski (2003)

^eKibble (1976); Zel'dovich (1980); Vilenkin (1981)

Topological defects - cosmic strings

- Cosmic strings^a are linear distributions of mass-energy in the universe.
- Mass per unit length μ , tension T . Normally $\mu = T/c^2$
- Dynamics: **acceleration** $\propto$ **curvature**: wave equation
- In theories of high energy physics they may be
 - **Elementary** (string theory): **zero width**
 - **Solitonic** (field theory): **non-zero width**
- Made in the early universe?^b $t \sim 10^{-36}$ s, $\mu \sim 10^{32}$ GeV², $w \sim 10^{-30}$ m
- If formed, still here: O(1) **“infinite” string**, unknown distribution of **closed loops**

^aHindmarsh & Kibble (1994); Vilenkin & Shellard (1994); Kibble (2004)

^bKibble (1976); Zurek (1996); Rajantie (2002); Yokoyama (1989); Kofman, Linde, Starobinski (1996); Jones, Stoica, Tye (2002); Sarangi & Tye (2003); Copeland, Myers, Polchinski (2003); Dvali & Vilenkin (2003)

Other defects: global monopoles, textures and semilocal strings**Global monopoles and textures^a**

- Self-ordering scalar fields (Goldstone modes) from global symmetry-breaking.
- Global monopoles: point-like, with attractive force proportional to distance.
- Symmetry-breaking scale $v \sim 10^{16}$ GeV: observable perturbations.

Semilocal strings^b

- Self-ordering scalar and vector fields from “semilocal” symmetry-breaking
- Semilocal: non-trivial combination of local and global symmetries
- Symmetry-breaking scale $v \sim 10^{16}$ GeV: observable perturbations.

^aTurok 1989; Spergel et al 1991; Pen, Spergel, Turok 1995; Durrer, Kunz, Melchiorri 1999,2002.

^bVachaspati, Achúcarro 1991; Hindmarsh 1992,1993; Urrestilla et al. 2008.

Observational signals from defects

Predictions under good theoretical control:

- Cosmic Microwave Background (power spectrum, bi/tri-spectrum) ^a

Some theoretical uncertainties:

- Density perturbation (power spectrum); ^b [strings: Gravitational lensing^c]

Large theoretical uncertainties:

- Gravitational radiation^d [Global strings: axion radiation]; ^e [strings: cosmic rays^f]

^aGott (1984); Kaiser & Stebbins (1984); Bouchet et al. (1989); Allen et al (1996,7); Landriau & Shellard (2004,10); Wyman et al (2005); Pogosian, [Tye], Wyman (2008,9); Bevis et al (2006,7,8,10); Hindmarsh, Ringeval, Suyama (2009,10); Regan, Shellard (2010).

^bZel'dovich (1980); Vilenkin (1981); Avelino, Shellard, Wu, Allen (1998); ...

^cVilenkin (1984); Hindmarsh (1989); de Laix & Vachaspati (1996,1997); Mack, Wesley, King (2007); ...

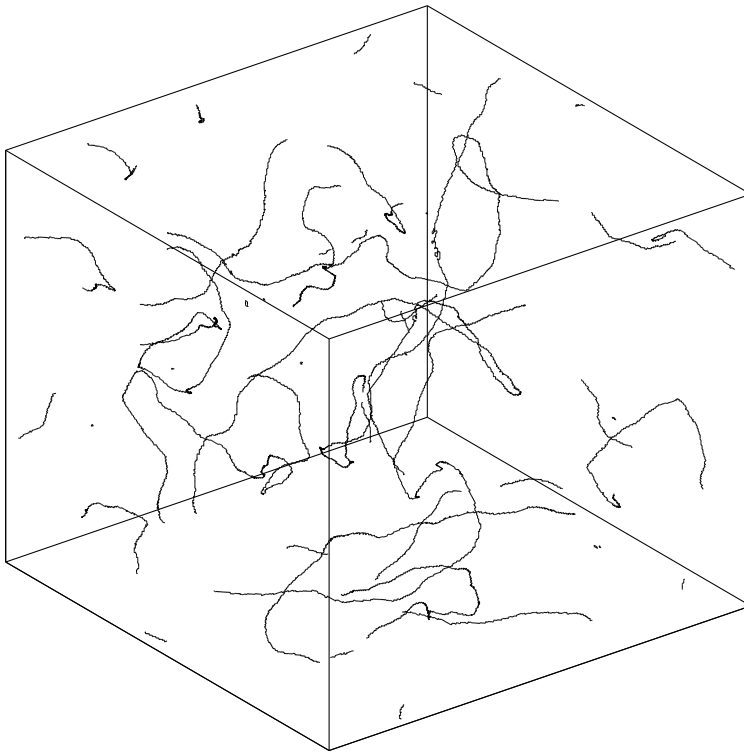
^dVachaspati & Vilenkin (1985); Hindmarsh (1990); Allen & Shellard (1992); Damour & Vilenkin (2000,2001,2005); ...

^eBattye & Shellard (1994-9); Yamaguchi, Kawasaki, Yokoyama (1998); Wantz & Shellard (2009).

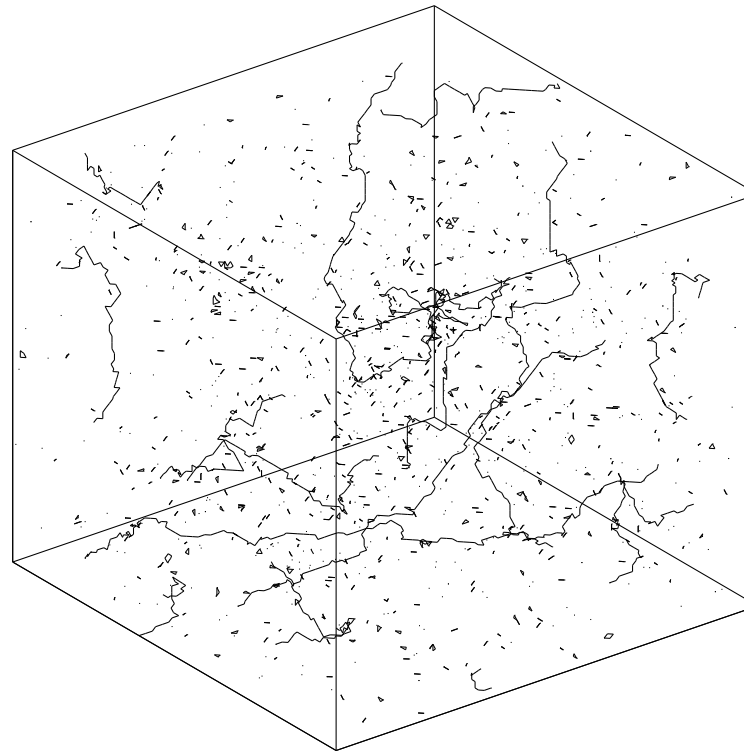
^fBhattacharjee (1990); Sigl (1996); Protheroe (1996); Berezhinski (1997); Vincent, M.H., Sakellariadou (1998); Wichowski, MacGibbon, Brandenberger (1998)

