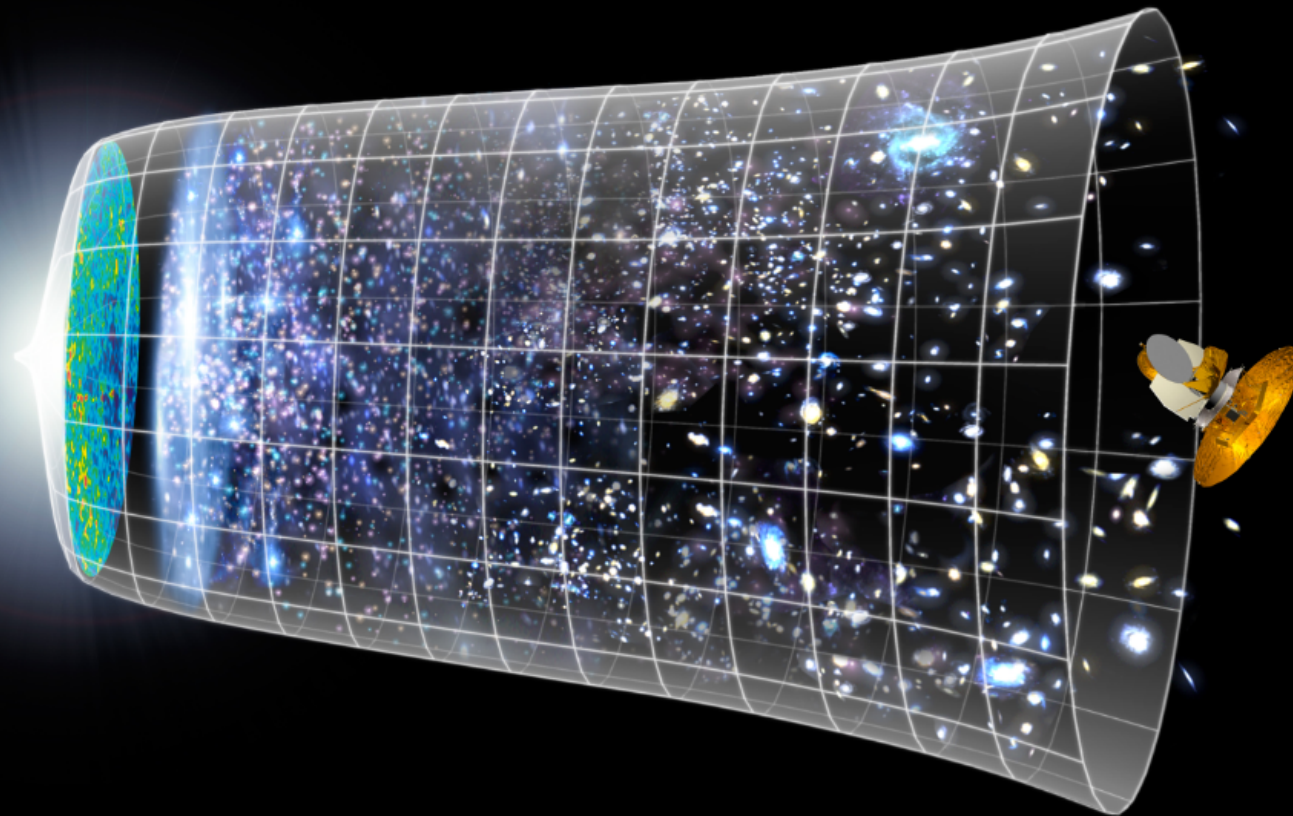
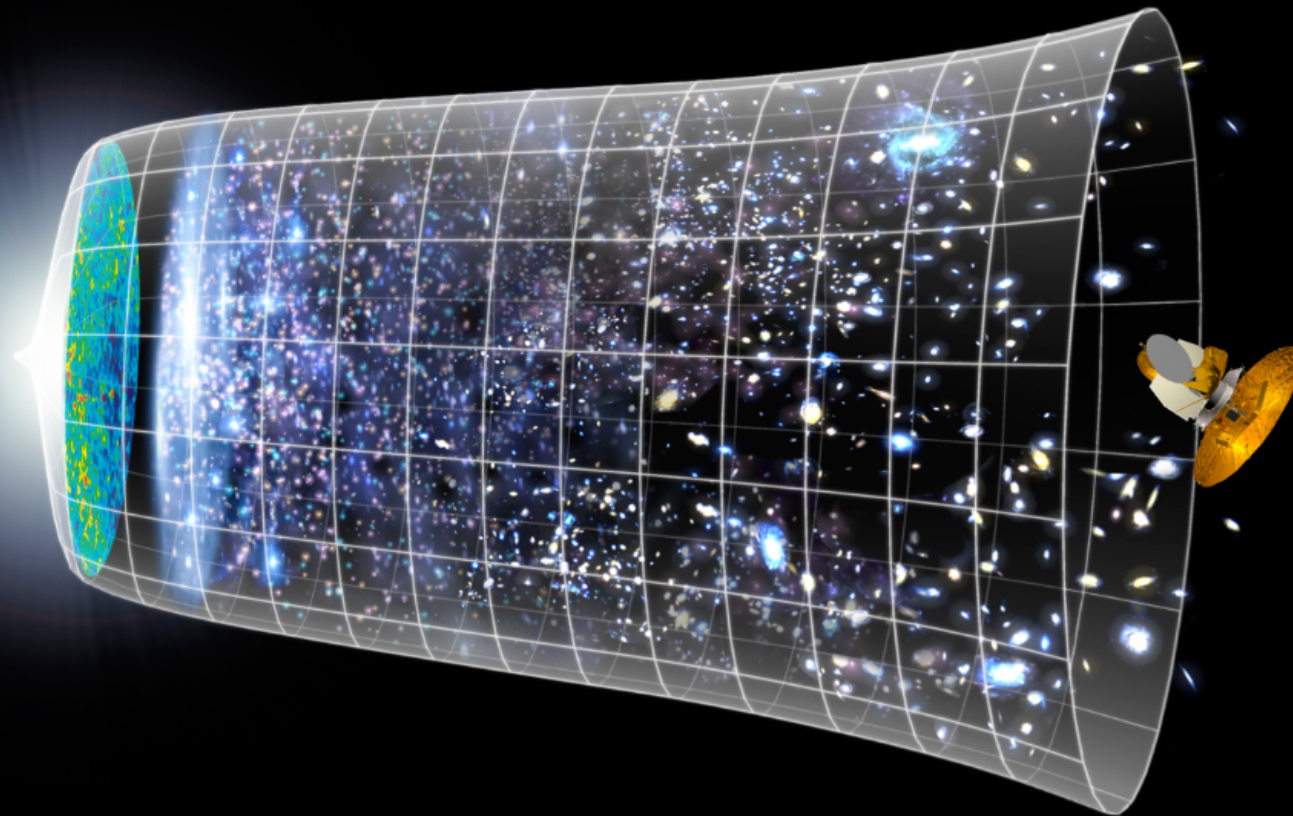


Primordial Non-Gaussianity in the CMB



Amit Yadav
Institute for Advanced Study, Princeton

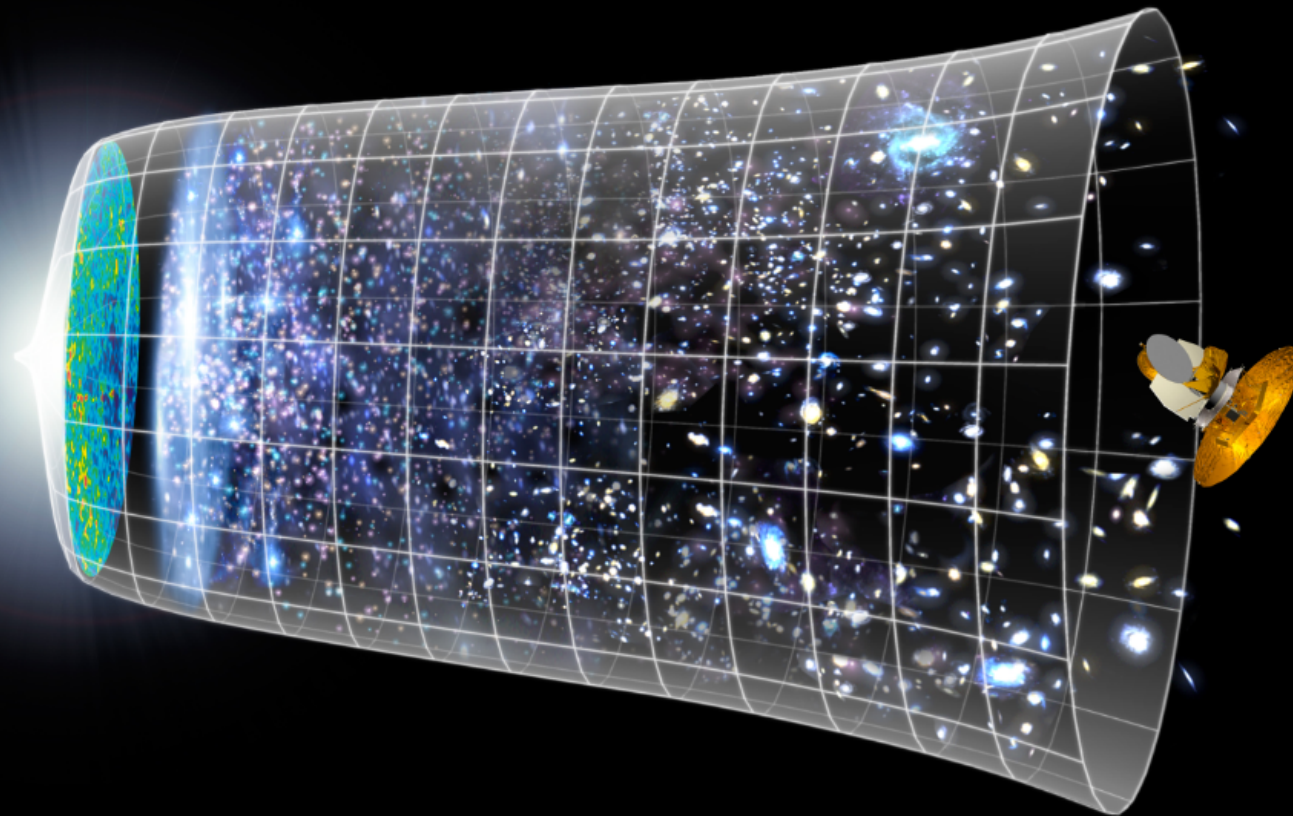
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It is well established that the observed structure originated from seed perturbations imprinted in the early universe

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The physics responsible for generating the seed perturbations is largely unknown.

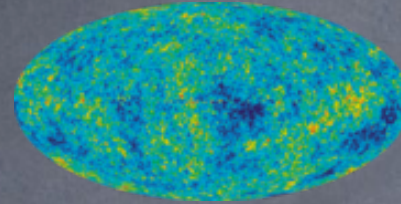
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Inflation (Exponentially Expanding Phase)

1) Scalar Fluctuations

🌀 Adiabatic

5 %



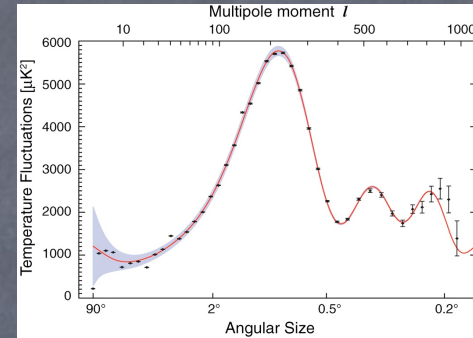
🌀 Nearly Gaussian

0.2 %

🌀 Nearly Scale Invariant

5%

2) Tensor Fluctuations (gravitational waves)

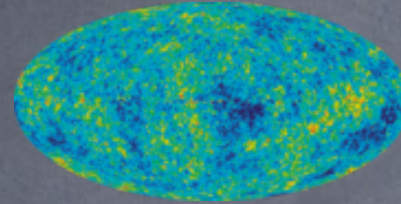


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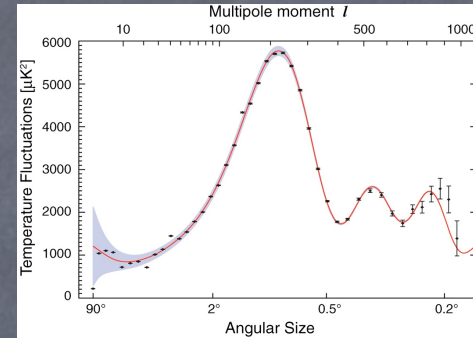
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(1) Current data still allows hundreds of inflationary models

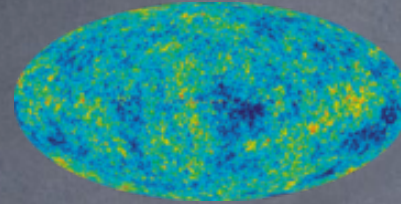
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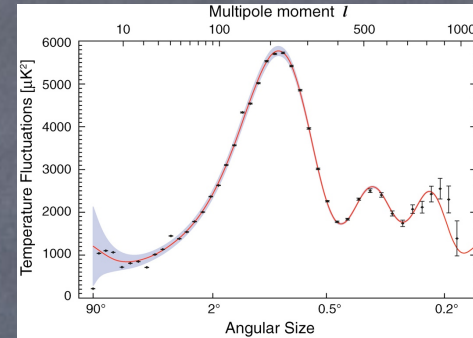


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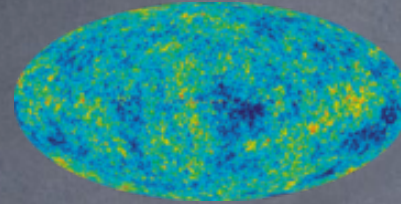
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- kinetic and potential terms
- interactions
- Energy scale

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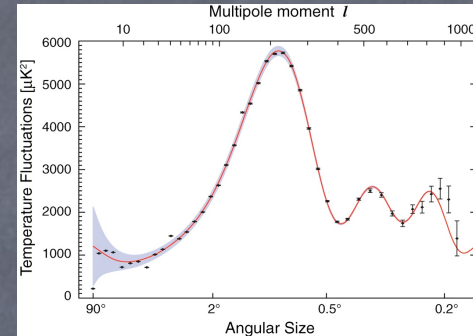


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- How many fields
- kinetic and potential terms
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- Non-Gaussianity in the CMB
- Gravitational waves (B-modes of CMB)
- Isocurvature modes

Non-Gaussianity= deviation from Gaussianity

For a Gaussian random field $\Phi(x)$

the statistical properties are completely characterized by its two point correlation function (or its power spectrum): $\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2) \rangle = \delta(k_1 + k_2)P_\Phi(k_1)$

All higher order connected correlation functions are vanishing

3-point/bispectrum $\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2)\Phi(\mathbf{k}_3) \rangle = 0$

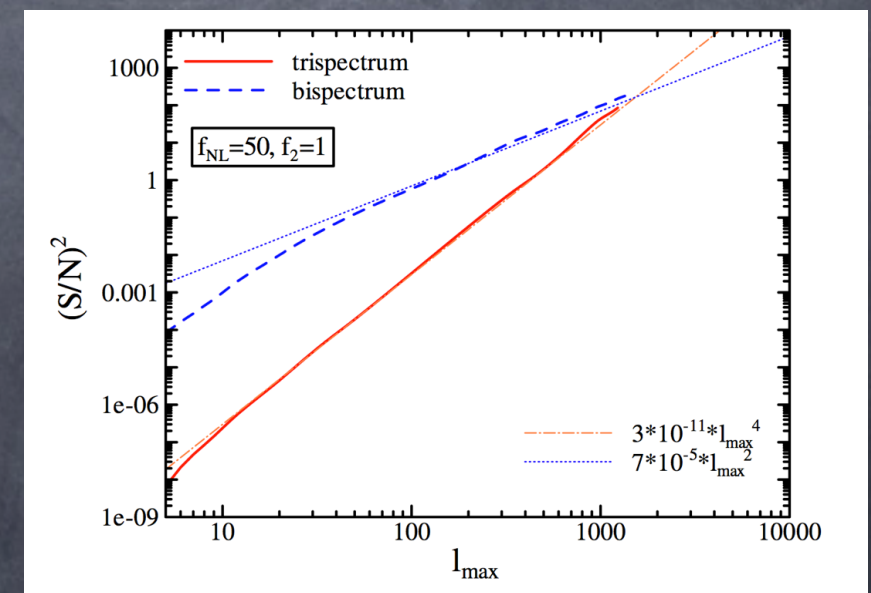
4-point/trispectrum $\langle \Phi(\mathbf{k}_1)\Phi(\mathbf{k}_2)\Phi(\mathbf{k}_3)\Phi(\mathbf{k}_4) \rangle_c = 0$