A problem for strings: small scales and loops

Classical Abelian Higgs model

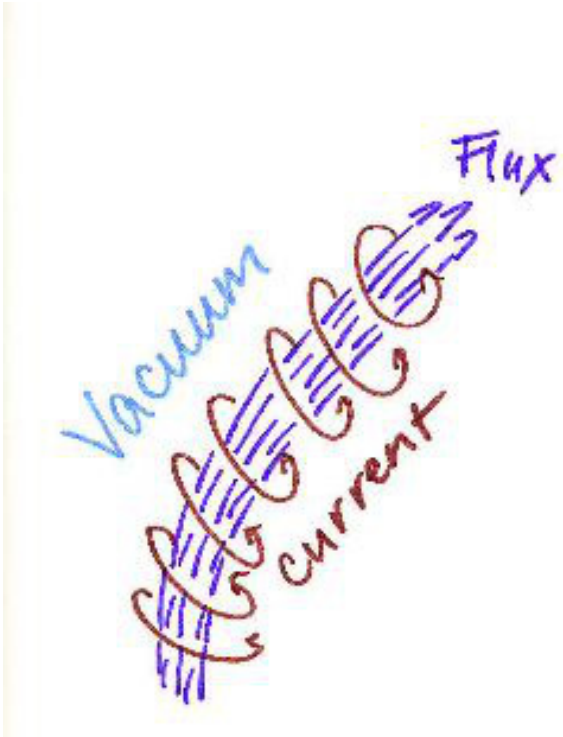


Nambu-Goto strings



How much string? How many loops?

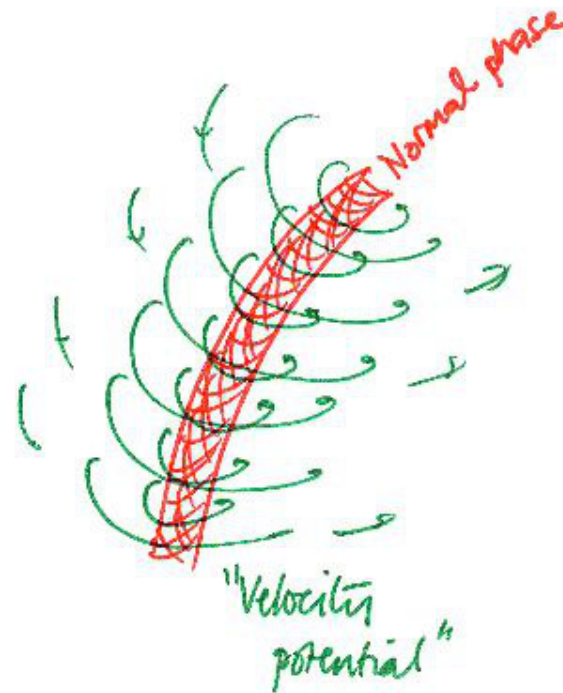
Defect zoo: strings from field theory



Gauge/local string

Nearest living relative:

Type II superconductor flux tube



Global string

Nearest living relative:

superfluid vortex

String solutions in the Abelian Higgs model

Complex **scalar** field $\phi(\mathbf{x}, t)$, **vector** field $A_\mu(\mathbf{x}, t)$

$$\mathcal{L} = (D\phi)^\dagger \cdot (D\phi) - V(|\phi|) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

Static finite energy (2D) cylindrically symmetric:

$$\phi = v f(r) e^{i\theta}, \quad A_i = \frac{1}{er} A_\theta(r) \hat{\theta}_i$$

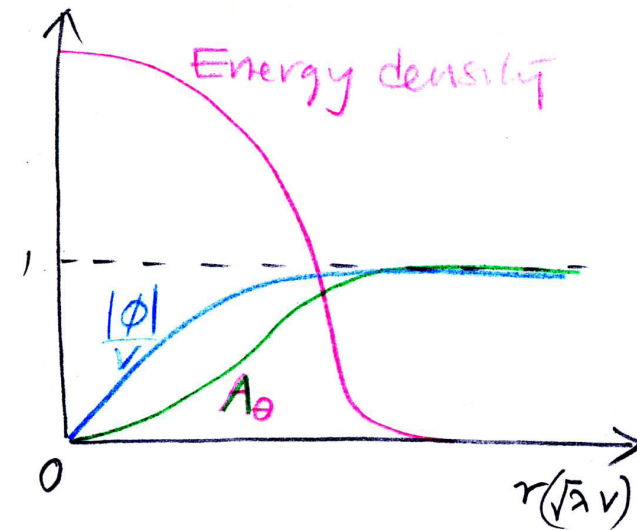
Energy density:

$$\rho = |D_i\phi|^2 + V_\lambda(\phi) + \frac{1}{2}B^2$$

ρ confined to region $r < \max(1/\sqrt{\lambda}v, 1/ev)$

String mass per unit length: $\mu = 2\pi v^2 E(\lambda/e^2)$

[E - slow function, $E(1) = 1$]



Textures

- N scalar fields $\phi^A(\mathbf{x}, t)$
- $O(N)$ symmetry: $\mathcal{L} = \frac{1}{2} \partial \phi^A \cdot \partial \phi^A - V(|\phi|)$
- Potential: $V(|\phi|) = \frac{1}{4} \lambda (\phi^A \phi^A - v^2)^2$

- Energy density:

$$\rho = \frac{1}{2} \dot{\phi}^2 + (\nabla \phi)^2 + V$$

- At low energy density, $|\phi| \simeq v$ [global symmetry is broken to $O(N - 1)$].
- Cosmic texture: field configuration with wavelength $\xi \sim t$
- Evolution timescale t .
- Energy density is $\rho \sim v^2/t^2$
- Textures exist in any kind of non-Abelian global symmetry-breaking.

Semilocal strings

- N complex scalar fields $\phi^A(\mathbf{x}, t)$
- $U(N)$ symmetry: $\mathcal{L} = (D\phi^A)^\dagger \cdot (D\phi^A) - V(|\phi|) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$
- $U(N) = SU(N)_{\text{global}} \times U(1)_{\text{local}}$.
- Potential: $V(\phi^A\phi^A) = \frac{1}{2}\lambda((\phi^A)^\dagger\phi^A - v^2)^2$
- At low energy density, $|\phi| \simeq v$ [symmetry is broken to $SU(N-1)_{\text{global}}$].
- For gauge coupling $\gg$ scalar coupling: like Abelian Higgs model
- For gauge coupling $\ll$ scalar coupling: Semilocal string
- A semilocal string is like texture, a field configuration with wavelength $\xi \sim t$
- Evolution timescale t .
- Energy density is $\rho \sim v^2/t^2$

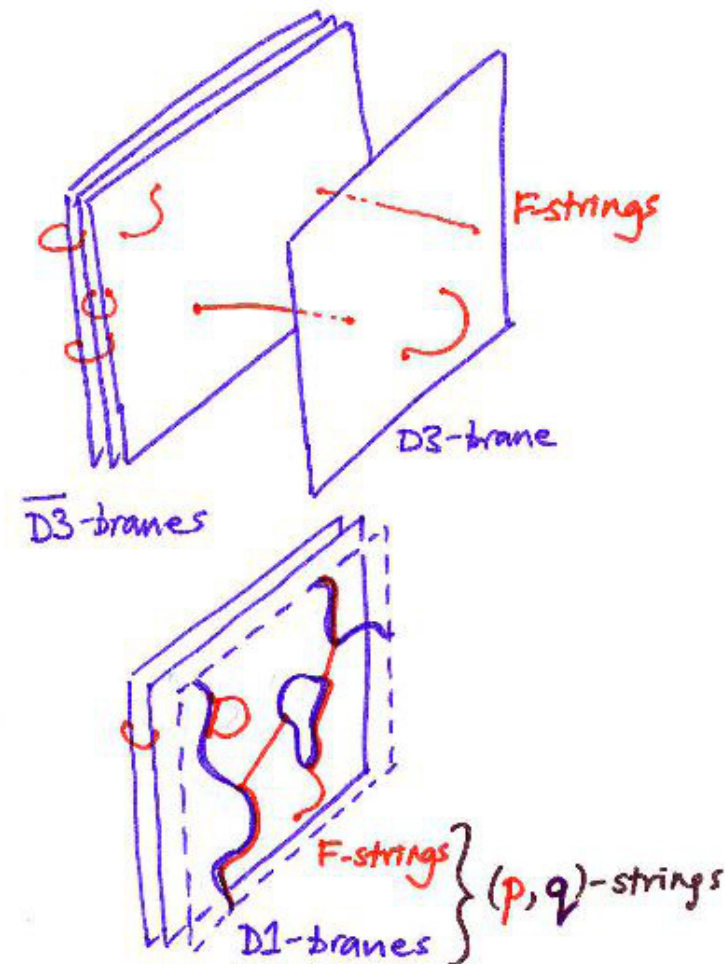
Cosmic strings from string theory

- Fundamental strings **F-strings**
- Extended objects **D-branes**
- F-strings end on D-branes, (**D1 = D-string**)
- Bound states, junctions: **(p, q) -strings^a**
- String tension $\mu_{p,q} = \frac{V_{6,\text{eff}}}{2\pi\alpha'} \sqrt{p^2 + \frac{q^2}{g_s^2}}$ ^b
- Formation: collisions of D3- $\overline{\text{D3}}$ branes^c

^aCopeland, Myers, Polchinski (2004); Firouzjahi, Leblond, Tye (2006); Dasgupta, Firouzjahi, Gwyn (2007)

^bCopeland, Myers, Polchinski (2004); Firouzjahi, Tye (2005)

^cSarangi & Tye (2002); Dvali & Vilenkin (2004); Barnaby, Berndsen, Cline, Stoica (2005)



Defect scaling hypothesis

- Defects have a characteristic scale ξ (e.g. string curvature radius $\sqrt{V/L}$)
- Scaling hypothesis: $\xi = x_* t$ (x_* constant $O(1)$)
- Energy density: $\rho_s \simeq v^2/\xi^2$
- Total energy density: $\rho_t \sim 1/Gt^2$:
- Defect density fraction: $\Omega_d \sim Gv^2/x_*^2$ - is **constant**.
- Grand Unification: $Gv^2 \sim 10^{-6}$

Scaling: extrapolate from $t_i \sim 10^{-36}$ s to $t_0 \sim 3 \times 10^{17}$ s today

A supersymmetric model of (hybrid) inflation ...

$\nu\text{MSSM} + \phi + \bar{\phi} + s + W_{\text{mix}}$ [extra U(1)' global, gauged]^a

- For example: $W = W_{\nu\text{MSSM}} + \kappa s(\phi\bar{\phi} - M^2)$
- $V = V_D + V_{\text{soft}} + \kappa^2 |s|^2 (|\phi|^2 + |\bar{\phi}|^2) + \kappa^2 |\phi\bar{\phi} - M^2|^2 + V_{1\text{-loop}}$
- $V_D = \frac{e'^2}{2} (\xi_{\text{FI}} - |\phi|^2 + |\bar{\phi}|^2)^2$
- Inflation for large $|s|$, $\phi = \bar{\phi} = 0$. F-term^b or D-term^c can dominate.
- Inflation ends at $|s| = |s_c(M, \xi_{\text{FI}})|$, $\phi, \bar{\phi}$ get vevs: U(1)' symmetry breaks.

^ae.g. Garbrecht, Pallis, Pilaftsis (2006)

^bCopeland et al (1994)

^cBinétruy & Dvali (1996), Halyo (1996)

... is a model with cosmic strings

- $U(1)'$ symmetry is broken after inflation: $\langle |\phi|^2 \rangle \simeq \phi_0^2(M, \xi_{FI})$
- Gauged $U(1)'$: Abelian Higgs cosmic strings with tension $\mu = E(\kappa^2/e'^2)2\pi\phi_0^2$
- Global $U(1)'$: Global cosmic strings with tension $\mu \sim \phi_0^2 \log(t\phi_0)$
- Inflation + strings - a natural paradigm
- Tight constraints from CMB^a

^aBattye, Garbrecht, Moss (2008, 2010); Battye, Moss (2010)

Formation and evolution: Abelian Higgs in FRW background

$$S = - \int d^4x \sqrt{-g} \left(g^{\mu\nu} D_\mu \phi^* D_\nu \phi + V(\phi) + \frac{1}{4e^2} g^{\mu\rho} g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma} \right),$$

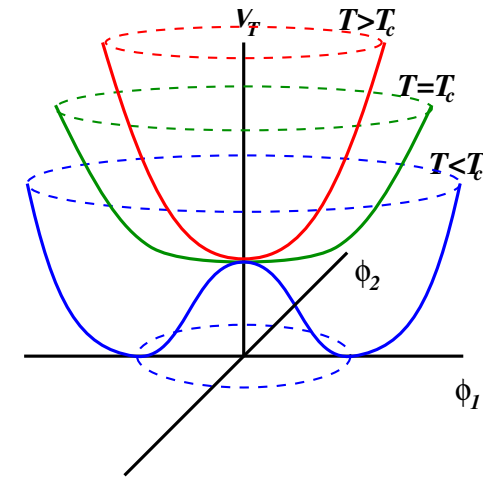
Complex **scalar** field $\phi(\mathbf{x}, t)$, **vector** field $A_\mu(\mathbf{x}, t)$

Covariant derivative $D_\mu = \partial_\mu - iA_\mu$.

Potential $V(\phi) = \frac{1}{2}\lambda(|\phi|^2 - v^2)^2$.

Metric $ds^2 = a^2(\tau)(-d\tau^2 + d\mathbf{x}^2)$

τ : **conformal time**, $\propto t, t^{\frac{1}{2}}$



Temporal gauge ($A_0 = 0$) field equations (index raised with Minkowski metric).

$$\ddot{\phi} + 2\frac{\dot{a}}{a}\dot{\phi} - D^2\phi + \lambda a^2(|\phi|^2 - v^2)\phi = 0,$$

$$\partial^\mu \left(\frac{1}{e^2} F_{\mu\nu} \right) - ia^2(\phi^* D_\nu \phi - D_\nu \phi^* \phi) = 0,$$

Visualising Abelian Higgs string simulations

Isosurfaces of constant energy density. Size: 256^3 , lattice spacing $0.5m^{-1}$