Minkowski Functionals

Wavelets

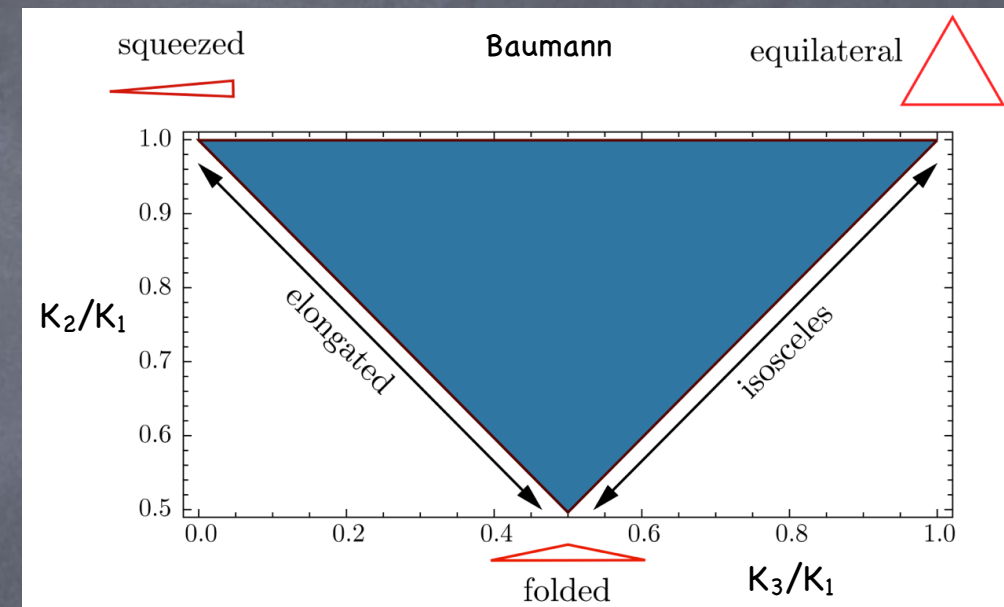
Bispectrum

- Bispectrum contains nearly all the information for local-Gaussianity
eg. Okamoto & Hu (2002), Babich (2005), Kogo & Komatsu (2006), Creminelli et al. (2007)
- For two scalar fields, it is possible to have large trispectrum
e.g. Byrnes et al. (2006), Engel et al. (2009)



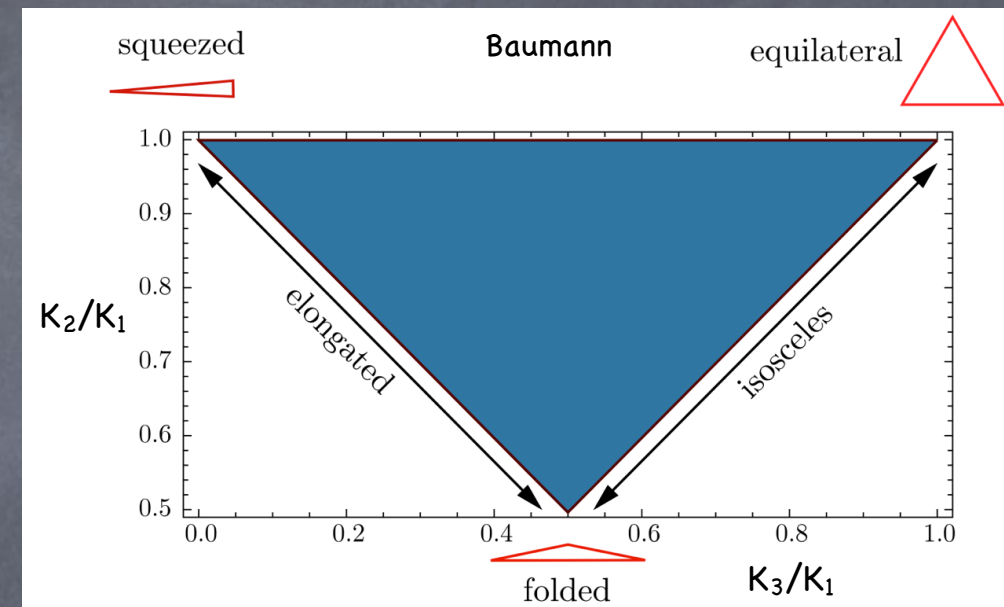
What does inflation predict?

$$\begin{aligned} B_{\Phi}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3) &= \langle \Phi(\mathbf{k}_1) \Phi(\mathbf{k}_2) \Phi(\mathbf{k}_3) \rangle \\ &= f_{NL} \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) F(k_1, k_2, k_3) \end{aligned}$$



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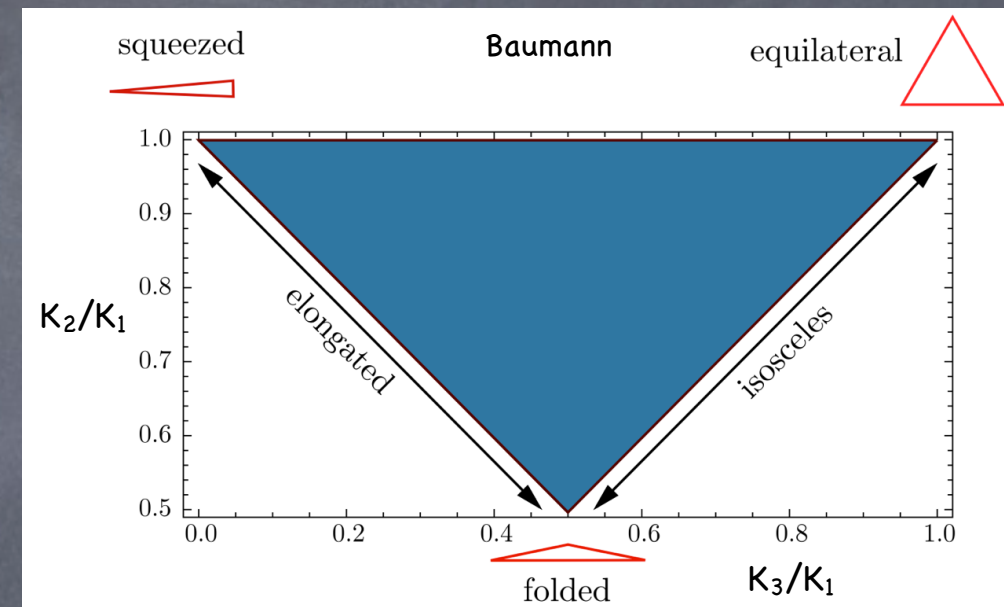
Maldacena (2003)

- Single field slow-roll inflation cannot produce large nonG

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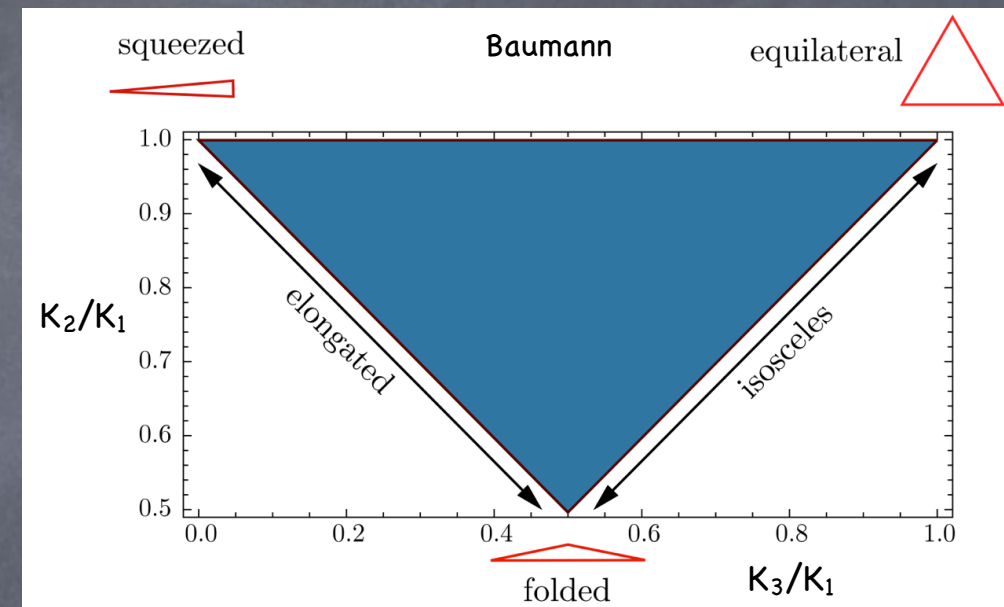
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- Single field inflation cannot produce squeezed nonG

$$f_{NL} = (1 - n_s) \quad \text{Maldacena (2003); Creminelli \& Zaldarriaga (2004)}$$

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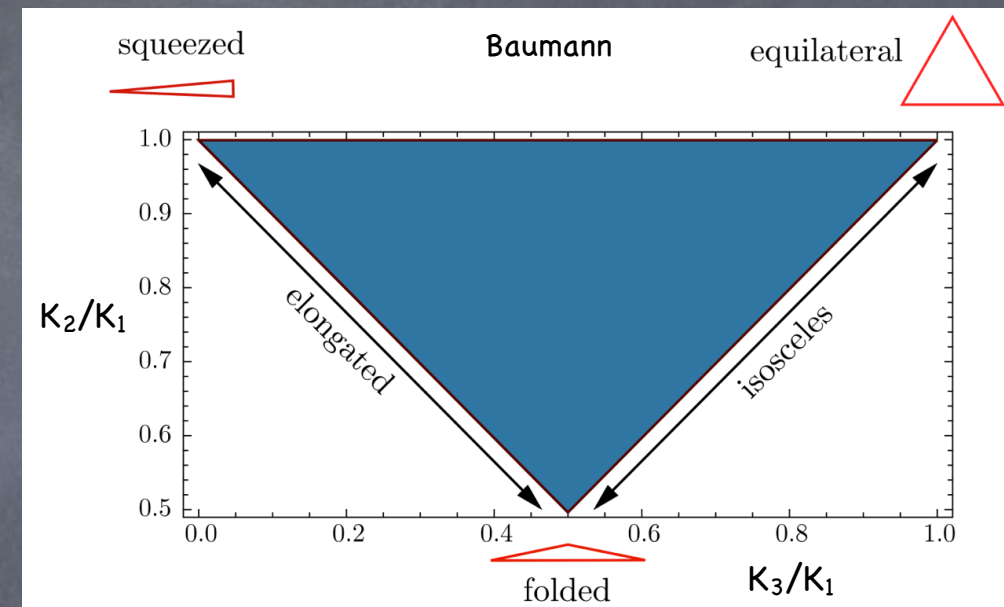
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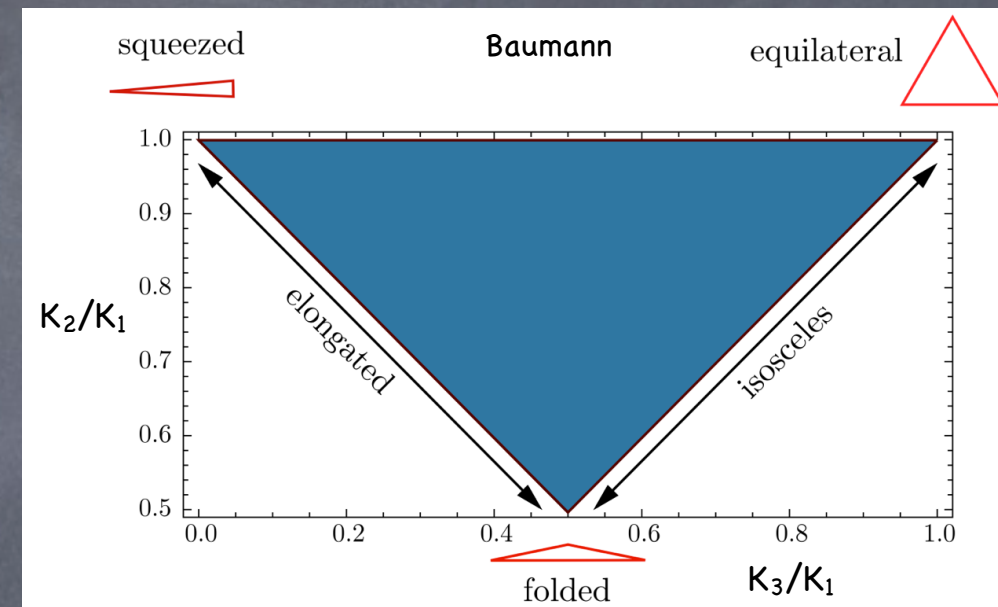
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4. Initial vacuum state

$f_{NL} \sim 0.05$ canonical inflation (single field, Maldacena 2003, Acquaviva et al. 2003)

$f_{NL} \sim 100$ DBI inflation (Alishahiha et al. 2004) Equilateral

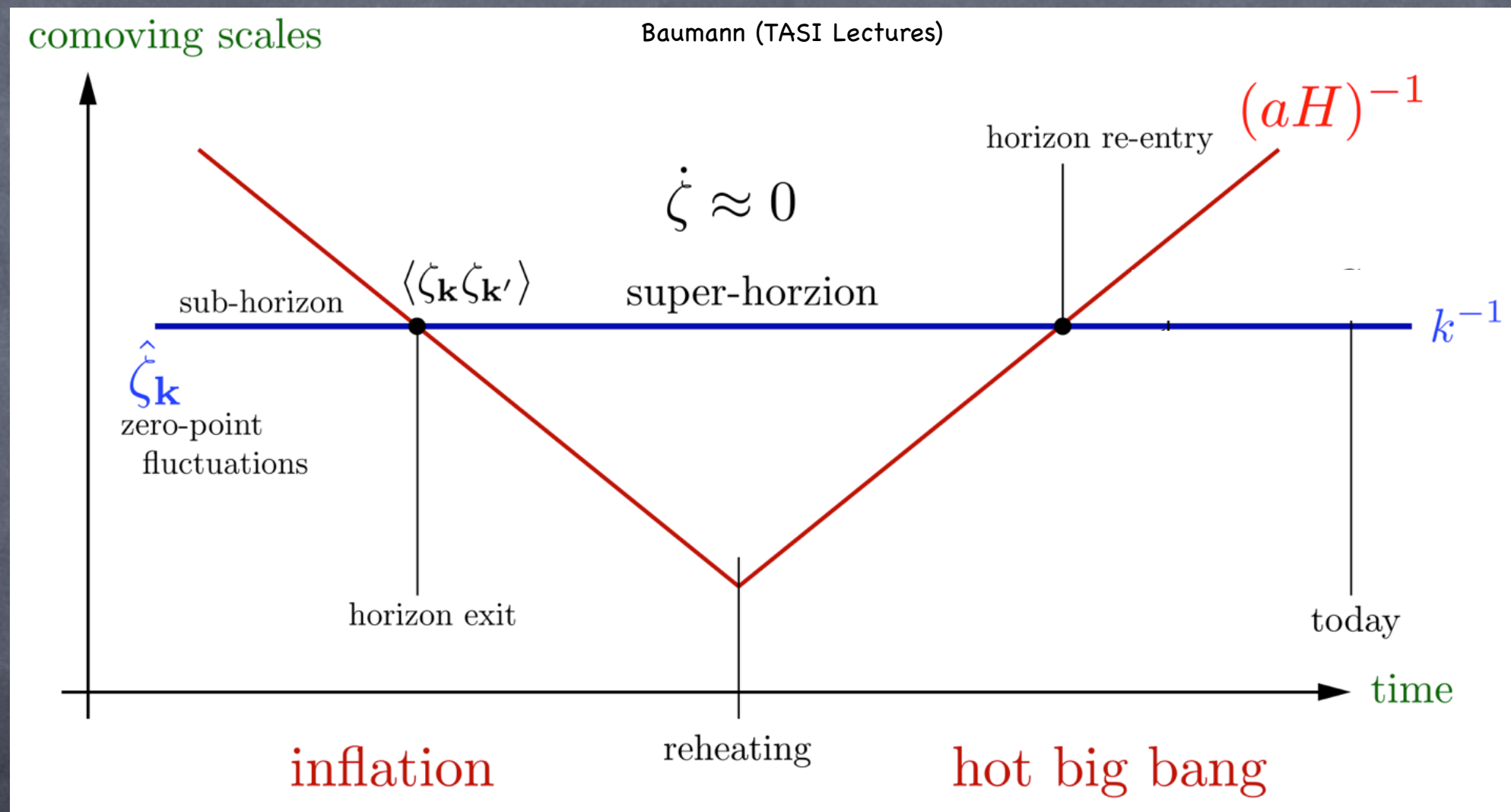
$f_{NL} > 10$ curvaton models (Lyth et al. 2003) Local

$f_{NL} \sim 100$ ghost inflation (Arkani-Hamed et al. 2004) Equilateral

$f_{NL} \sim 20-100$ Cyclic models (Creminelli et al. 07, Buchbinder et. al 07, Koyama et. al 07)

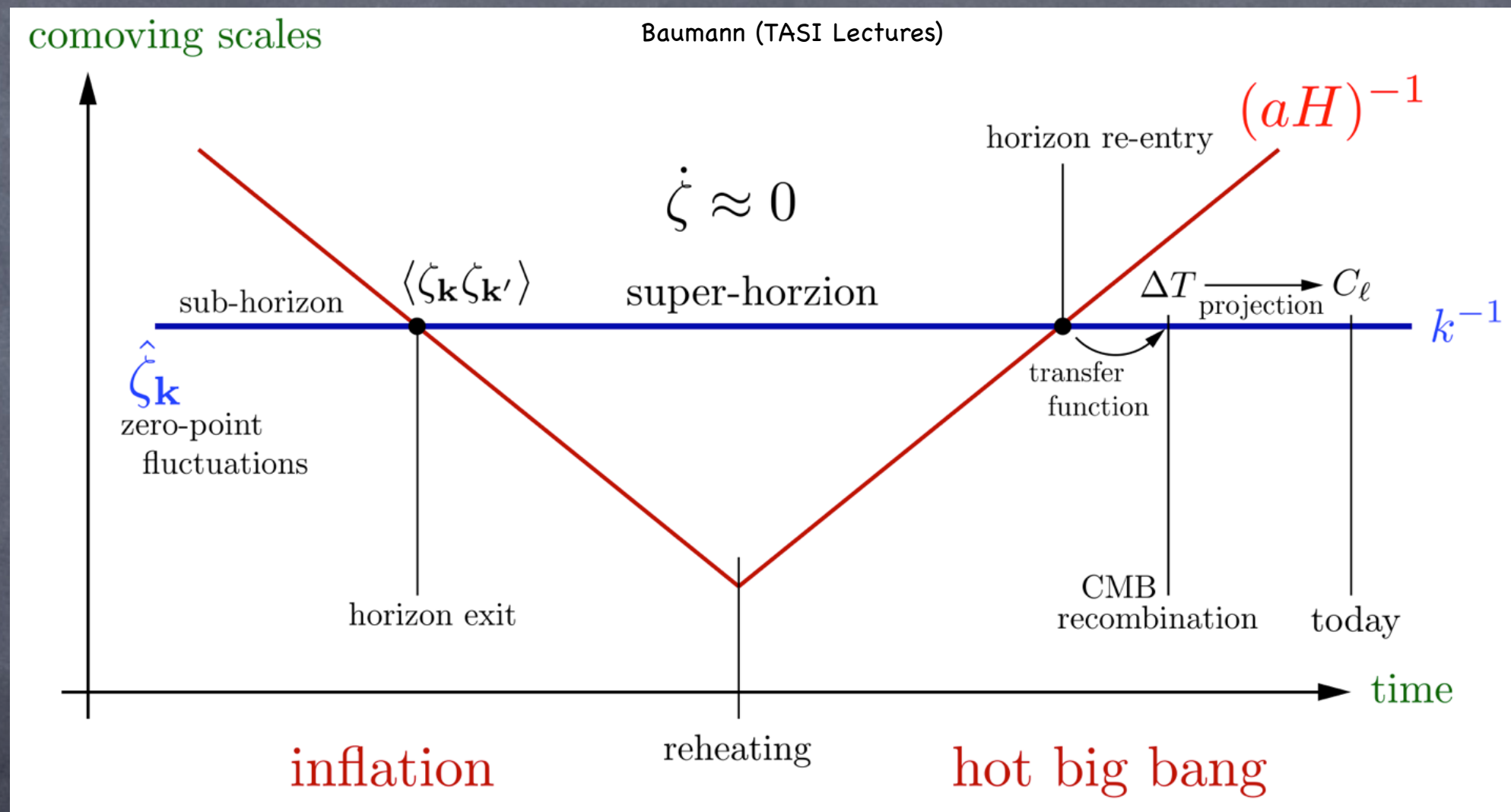
How is Cosmic Microwave Background
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- Gauge invariant
- Conserve outside the horizon

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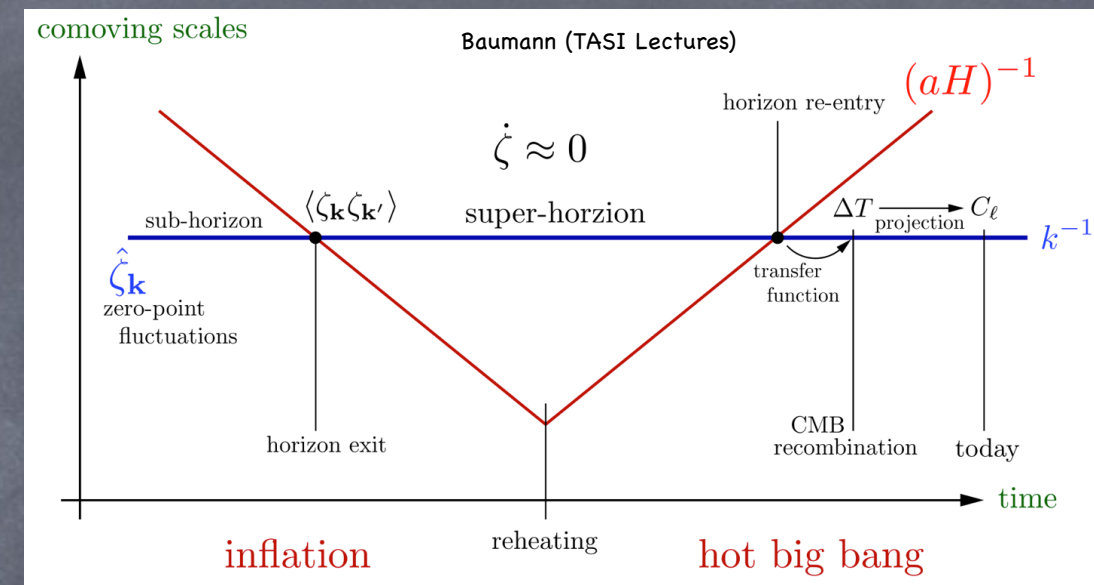


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CMB Bispectrum

Math people this slide
Picture people next slide

$$a_{\ell m} = 4\pi(-i)^\ell \int \frac{d^3 k}{(2\pi)^3} \Phi(\mathbf{k}) g_\ell(k) Y_{\ell m}^*(\hat{k})$$

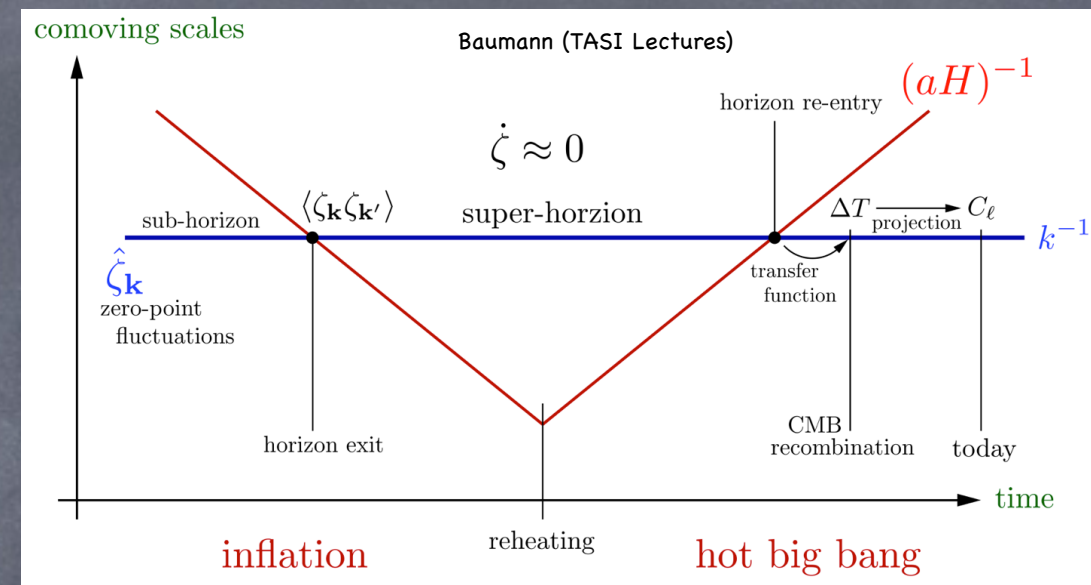


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Transfer function
CMBFast, CAMB



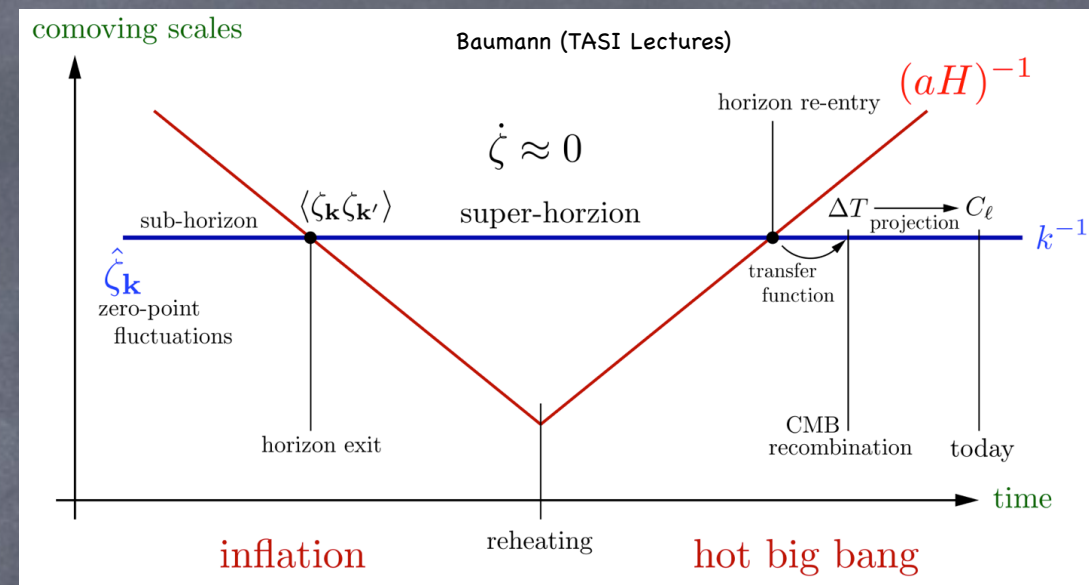
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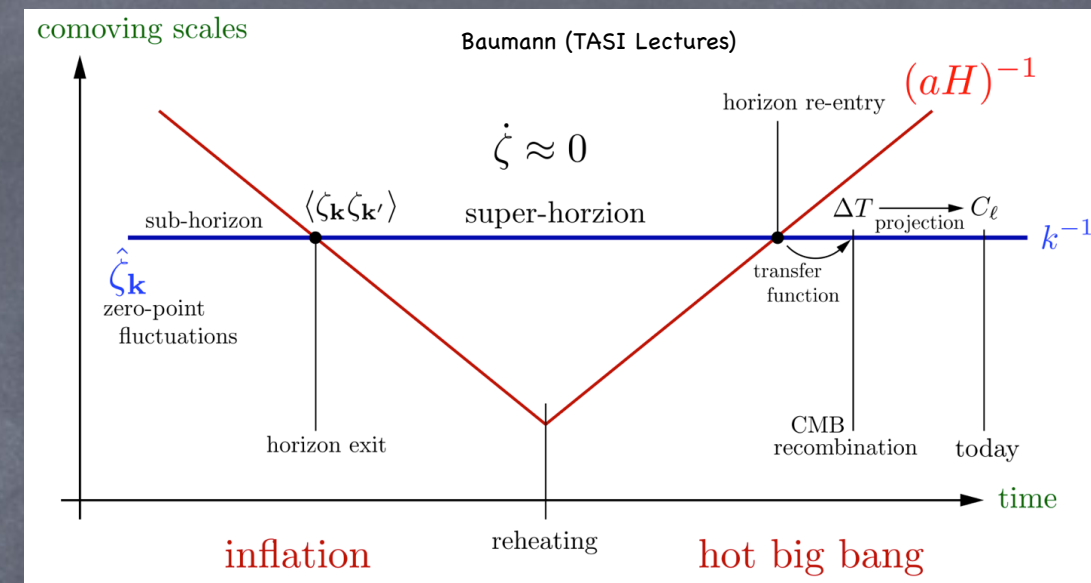
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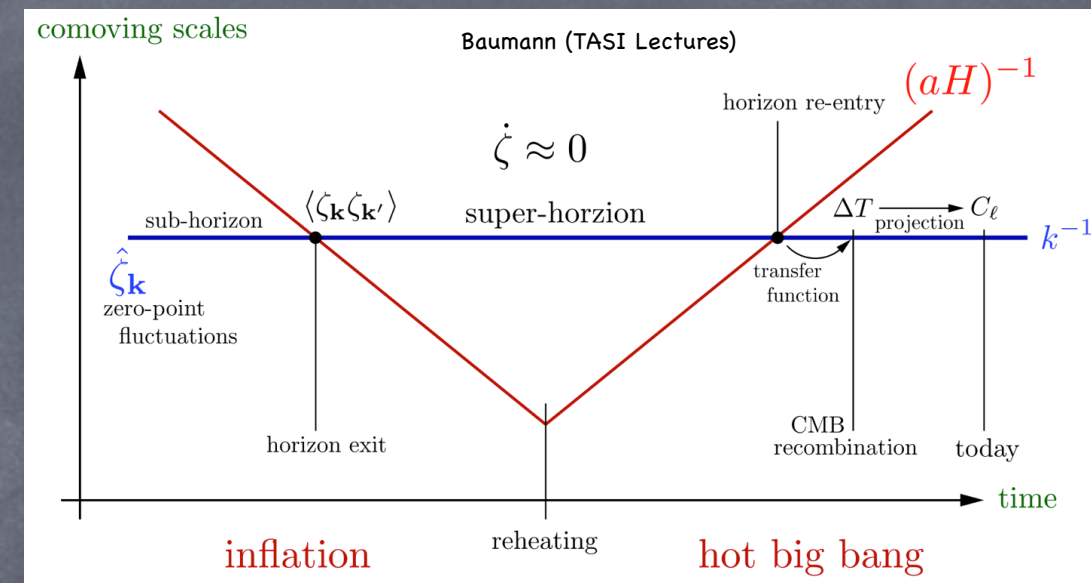
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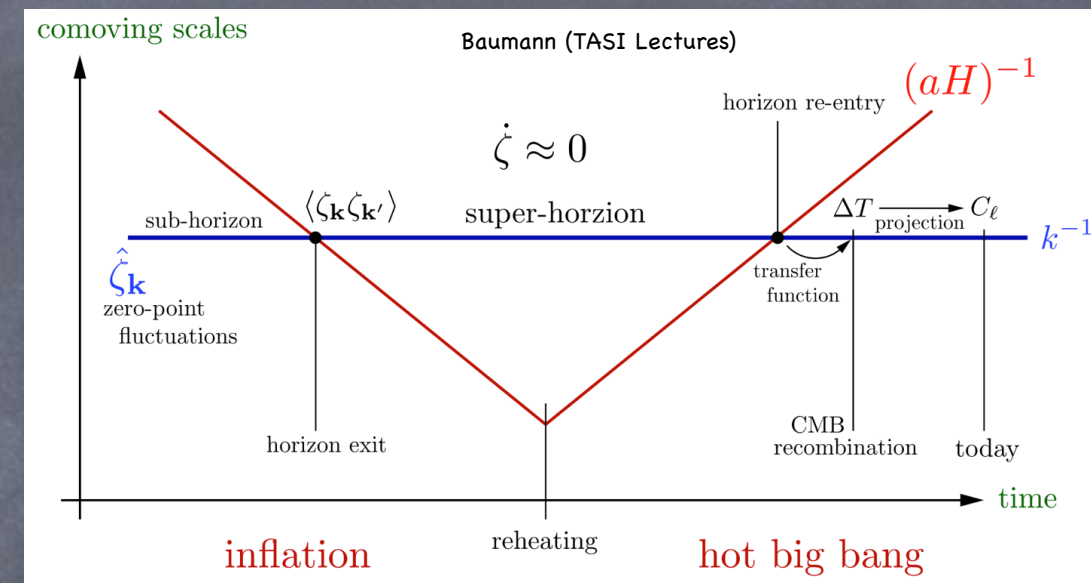
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Primordial
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$N_{pix}^{5/2}$

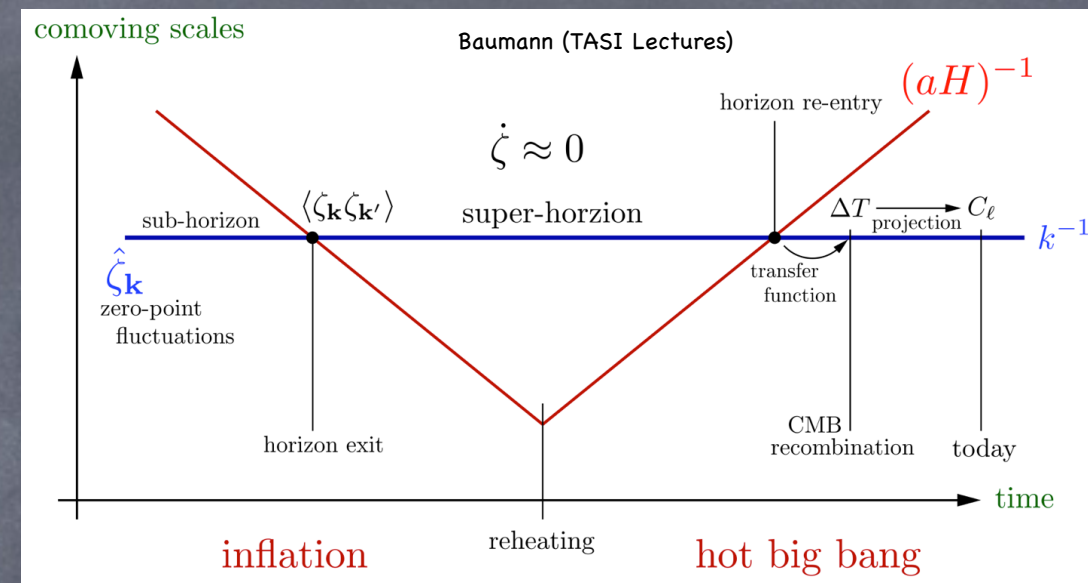
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Primordial
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$$F(k_1, k_2, k_3) = f_1(k_1) f_2(k_2) f_3(k_3)$$