Initial conditions: Gaussian random field ϕ

Boundary conditions: toroidal

Abelian Higgs model simulations: string length scale

Scaling: $L/V \propto \tau^{-2}$

Network scale: $\xi = \sqrt{(V/L)}$

Hence $\xi \propto \tau$

Couplings: $\lambda = 2, e = 1$

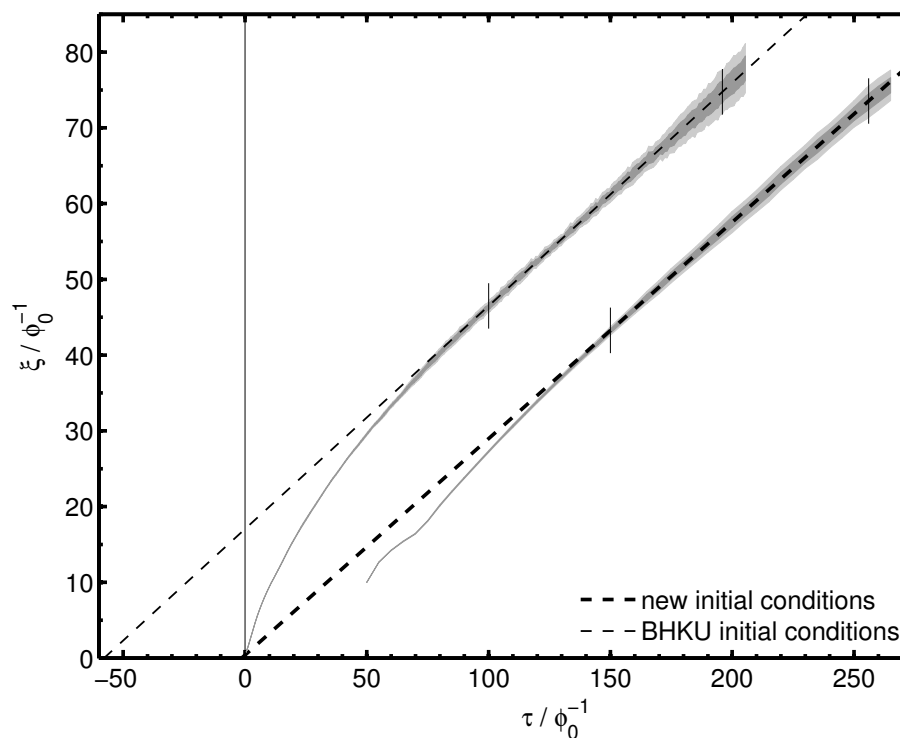
Masses: $m_s = m_v = 1$

Lattice spacing: $\Delta x \simeq 0.5$

Time step: $\Delta t = 0.1$

Volume: $768^3, 1024^3$

Matter era, different initial conditions



Semilocal strings simulations: length scale

Scaling: Energy density $T_{00} \propto \tau^{-2}$

Couplings: $\lambda = 2, e = 1$

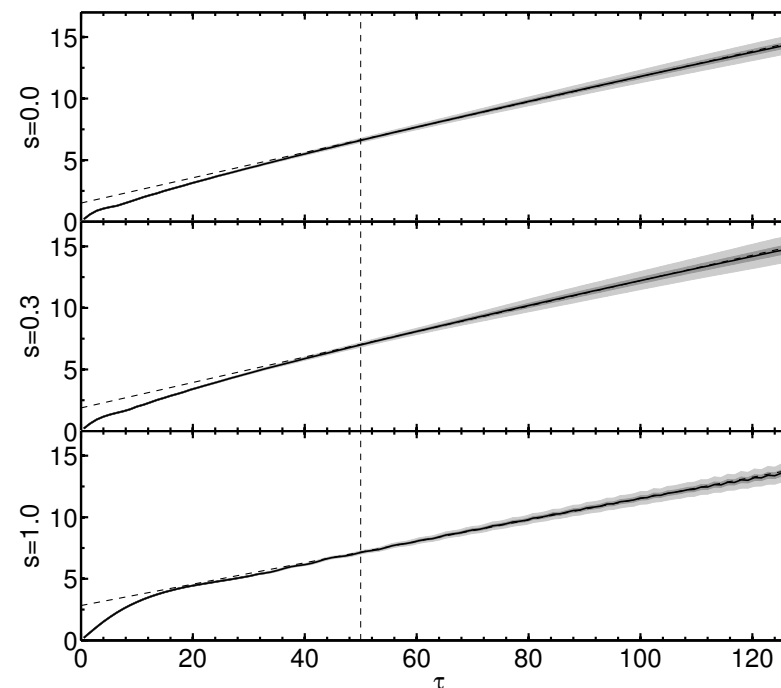
Masses: $m_s = m_v = 1$

Lattice spacing: $\Delta x \simeq 0.5$

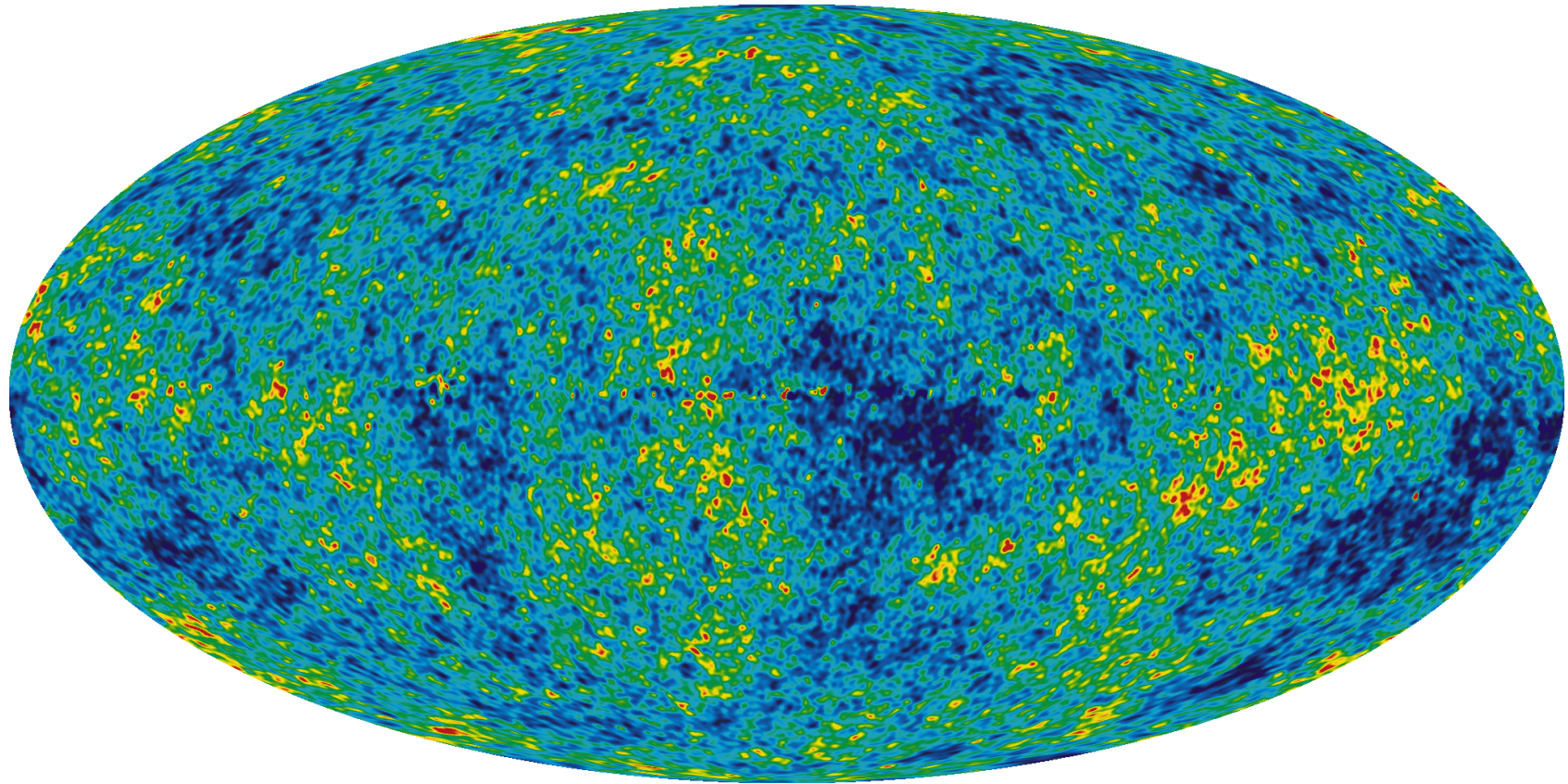
Time step: $\Delta t = 0.1$

Volume: 512^3

Radiation era, different algorithms



Cosmic Microwave Background constraints



$$-200 < \Delta T < 200 \mu\text{K}$$

Gott-Kaiser-Stebbins (GKS) effect for cosmic strings

Discontinuity^a $\Delta T \sim 8\pi(G\mu)vT_{\text{CMB}}$

Need high resolution & high sensitivity

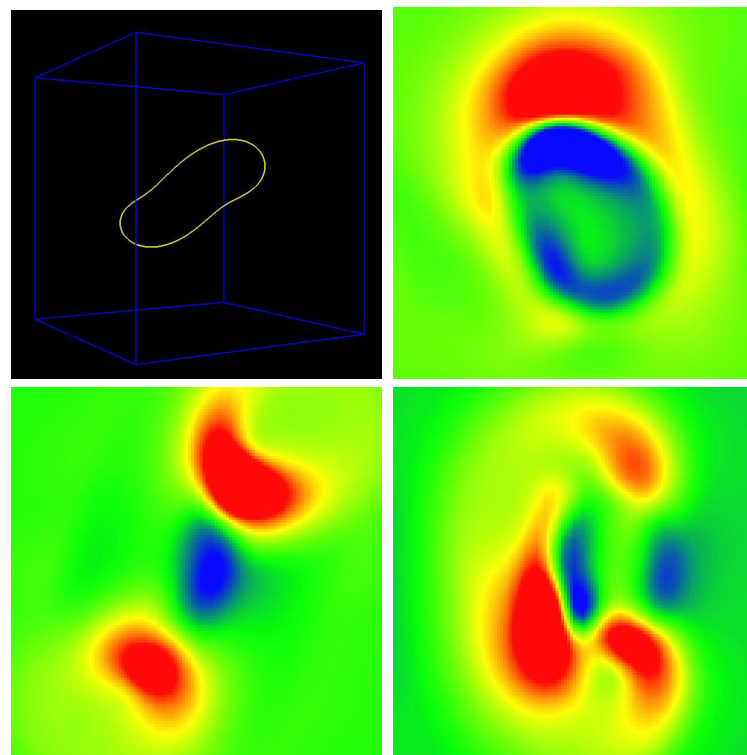
- OVRO + string correlators:^b
 $(G\mu) < x_* 11 \times 10^{-6}$
- WMAP + random edge model:^c
 $G\mu < 1.07 \times 10^{-5}$
- WMAP3 + random edge model:^d
 $G\mu < 3.7 \times 10^{-6}$

^aKaiser, Stebbins (1984)

^bHindmarsh (1993)

^cLo & Wright (2005)

^dJeong & Smoot (2006)



Landriau and Shellard (2002)

Calculating CMB power spectra from defects: UETC method

$h_\alpha(\tau, k)$: linear perturbation (metric, matter, temperature ...)