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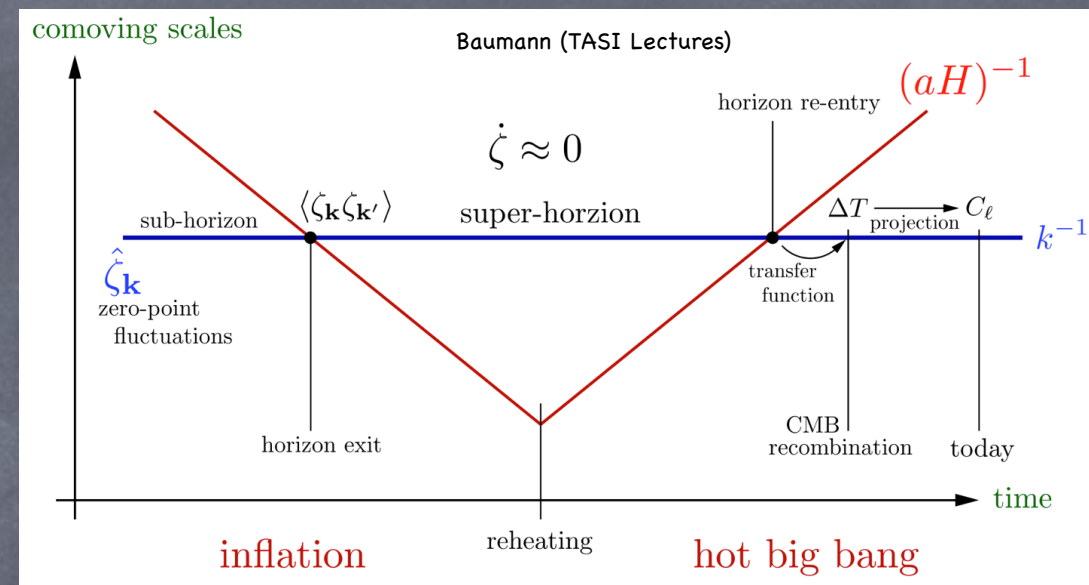
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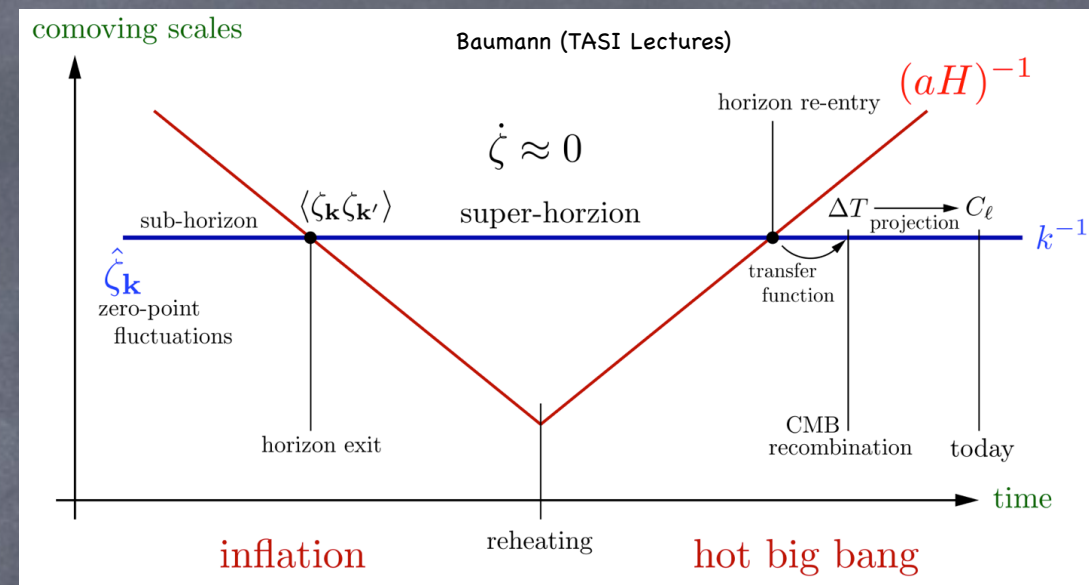
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$$\hat{f}_{NL} = \frac{1}{N} \sum_{\ell_i m_i} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \frac{B_{\ell_1 \ell_2 \ell_3}}{C_{\ell_1} C_{\ell_2} C_{\ell_3}} a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3}$$



$$\hat{f}_{NL} = \frac{1}{N} \cdot \sum_{l_i m_i} \int d^2 \hat{n} Y_{l_1 m_1}(\hat{n}) Y_{l_2 m_2}(\hat{n}) Y_{l_3 m_3}(\hat{n}) \int_0^\infty r^2 dr j_{l_1}(k_1 r) j_{l_2}(k_2 r) j_{l_3}(k_3 r) C_{l_1}^{-1} C_{l_2}^{-1} C_{l_3}^{-1}$$

$$\int \frac{2k_1^2 dk_1}{\pi} \frac{2k_2^2 dk_2}{\pi} \frac{2k_3^2 dk_3}{\pi} F(k_1, k_2, k_3) \Delta_{l_1}^T(k_1) \Delta_{l_2}^T(k_2) \Delta_{l_3}^T(k_3) a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3} ,$$

$N_{pix}^{5/2}$

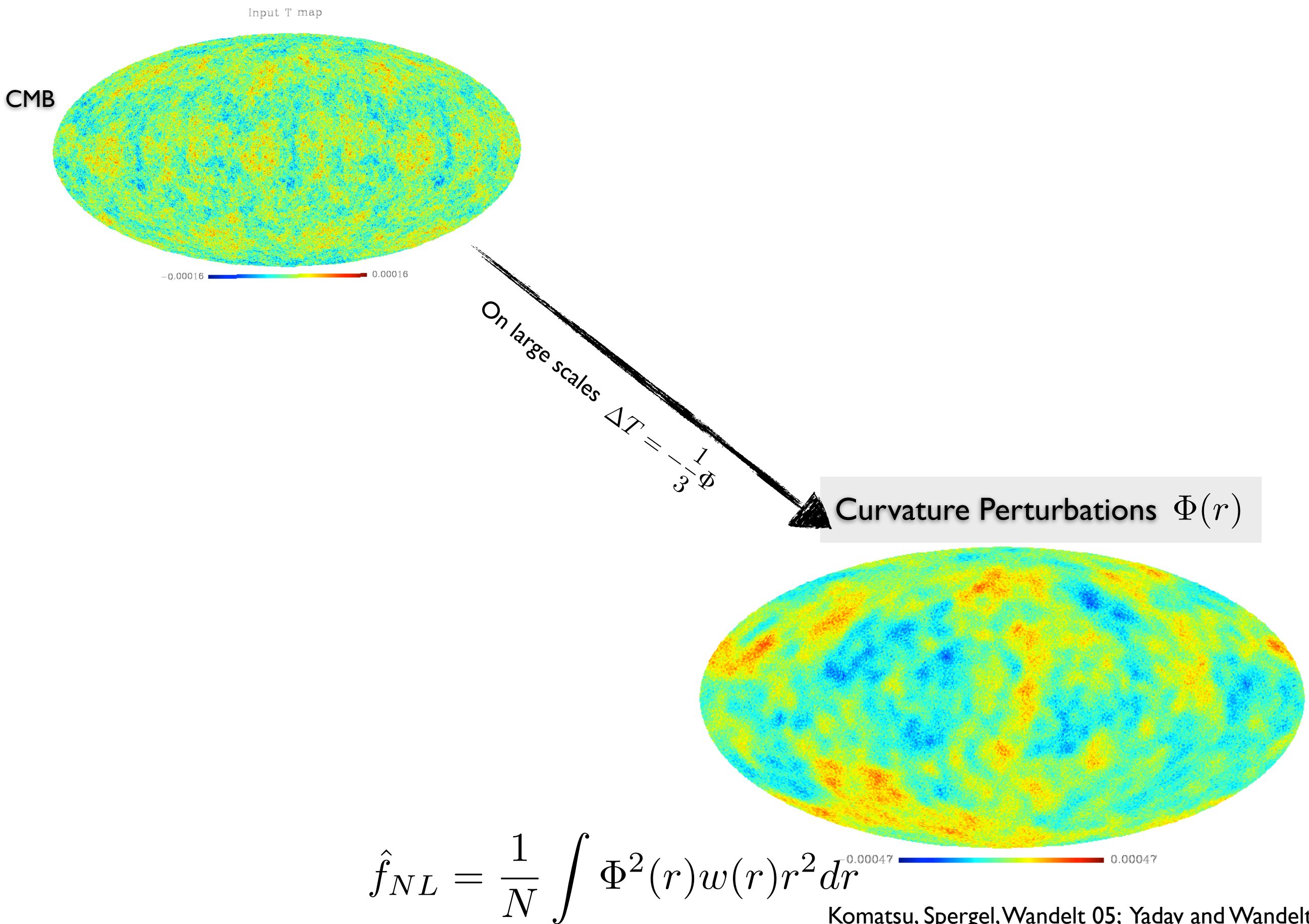
Primordial
Bispectrum

$$F(k_1, k_2, k_3) = f_1(k_1) f_2(k_2) f_3(k_3)$$

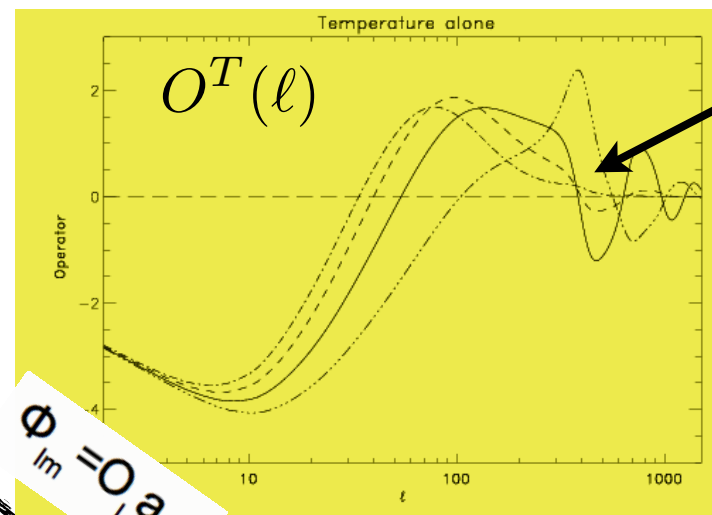
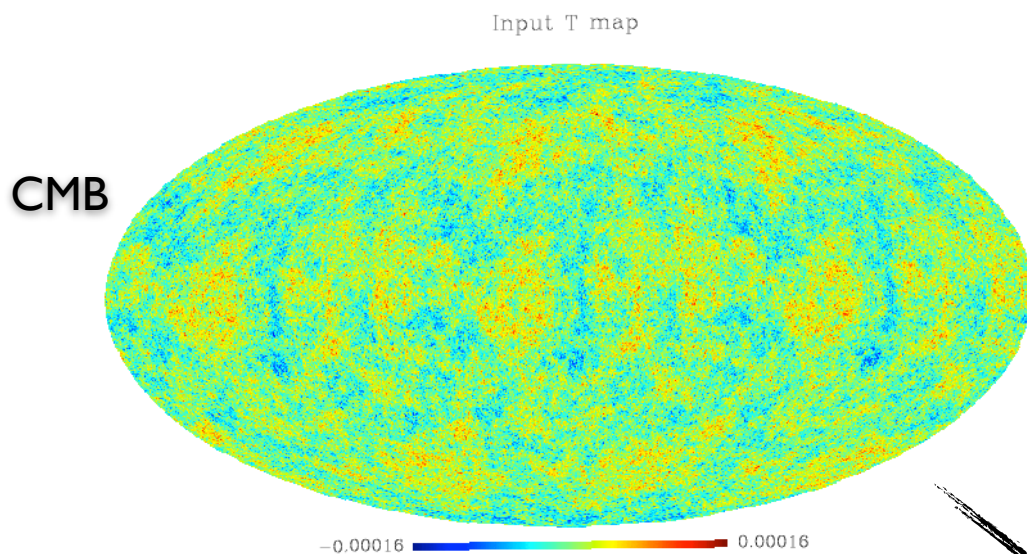
$$\hat{f}_{NL} = \frac{1}{N} \cdot \int d^2 \hat{n} \int_0^\infty r^2 dr \prod_{i=1}^3 \sum_{l_i m_i} \int \frac{2k^2 dk}{\pi} j_{l_i}(kr) f_i(k) \Delta_{l_i}^T(k) C_{l_i}^{-1} a_{l_i m_i} Y_{l_i m_i}(\hat{n})$$

$N_{pix}^{3/2}$

Primordial non-Gaussianity from CMB



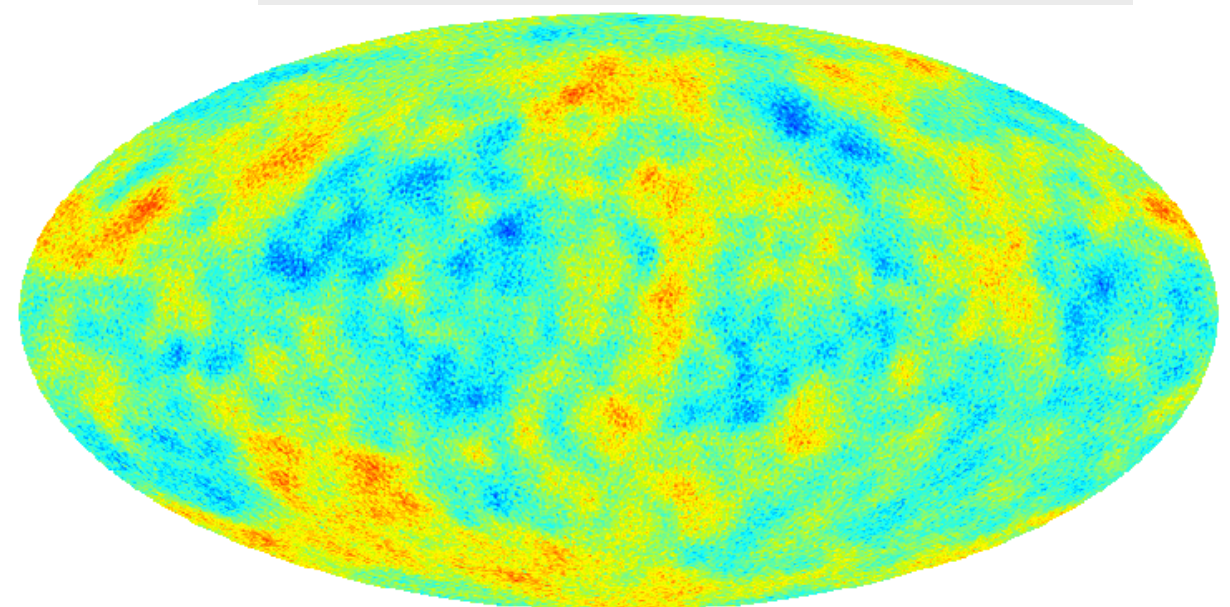
Primordial non-Gaussianity from CMB



Filter goes to zero
(reconstruction fails)

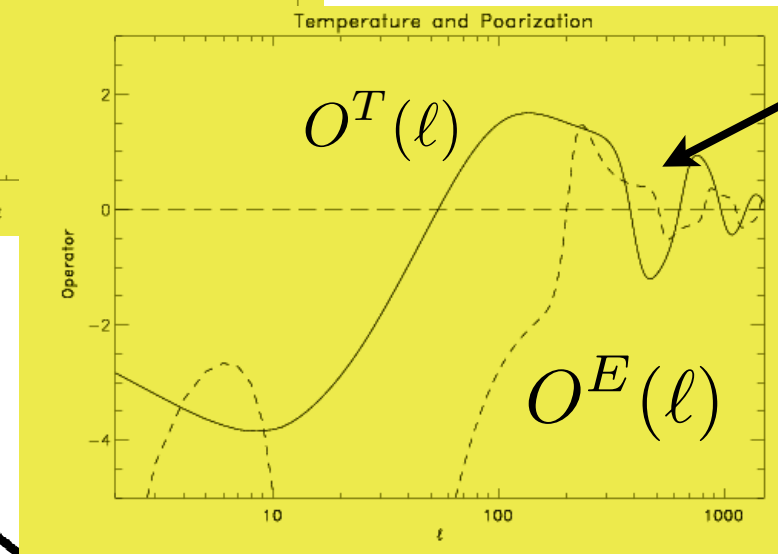
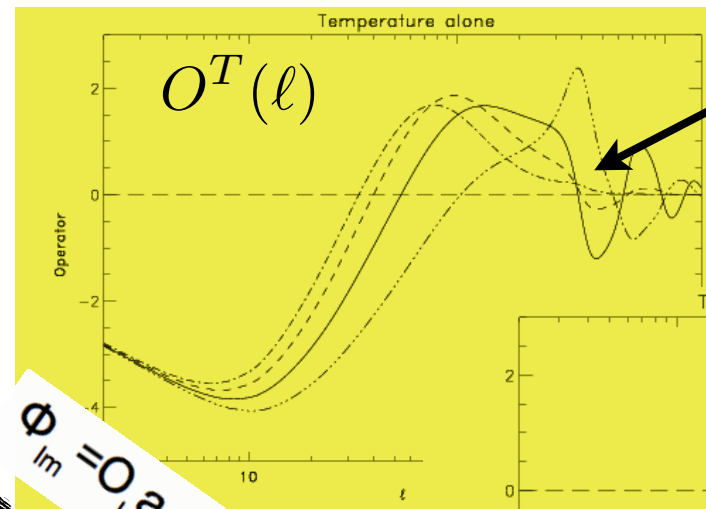
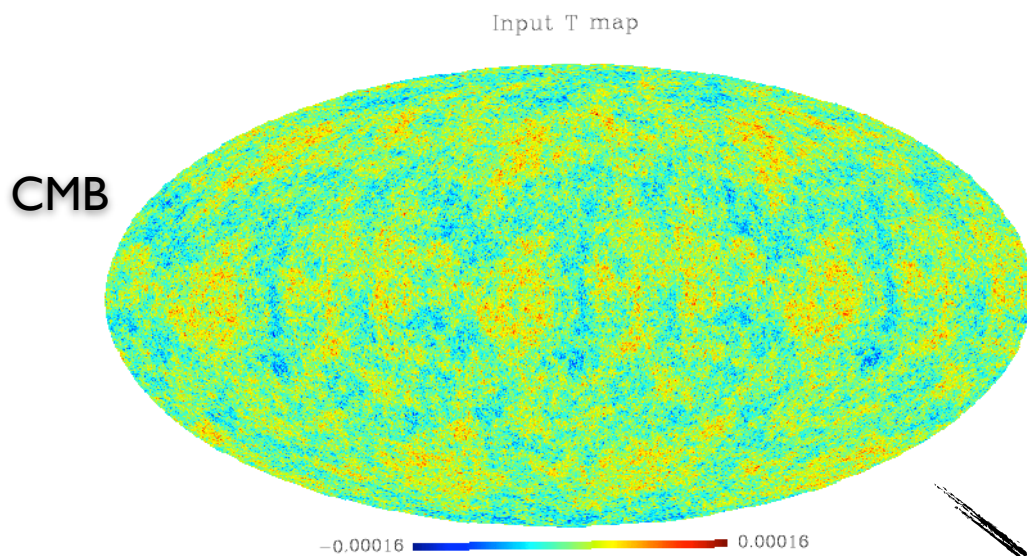
On large scales $\Delta T = -\frac{1}{3}\Phi$

Curvature Perturbations $\Phi(r)$



$$\hat{f}_{NL} = \frac{1}{N} \int \Phi^2(r) w(r) r^2 dr$$

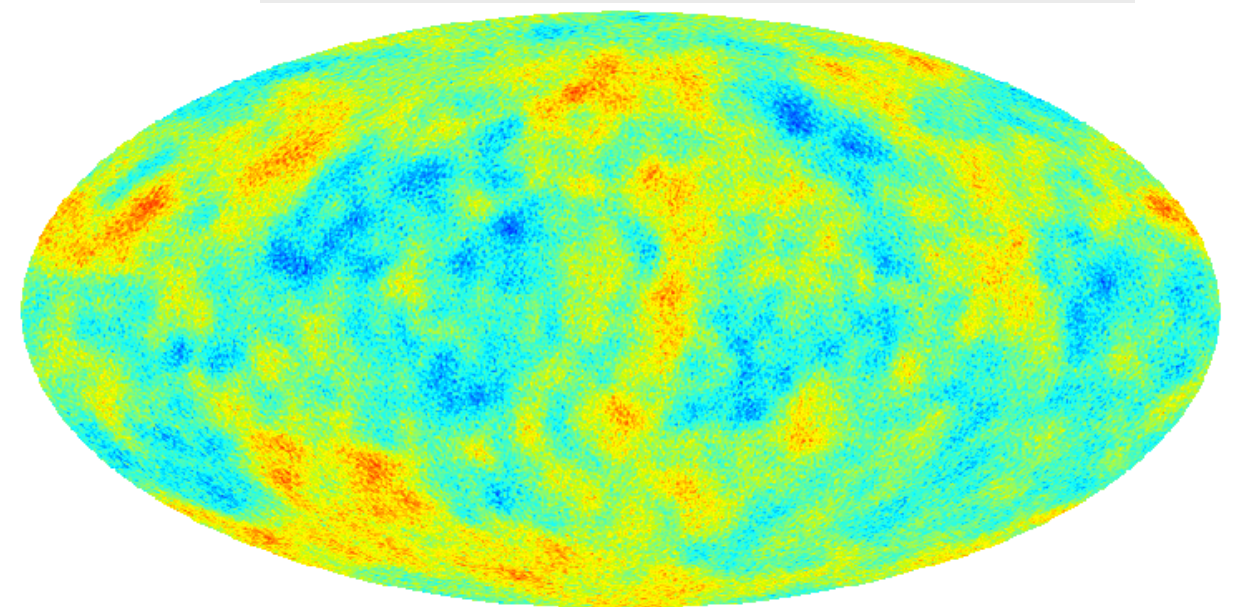
Primordial non-Gaussianity from CMB



On large scales

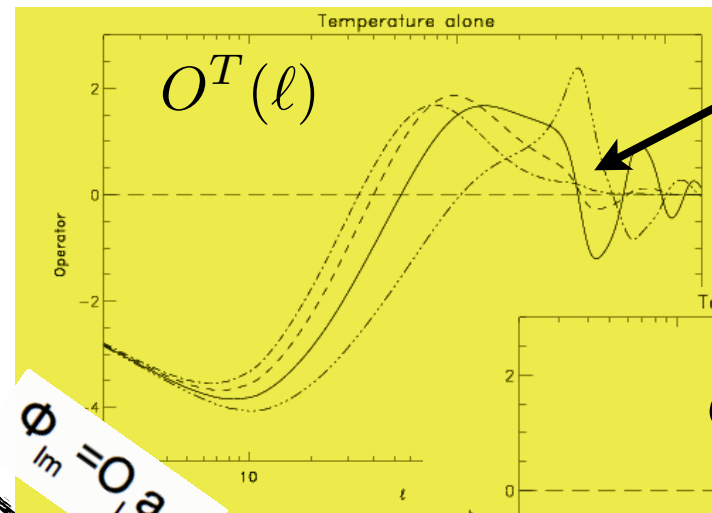
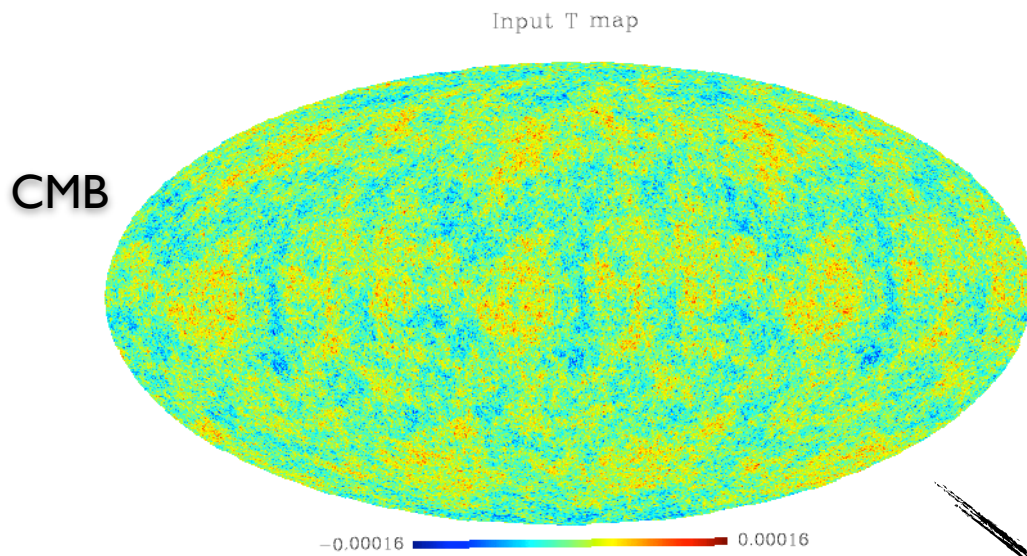
$$\Delta T = -\frac{1}{3} \Phi$$

Curvature Perturbations $\Phi(r)$

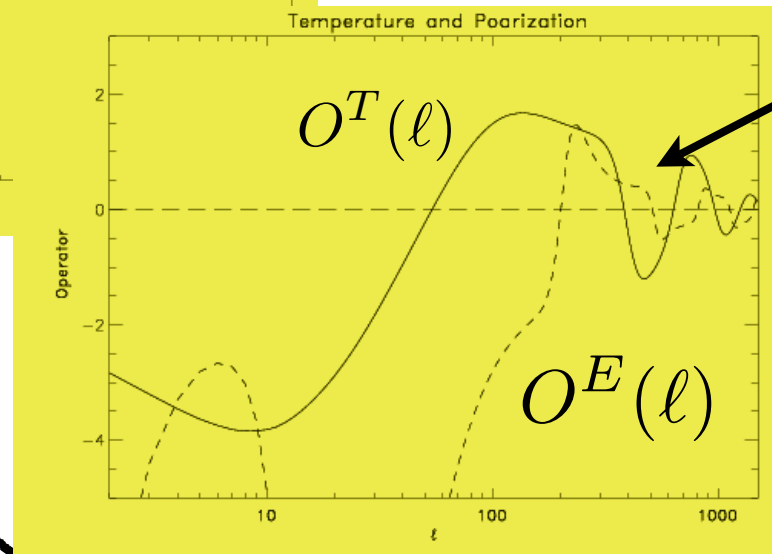


$$\hat{f}_{NL} = \frac{1}{N} \int \Phi^2(r) w(r) r^2 dr$$

Primordial non-Gaussianity from CMB



Filter goes to zero
(reconstruction fails)

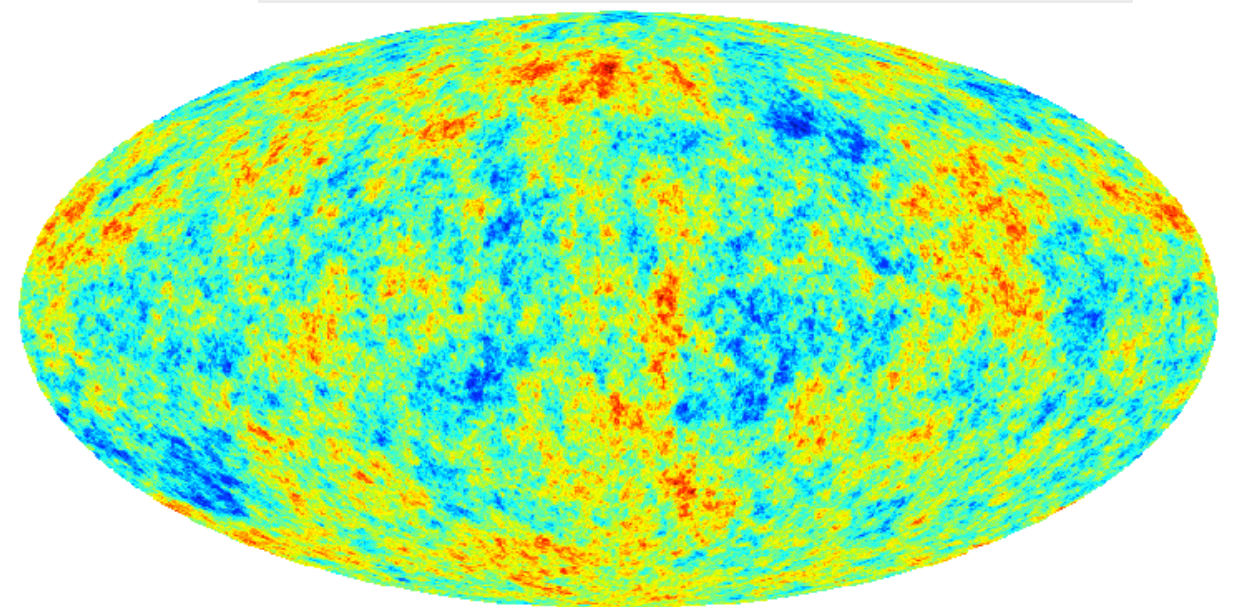


Out of phase

On large scales $\Delta T = -\frac{1}{3}\Phi$

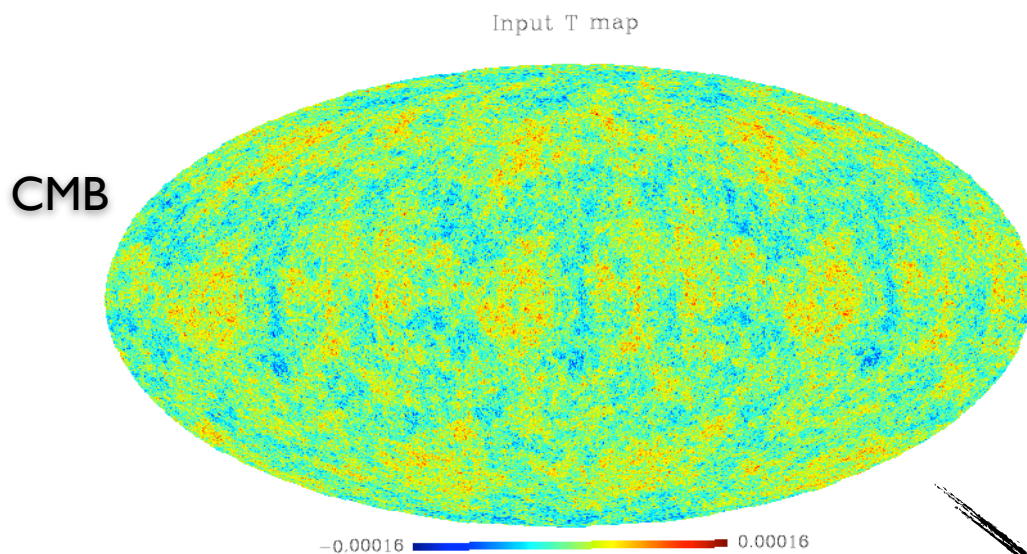
$$\Phi_{lm} = \frac{1}{l(l+1)} a_{lm}$$

Curvature Perturbations $\Phi(r)$

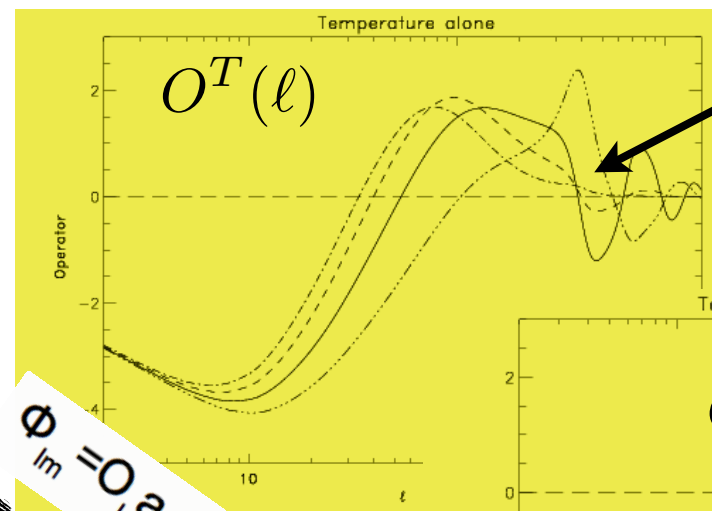


$$\hat{f}_{NL} = \frac{1}{N} \int \Phi^2(r) w(r) r^2 dr$$

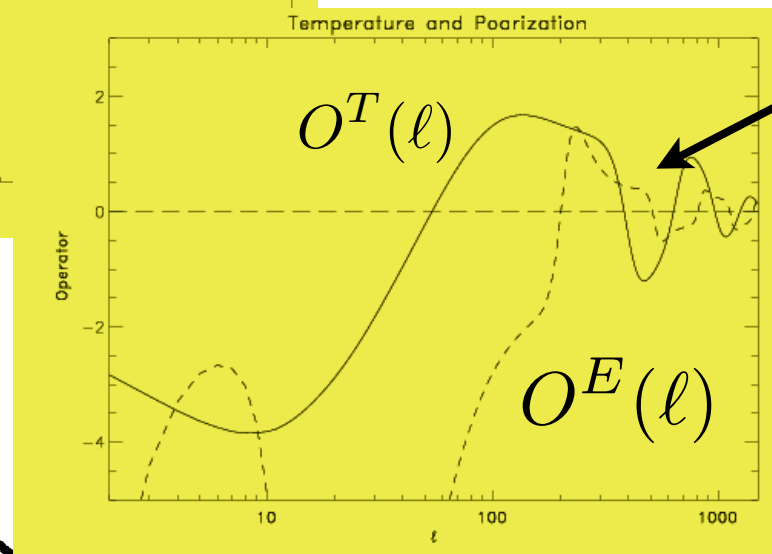
Primordial non-Gaussianity from CMB



Curvature Perturbations Movie $\Phi(r)$



Filter goes to zero
(reconstruction fails)