$S_\alpha(\tau, k)$: source (string energy-momentum, separately conserved)

$D_{\alpha\beta}(\tau, k)$: time dependent differential operator

Perturbation equation: $D_{\alpha\beta}(\tau, k)h_\beta(\tau, k) = S_\alpha(\tau, k)$

Power spectrum:^a $\langle |h_\alpha(\tau_0, k)|^2 \rangle = \int \int \mathcal{D}^{-1} \mathcal{D}^{-1} \langle S_\alpha(\tau, k) S_\alpha^*(\tau', k) \rangle$

Need **unequal-time correlators** (UETCs) of source (energy-momentum tensor)

$$C_{\mu\nu\rho\lambda}(k, \tau, \tau') = \langle T_{\mu\nu}(k, \tau) T_{\rho\lambda}^*(k, \tau') \rangle$$

5 independent UETCS [**3 scalar, 1 vector, 1 tensor**] from numerical simulations

$\mathcal{D}^{-1}$ is CMBEASY, applied to eigenvectors of UETCs

Scaling: small times, lengths $\rightarrow$ large times, lengths

^aPen, Seljak, Turok (1997); Dürer, Kunz, Melchiorri (1998,2002)

Fitting CMB with inflation & cosmic defects

- Two sources of perturbations: **incoherent** - **add in quadrature**
- Cosmological model with 1 more parameter: μ
- Gauge and semilocal strings: $\mu = 2\pi v^2 E(\lambda/e^2)$
[λ - scalar coupling, e - gauge coupling, E - slow function, $E(1) = 1$]
- Textures: $\mu = 2\pi v^2$
- Use $f_{10} = C_{10}^{\text{defect}} / C_{10}^{\text{total}}$. Proportional to $(G\mu)^2$.
- Modify COSMOMC and perform Monte Carlos.

CMB from gauge strings, textures, and semilocal strings

Multipole moments:

$$a_{lm} = \int d\Omega \Delta T(\mathbf{n}) Y_{lm}^*(\mathbf{n})$$

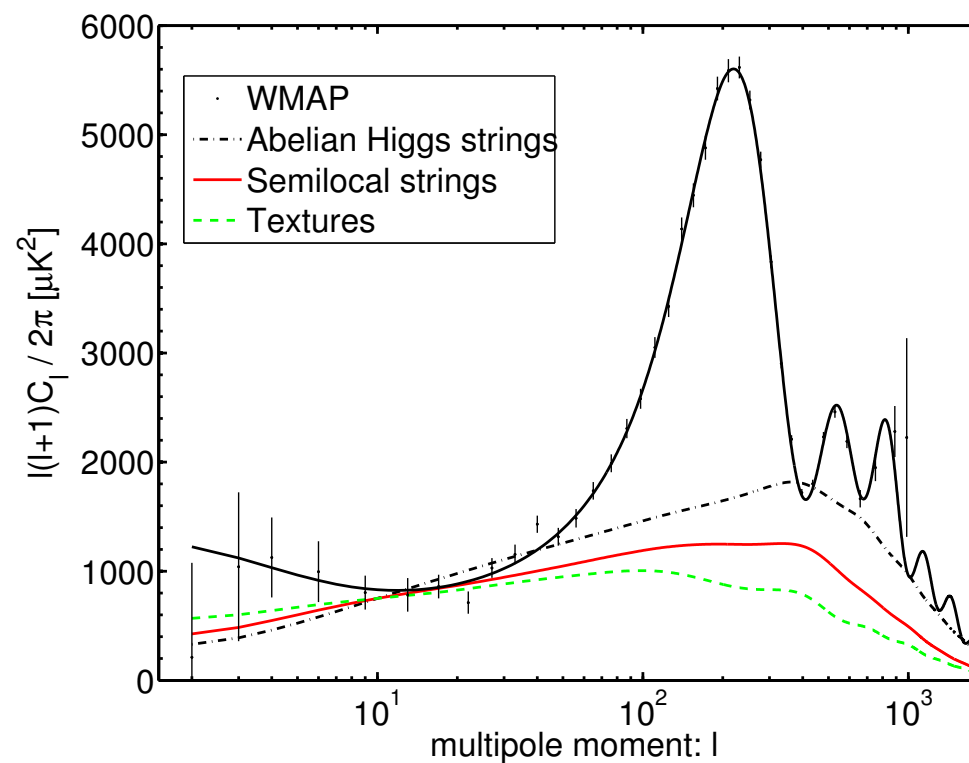
Angular power spectrum:

$$C_l = \sum_{m=-l}^l |a_{lm}|^2$$

Anisotropy power:

$$l(l+1)C_l/(2\pi)$$

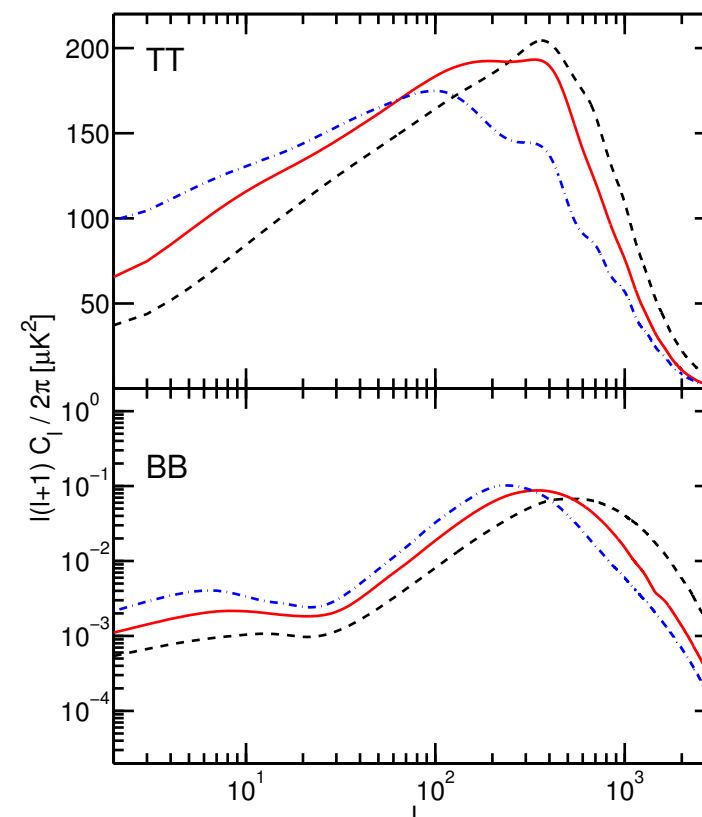
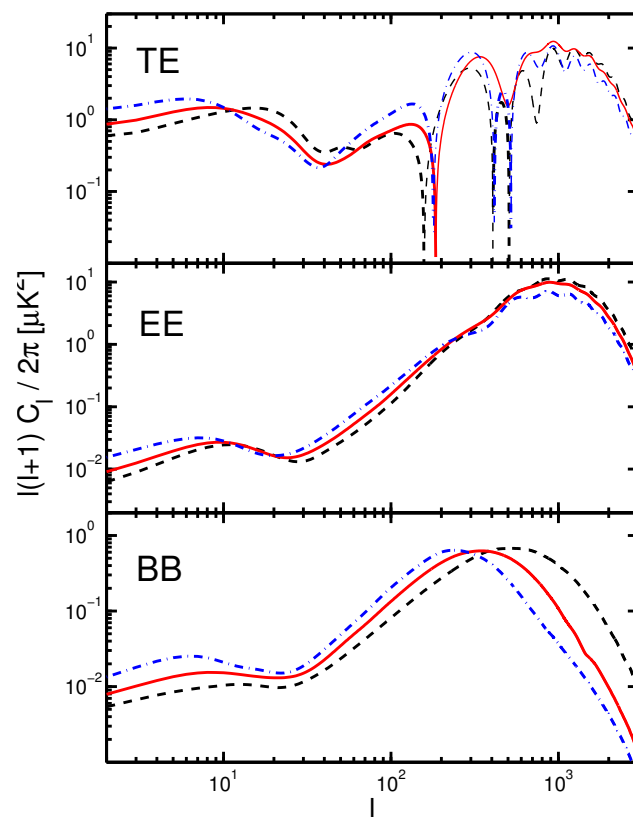
	$G\mu_{10}$
Abelian Higgs string	2.7×10^{-6}
Semilocal string	5.3×10^{-6}
Textures	4.5×10^{-6}



Defects normalised to WMAP3 ($\ell = 10$)^a

^aUrrestilla, Bevis, Hindmarsh, Kunz, Liddle (2008)

Polarisation from gauge strings, textures, and semilocal strings

Defects normalised to WMAP3 ($\ell = 10$)

Defects normalised to 95% c.i. upper bound

Results: WMAP3, ACBAR, BOOMERANG, CBI and VSA

SL: Semilocal strings

TX: textures

AH: Abelian Higgs strings

HKP: Gaussian prior on H_0 from Hubble Key Project

BBN: Gaussian prior on $\Omega_b h^2$ from Big Bang Nucleosynthesis

95% upper bounds on f_{10} , $G\mu$ and n_s for standard 6-parameter cosmology + X

Model	CMB only			CMB+BBN+HKP		
	f_{10}	$G\mu$	n_s	f_{10}	$G\mu$	n_s
SL	0.25	2.6×10^{-6}	1.09	0.14	2.0×10^{-6}	1.01
TX	0.33	2.5×10^{-6}	1.14	0.16	1.8×10^{-6}	1.02
AH	0.17	1.1×10^{-6}	1.06	0.10	0.9×10^{-6}	1.00

Results preview

8 parameter model (Standard Cosmology + $G\mu$ + Sunyaev-Zeldovich)

MCMC fit to CMB data (WMAP7+QUAD+ACBAR)

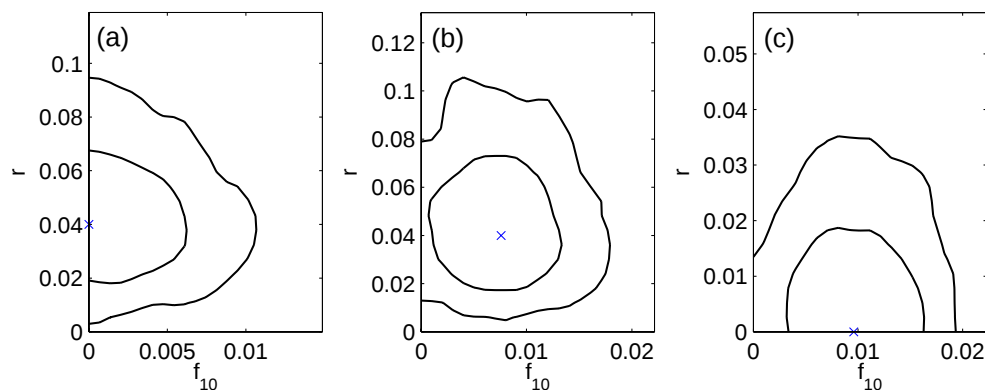
Cosmic string perturbations from Classical Abelian Higgs model^a

$$f_{10} < 0.056 \text{ (95\%)} \quad G\mu < 0.58 \times 10^{-6} \text{ (95\%)}$$