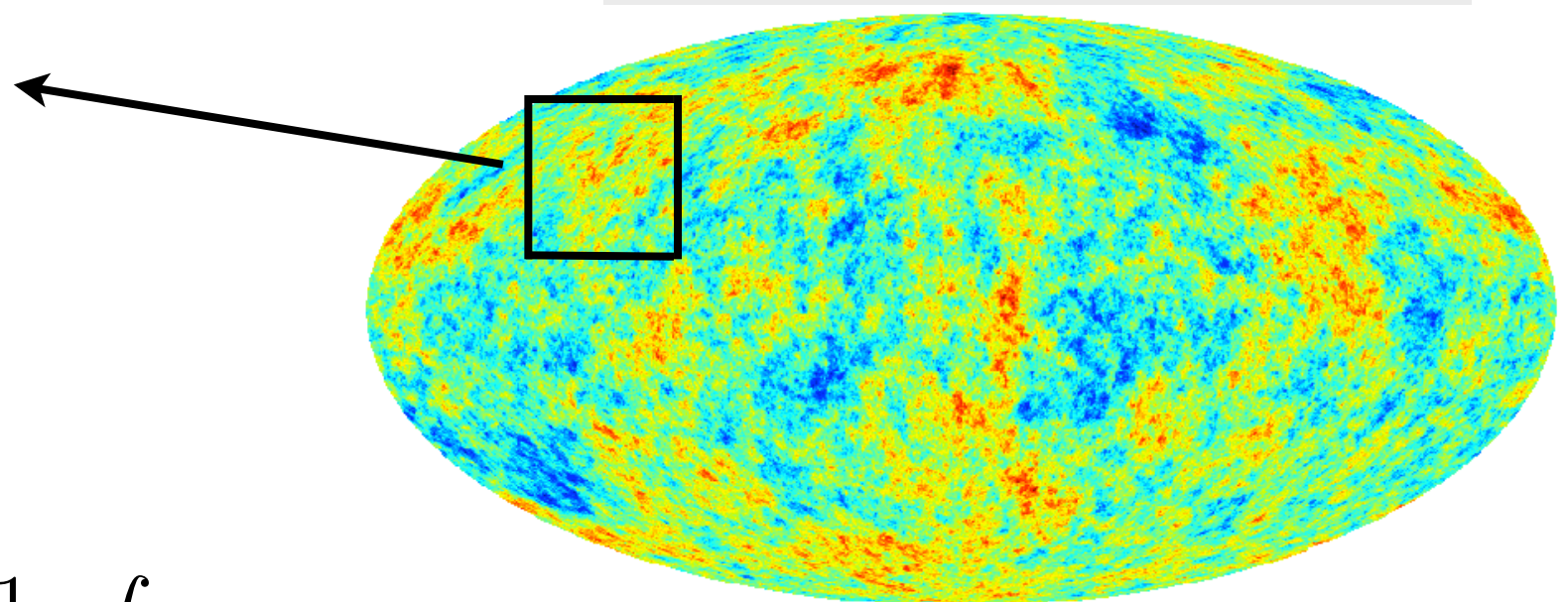


Out of phase

On large scales

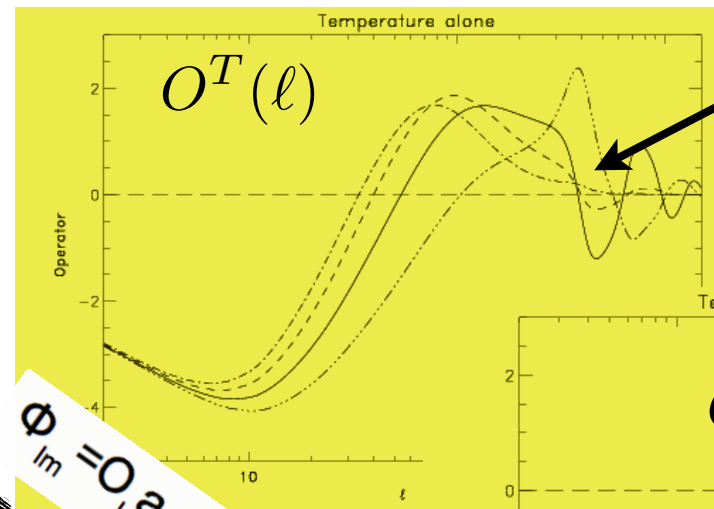
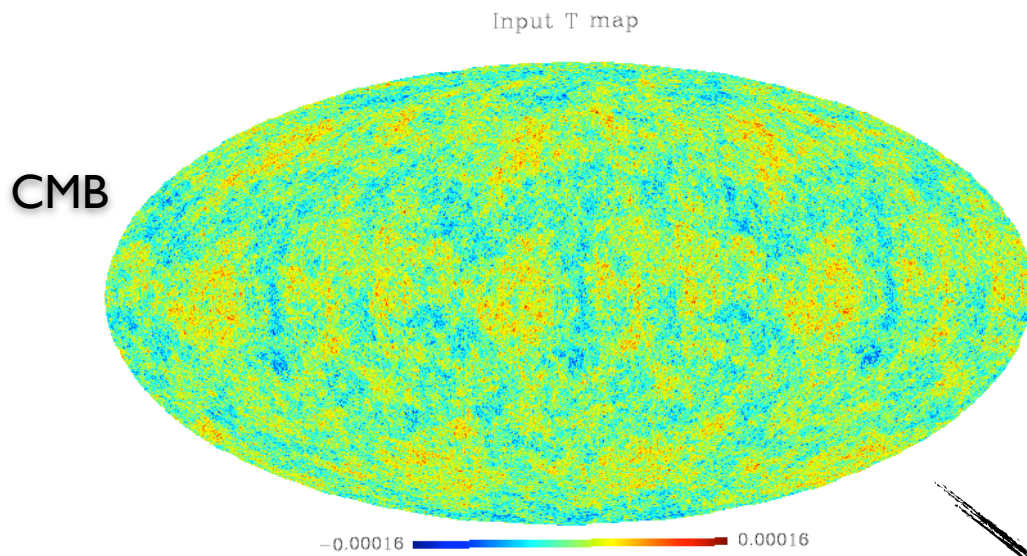
$$\Delta T = -\frac{1}{3} \Phi$$

Curvature Perturbations $\Phi(r)$

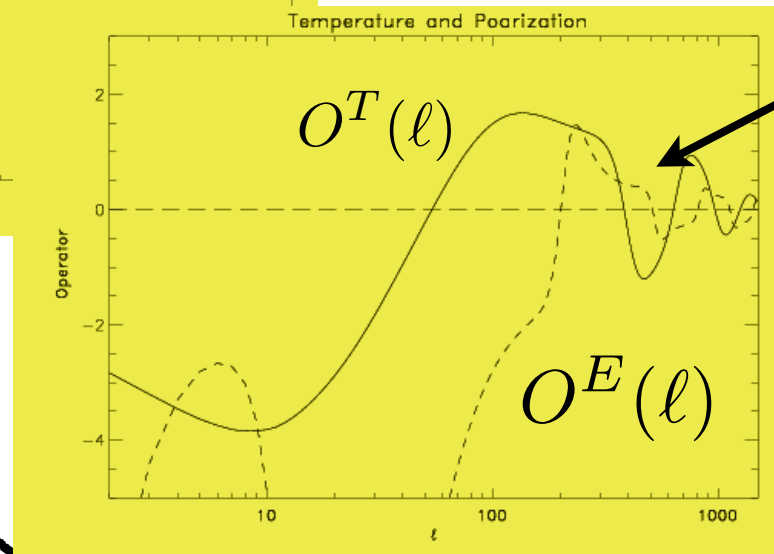


$$\hat{f}_{NL} = \frac{1}{N} \int \Phi^2(r) w(r) r^2 dr$$

Primordial non-Gaussianity from CMB

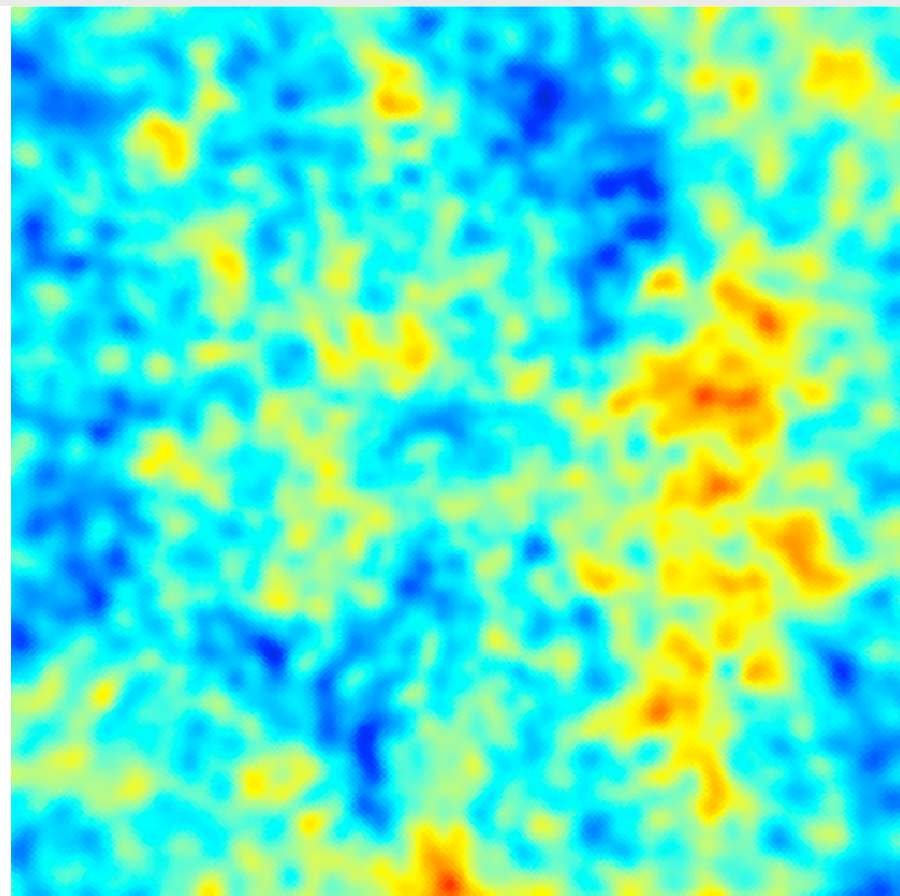


Filter goes to zero
(reconstruction fails)



Out of phase

Curvature Perturbations Movie $\Phi(r)$

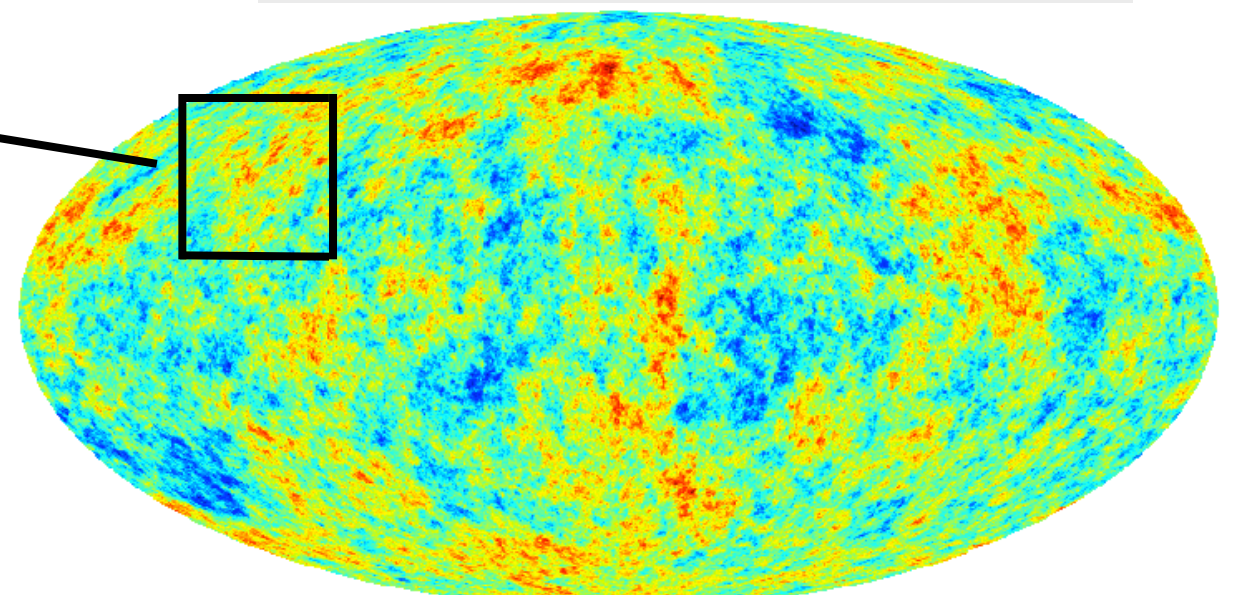


(0.0, 0.0) Galactic

On large scales

$$\Delta T = -\frac{1}{3}\Phi$$

Curvature Perturbations $\Phi(r)$

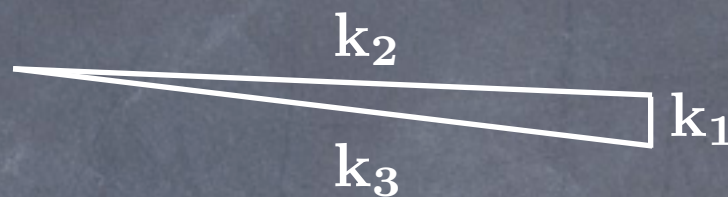
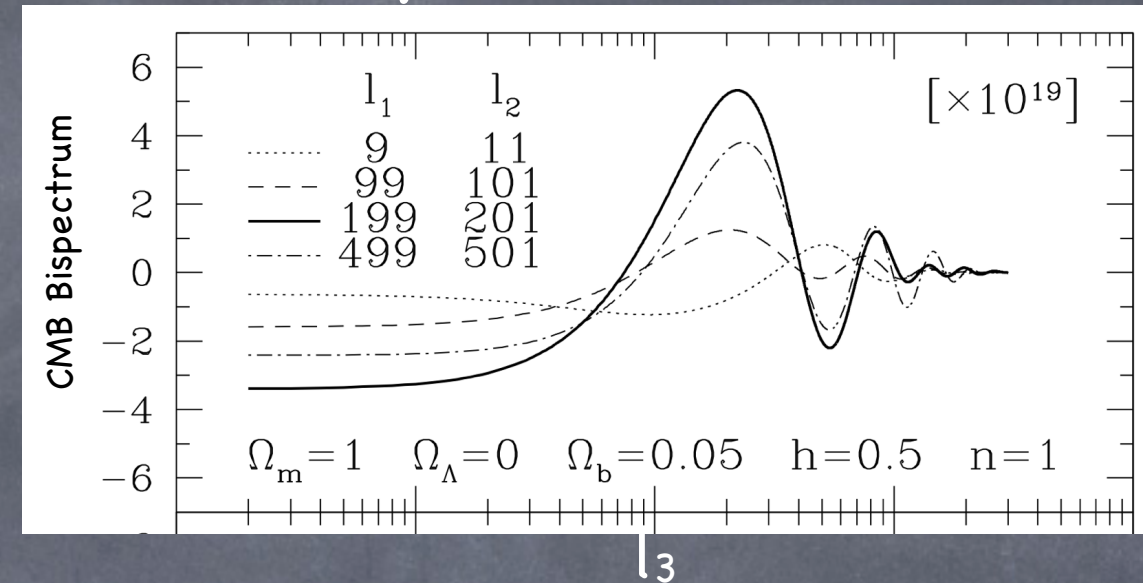


$$\hat{f}_{NL} = \frac{1}{N} \int \Phi^2(r) w(r) r^2 dr$$

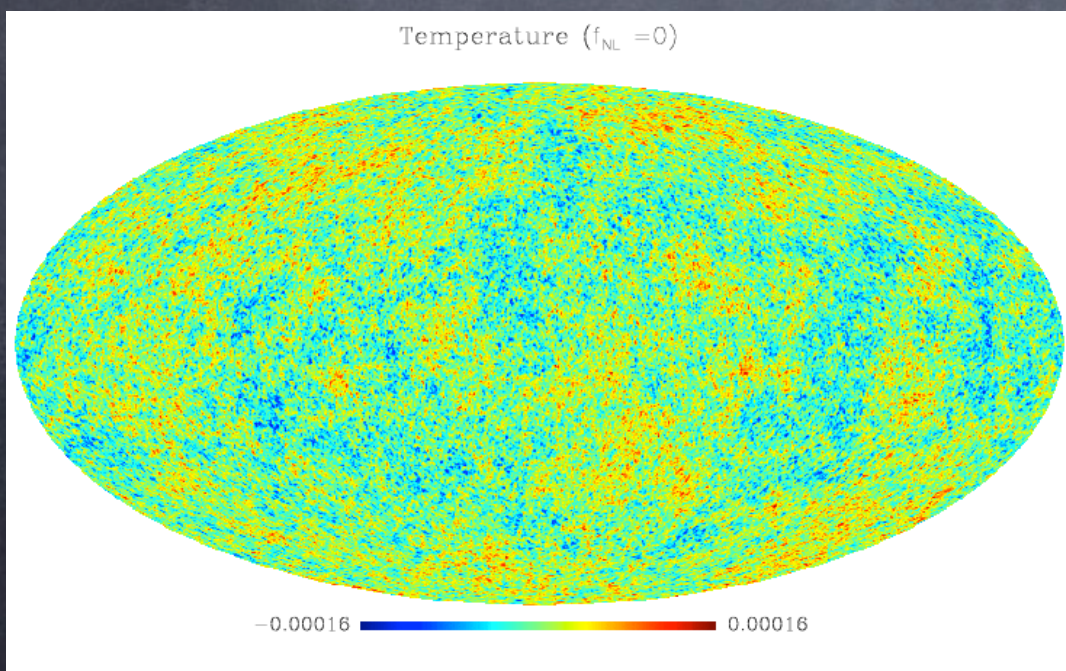
Squeezed Bispectrum/Local Non-Gaussianity

$$\Phi(x) = \Phi_G(x) + f_{NL}\Phi_G^2(x)$$

$$F(k_1, k_2, k_3) = \Delta_\Phi^2 \cdot \left(\frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3} \right)$$



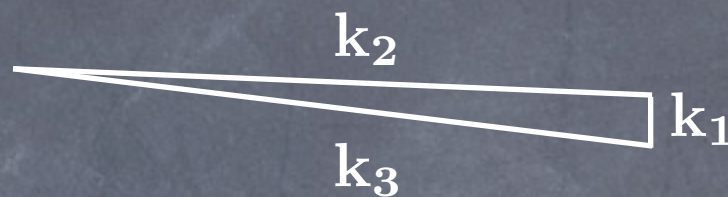
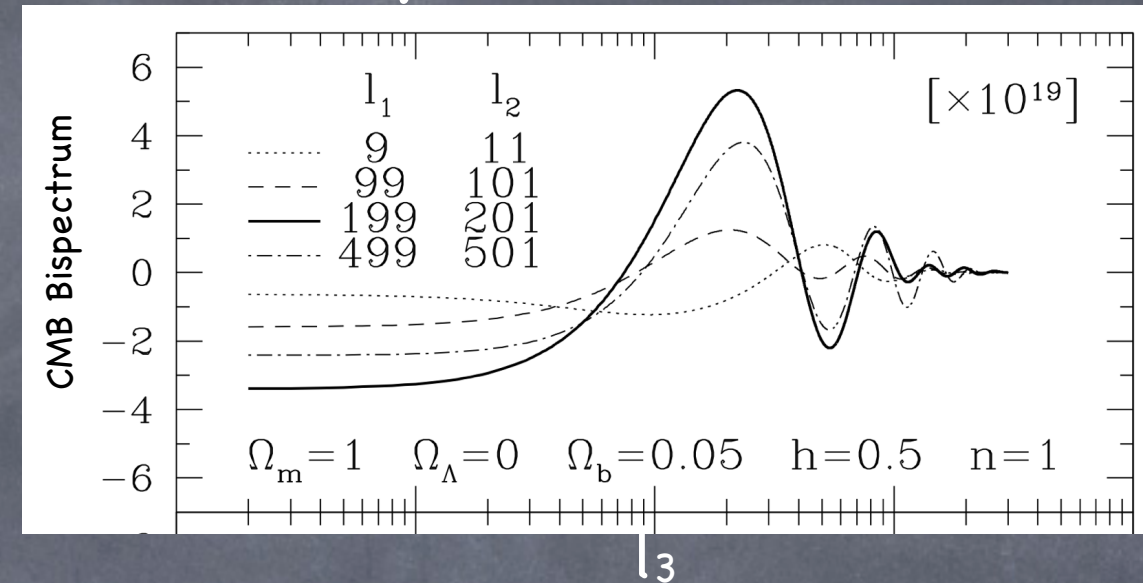
- Large values for squeezed triangle configuration
- Represents inflationary models where perturbations are generated outside the horizon
 - Multifield inflation, Curvaton
 - Inhomogeneous reheating...



Squeezed Bispectrum/Local Non-Gaussianity

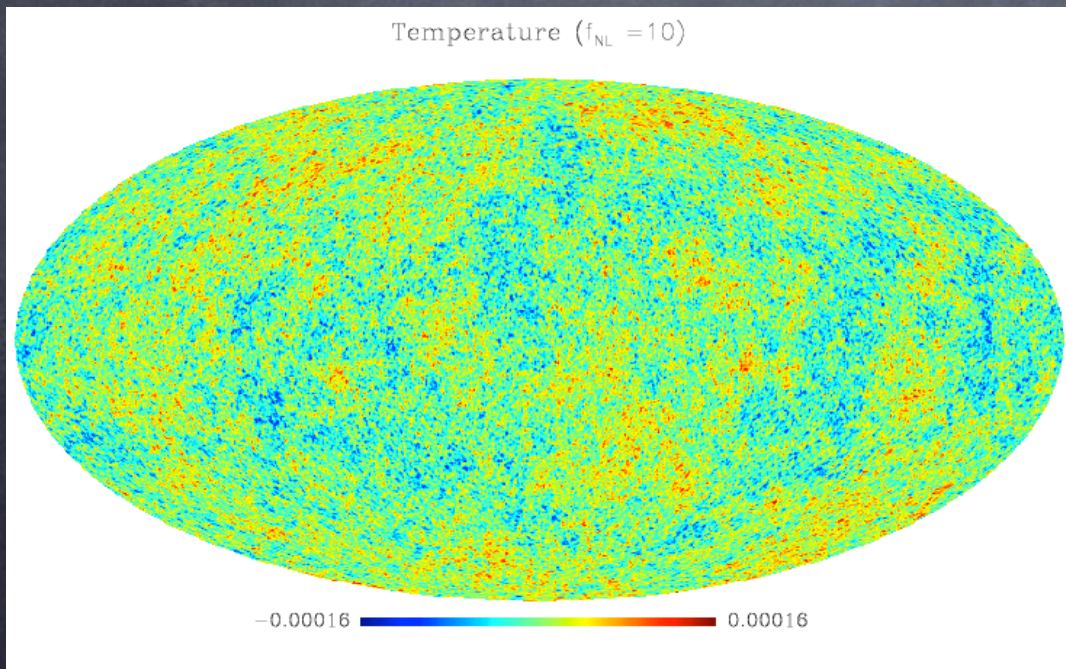
$$\Phi(x) = \Phi_G(x) + f_{NL}\Phi_G^2(x)$$

$$F(k_1, k_2, k_3) = \Delta_\Phi^2 \cdot \left(\frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3} \right)$$



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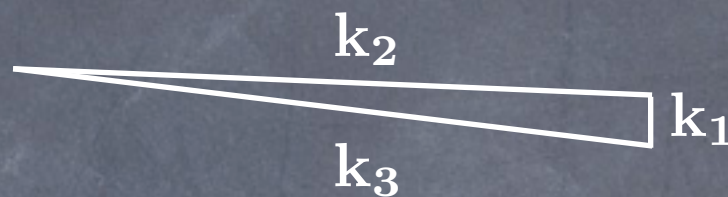
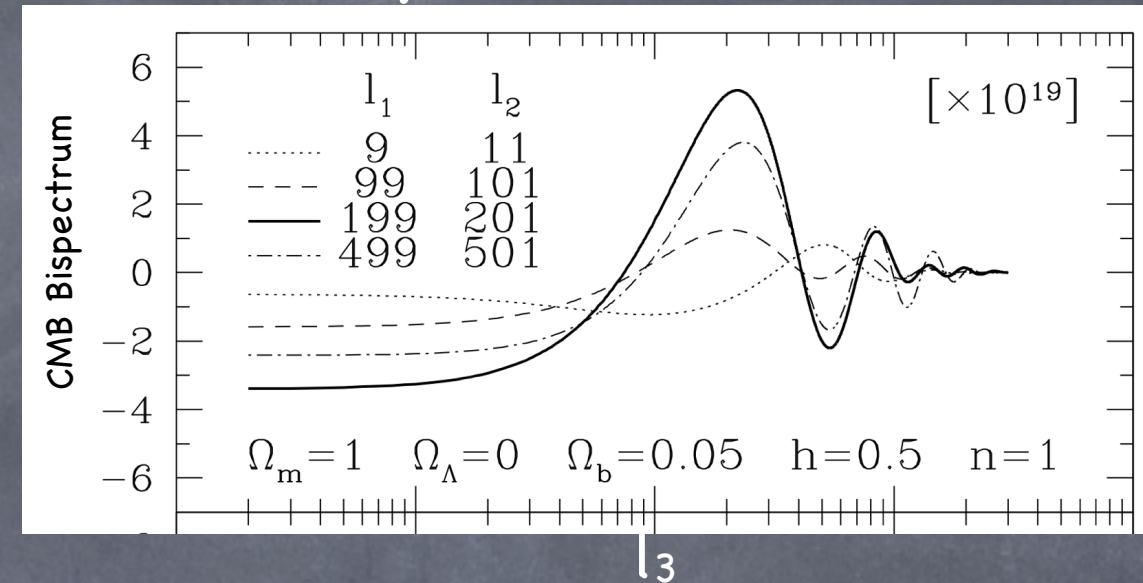
Temperature ($f_{NL} = 10$)



Squeezed Bispectrum/Local Non-Gaussianity

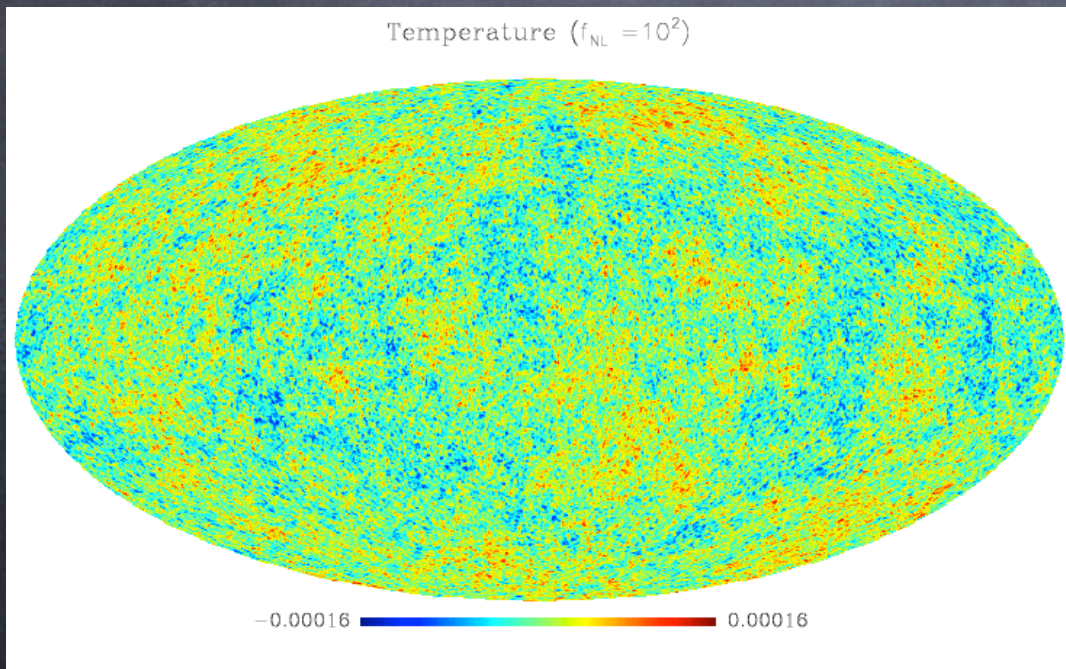
$$\Phi(x) = \Phi_G(x) + f_{NL}\Phi_G^2(x)$$

$$F(k_1, k_2, k_3) = \Delta_\Phi^2 \cdot \left(\frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3} \right)$$



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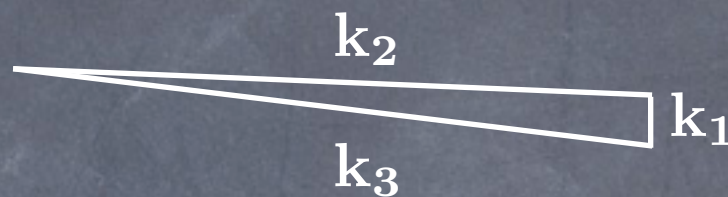
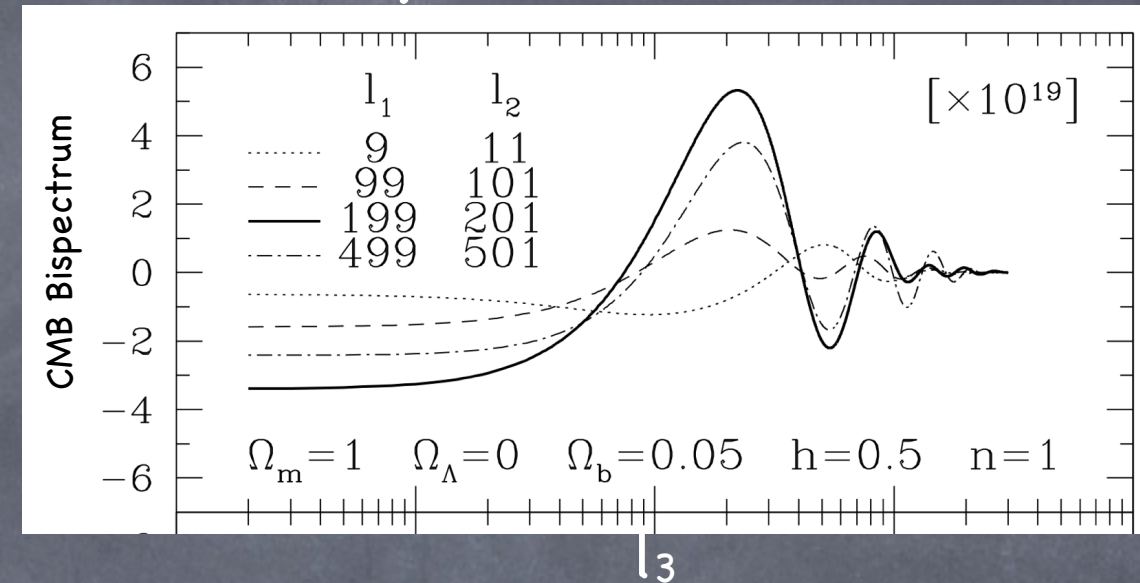
Temperature ($f_{NL} = 10^2$)



Squeezed Bispectrum/Local Non-Gaussianity

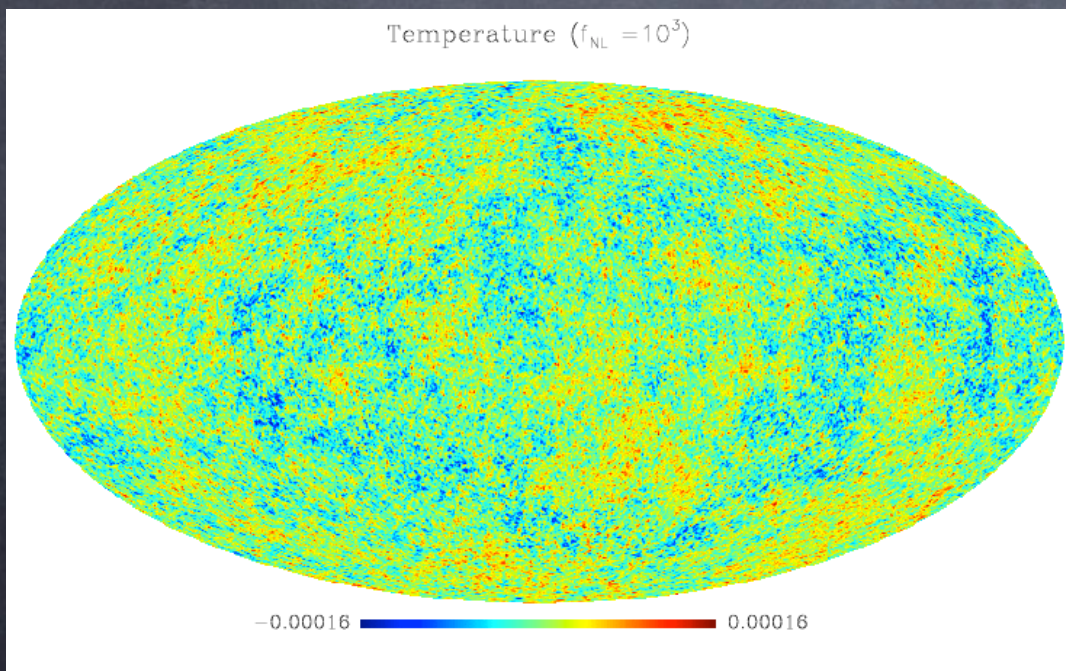
$$\Phi(x) = \Phi_G(x) + f_{NL}\Phi_G^2(x)$$