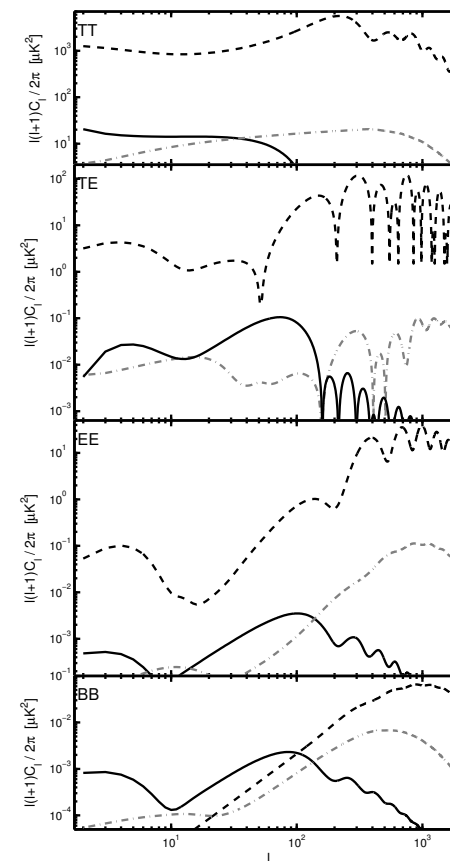
^aBevis et al. (2010); Bevis et al (in preparation)

Planck: distinguishing defects & tensors

- Inflation - **B-modes** from **gravitational waves**
- Defects - **B-modes** from **vector modes**
- Parameters: $r = \frac{|A_{\text{tensor}}|^2}{|A_{\text{scalar}}|^2}$, $f_{10}^{\text{string}} = \frac{C_{10}^{\text{TT,string}}}{C_{10}^{\text{TT,total}}}$
- Planck can distinguish, $f_{10} \gtrsim 0.02^a$



^aUrrestilla et al. 2008



scalar, black-dashed

$r = 0.04$, solid

$f_{10} = 0.01$, grey dot-dash

Future B-mode satellite: detecting strings/textures/tensors

- Threshold detection, assuming true model

known (3σ):^a [for CMBpol. COrE similar]

- $f_{10}^{\text{string}} \gtrsim 0.0012$ ($G\mu = 7 \times 10^{-8}$)
- $f_{10}^{\text{texture}} \gtrsim 0.0005$ ($G\mu = 1.0 \times 10^{-7}$)
- $r \gtrsim 0.0018$

- No foreground, iterative delensing:

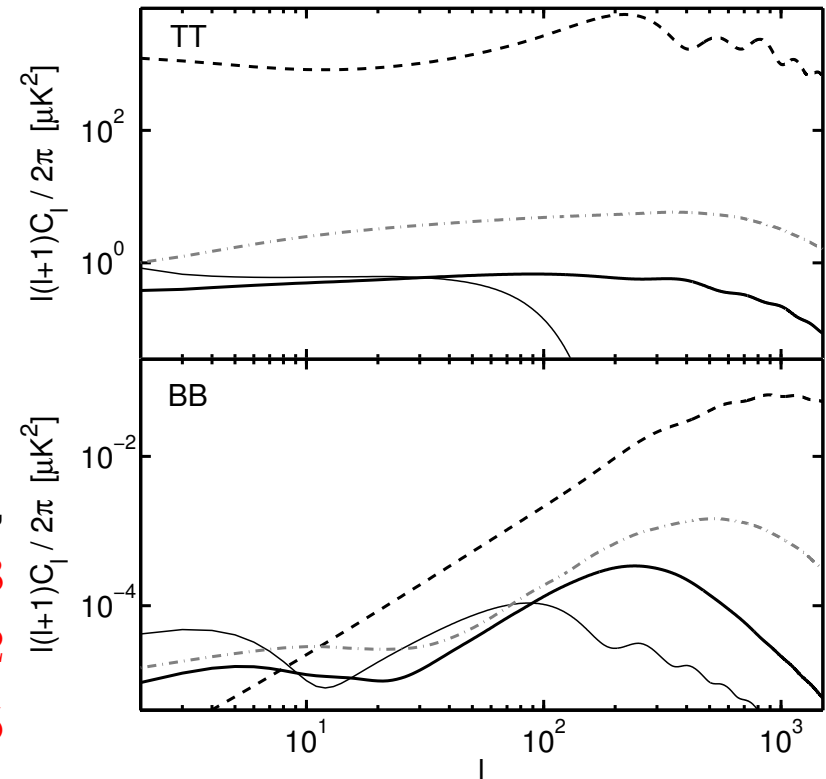
- $f_{10}^{\text{string}} \gtrsim 10^{-4}$ ($G\mu \simeq 8 \times 10^{-9}$)^b

scalar (black dashed),

tensor (black thin), $r = 0.0018$

cosmic strings (grey dot-dashed), $f_{10}^{\text{string}} = 0.0012$

textures (black thick) $f_{10}^{\text{texture}} = 0.0005$



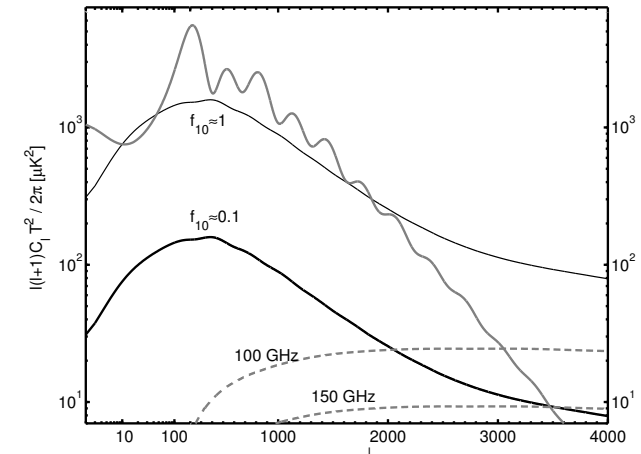
^aMukherjee et al (2010)

^bSeljak, Slosar (2006); Garcia-Bellido et al (2010)

Cosmic string CMB power spectrum at small angular scales

- Inflationary scalar fluctuations damped at high ℓ
- Strings have ℓ^{-1} behaviour for $\ell > 3000^a$
- Sunyaev-Zeldovich dominates both at high ℓ
- SZ may be removable - frequency dependence

^aHindmarsh (1994); Fraisse et al (2007); Pogosian, Tye, Wyman (2009); Bevis et al (2010); Yamauchi et al (2010)



Bevis et al (2010)

Cosmic string CMB bi & trispectrum at small angular scales

Temperature fluctuation $\Theta_{\mathbf{k}}$, flat sky, 2D wave vector $\mathbf{k}$: small angle approx

$$\langle \Theta_{\mathbf{k}_1} \Theta_{\mathbf{k}_2} \Theta_{\mathbf{k}_3} \rangle = b_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3} (2\pi)^2 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3).$$

$$\langle \Theta_{\mathbf{k}_1} \Theta_{\mathbf{k}_2} \Theta_{\mathbf{k}_3} \Theta_{\mathbf{k}_4} \rangle = T_{\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \mathbf{k}_4} (2\pi)^2 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3 + \mathbf{k}_4).$$

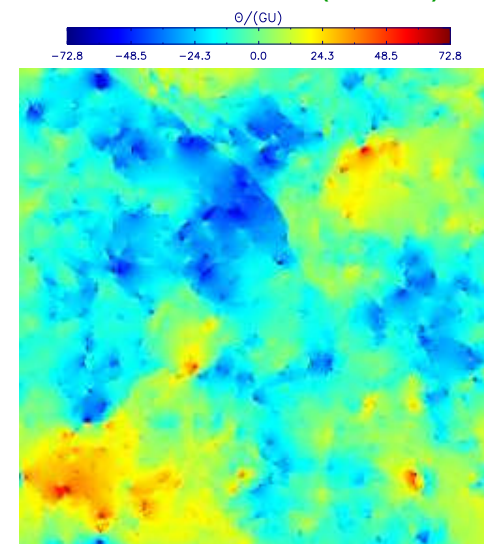
Bi & trispectrum from Gott-Kaiser-Stebbins effect only^a

- $\left[\frac{\ell(\ell+1)}{2\pi} \right]^{3/2} b_{\ell\ell\ell} \simeq (-0.5) (G\mu)^3 \left(\frac{1000}{\ell} \right)^{-3},$
- GKS bispectrum **negative**^b, $O(C_\ell^{3/2})$
- $T_{\ell\ell\ell\ell} \sim (G\mu)^4 \ell^{-\rho}$, with $6 < \rho(v_{\text{rms}}) < 7$
- GKS trispectrum $O(C_\ell^2)$
- **Must include fluid for $\ell \lesssim 3000$**

^aHindmarsh, Ringeval, Suyama (2009,10); Regan, Shellard (2009)

^bNegative skewness Fraise et al (2007); Yamauchi et al (2010)

Fraise et al (2007);



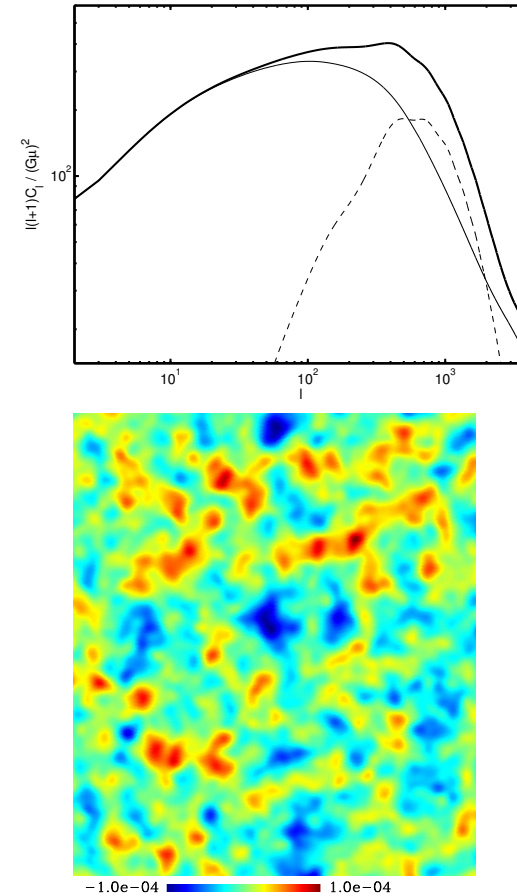
Size 7.2° , resol'n $0.41'$

Cosmic string CMB non-Gaussianity $100 < \ell < 3000$

- C_ℓ s dominated by last scattering (right).^a
- GKS tricks don't work for fluid
- String CMB maps needed^b
- e.g. 3° (right)
- GKS “edges” smeared by fluid response
- Skewness **positive**

^aBevis et al (2010)

^bLandriau & Shellard 2010



Strings and 21cm radiation background

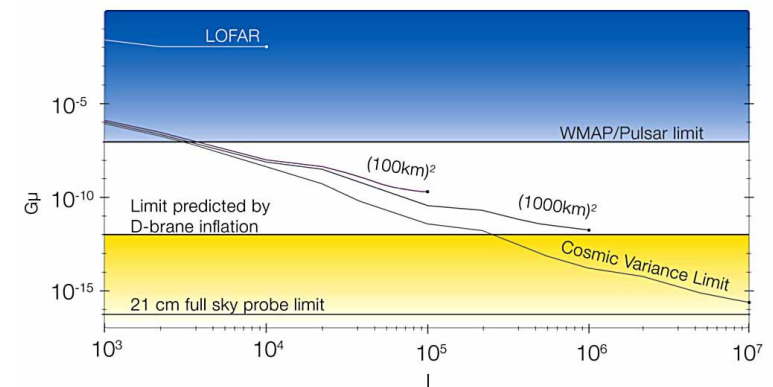
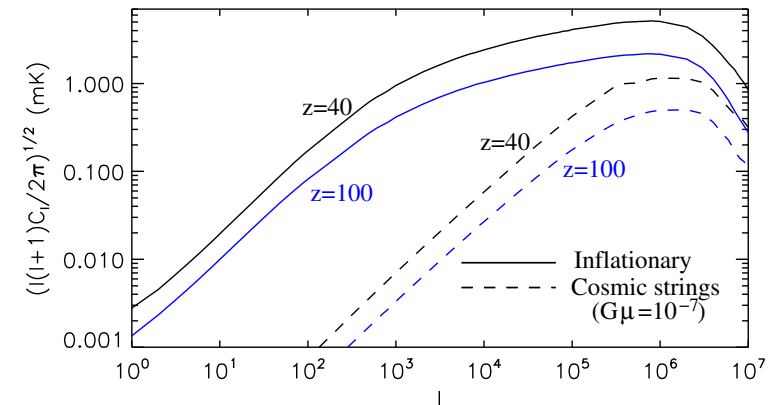
- Strings add to 21cm power spectrum^a
 - Need huge collecting area: $\gg 1 \text{ km}^2$
- Weakly cross correlated with CMB^b
 - X-correlation also needs huge collecting area to see $G\mu \sim 5 \times 10^{-7}$
- Cosmic string wakes observable at $z \sim 30$ with $G\mu \simeq 6 \times 10^{-7}$ ^c

^aKhatri, Wandelt (2008)

^bBerndsen, Pogosian, Wyman (2010)

^cBrandenberger et al (2010)

Khatri & Wandelt 2008:



Conclusions

- Well-motivated inflation models produce defects
- (Inflation + defect) cosmological models have one extra parameter, $G\mu$
- CMB indicates that $G\mu \lesssim (5 \times 10^{-7}) - 10^{-6}$
- Future:
 - CMB:** Planck, CMBpol/COrE/ground-based B-modes
 - Gravitational radiation (strings):** AdvLIGO, LISA
 - Cosmic rays (strings):** Fermi
 - Lensing (strings):** SKA (Compact Radio Sources)
 - 21cm (strings):** SKA