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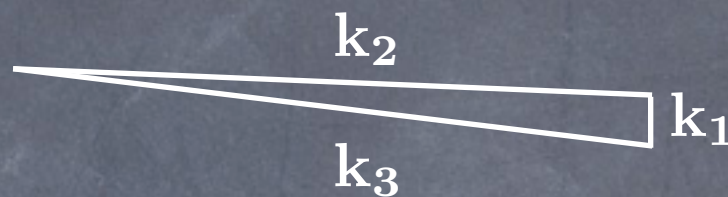
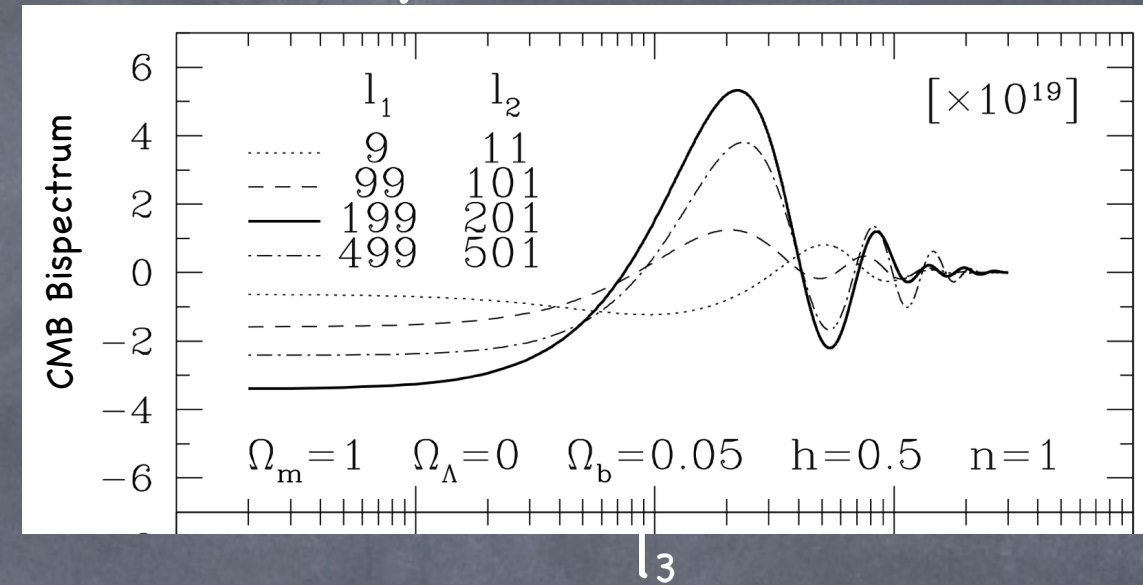
Temperature ($f_{NL} = 10^3$)



Squeezed Bispectrum/Local Non-Gaussianity

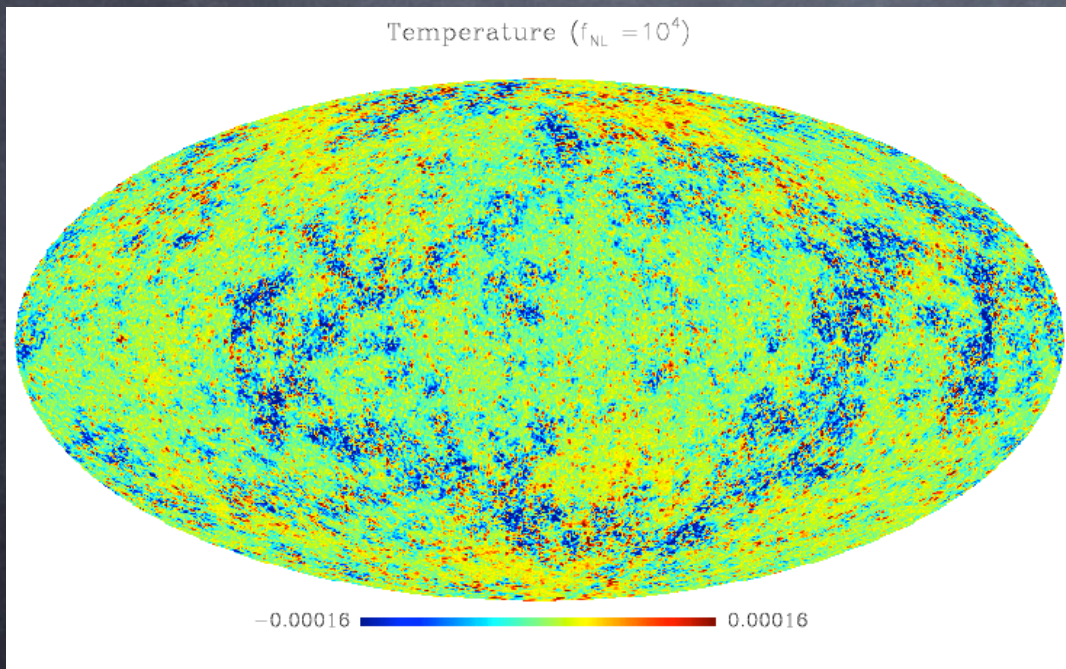
$$\Phi(x) = \Phi_G(x) + f_{NL}\Phi_G^2(x)$$

$$F(k_1, k_2, k_3) = \Delta_\Phi^2 \cdot \left(\frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3} \right)$$



- Large values for squeezed triangle configuration
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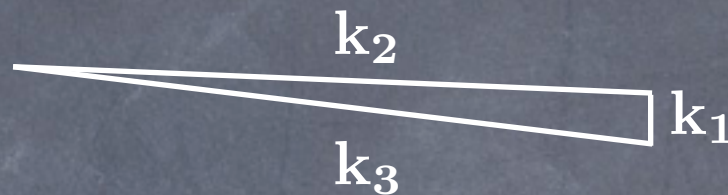
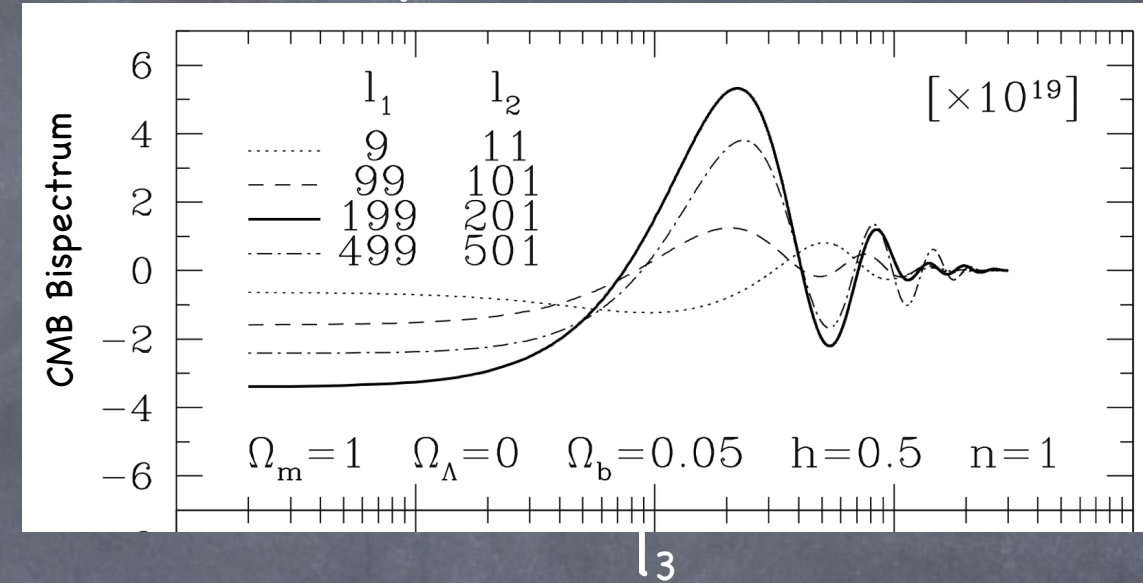
Temperature ($f_{NL} = 10^4$)



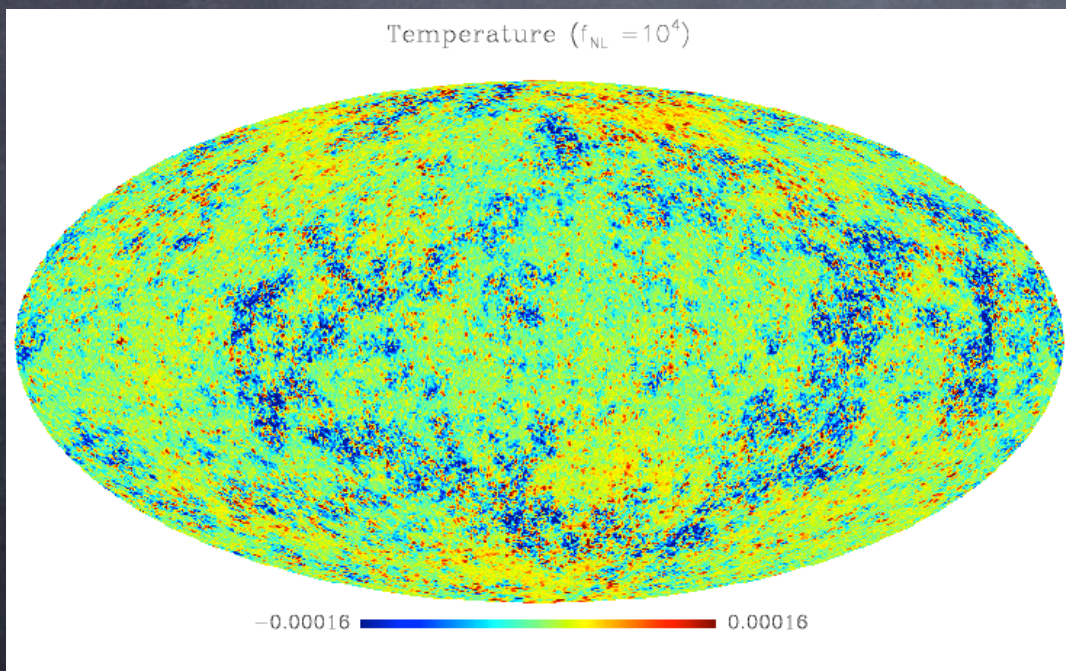
Squeezed Bispectrum/Local Non-Gaussianity

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$$F(k_1, k_2, k_3) = \Delta_\Phi^2 \cdot \left(\frac{1}{k_1^3 k_2^3} + \frac{1}{k_1^3 k_3^3} + \frac{1}{k_2^3 k_3^3} \right)$$



- Large values for squeezed triangle configuration
- Represents inflationary models where perturbations are generated outside the horizon
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data	Method	$f_{NL}^{local} \pm 2\sigma$ error	
COBE	Bispectrum non-optimal	$ f_{NL} < 1500$	Komatsu et al. (2002)
WMAP 1-year	Bispectrum non-optimal	39.5 ± 97.5	E. Komatsu et al. (2003)
WMAP 1-year	Bispectrum near-optimal-v1	47 ± 74	Creminelli et al. (2005)
WMAP 3-year	Bispectrum non-optimal	30 ± 84	Spergel et al. (2006)
WMAP 3-year	Bispectrum near-optimal-v1	32 ± 68	Creminelli et al. (2006)
WMAP 3-year	Bispectrum near-optimal	87 ± 62	Yadav and Wandelt (2008)
WMAP 3-year	Bispectrum near-optimal	69 ± 60	Smith et al. (2009)
WMAP 3-year	Bispectrum optimal	58 ± 46	Smith et al. (2009)
WMAP 5-year	Bispectrum near-optimal	51 ± 60	Komatsu et al. (2008)
WMAP 5-year	Bispectrum optimal	38 ± 42	Smith et al. (2009)
WMAP 7-year	Bispectrum optimal	32 ± 42	Komatsu et al. (2010)

$$f_{NL}^{equil} = 26 \pm 140 \text{ (68\% CL).}$$

$$f_{NL}^{orthog} = -202 \pm 104 \text{ (68\% CL).}$$

Temperature + E-Polarization

1. Cross Check: independent analysis for E and T
2. Combined T+E analysis

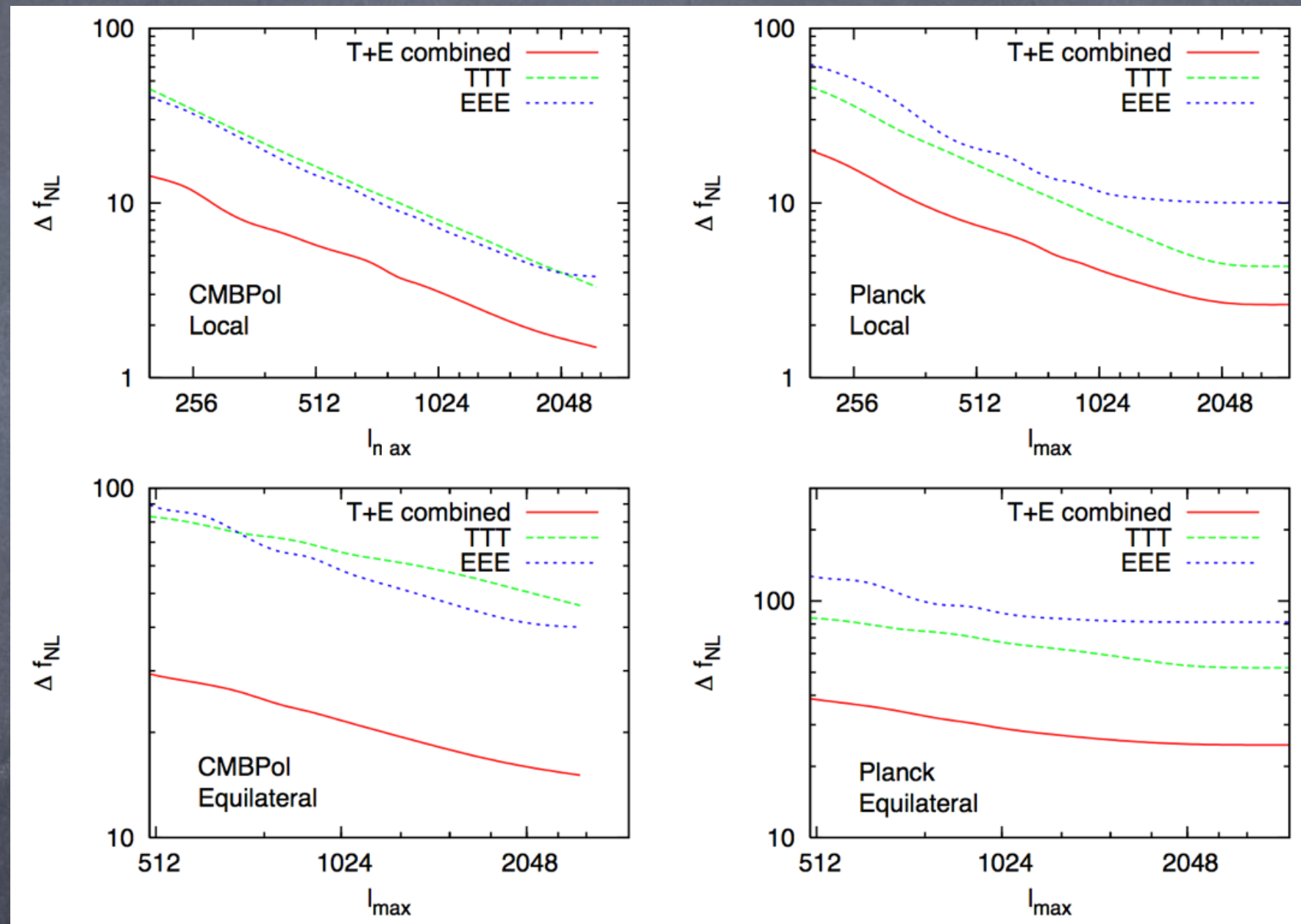
Concerns:

secondary anisotropies

ISW-lensing $f_{\text{NL}} \sim 10$ (Smith & Zaldarriaga 06)

Recombination $f_{\text{NL}} \sim \text{few}$ (Pitrou 2010)

Instrumental systematics



Babich & Zaldarriaga (2004); Yadav, Komatsu, Wandelt (2007)
Yadav & Wandelt (2005)

Beyond f_{NL} : the case for running non-Gaussianity

- f_{NL} has been shown to be scale-dependent in several models of with a variable speed of sound, such as DBI inflation.

- Non-Gaussianity can be larger (or smaller) at different scales and for different observables

$$f_{NL} \rightarrow f_{NL} \left(\frac{K}{k_{\star}} \right)^{n_{NG}}$$

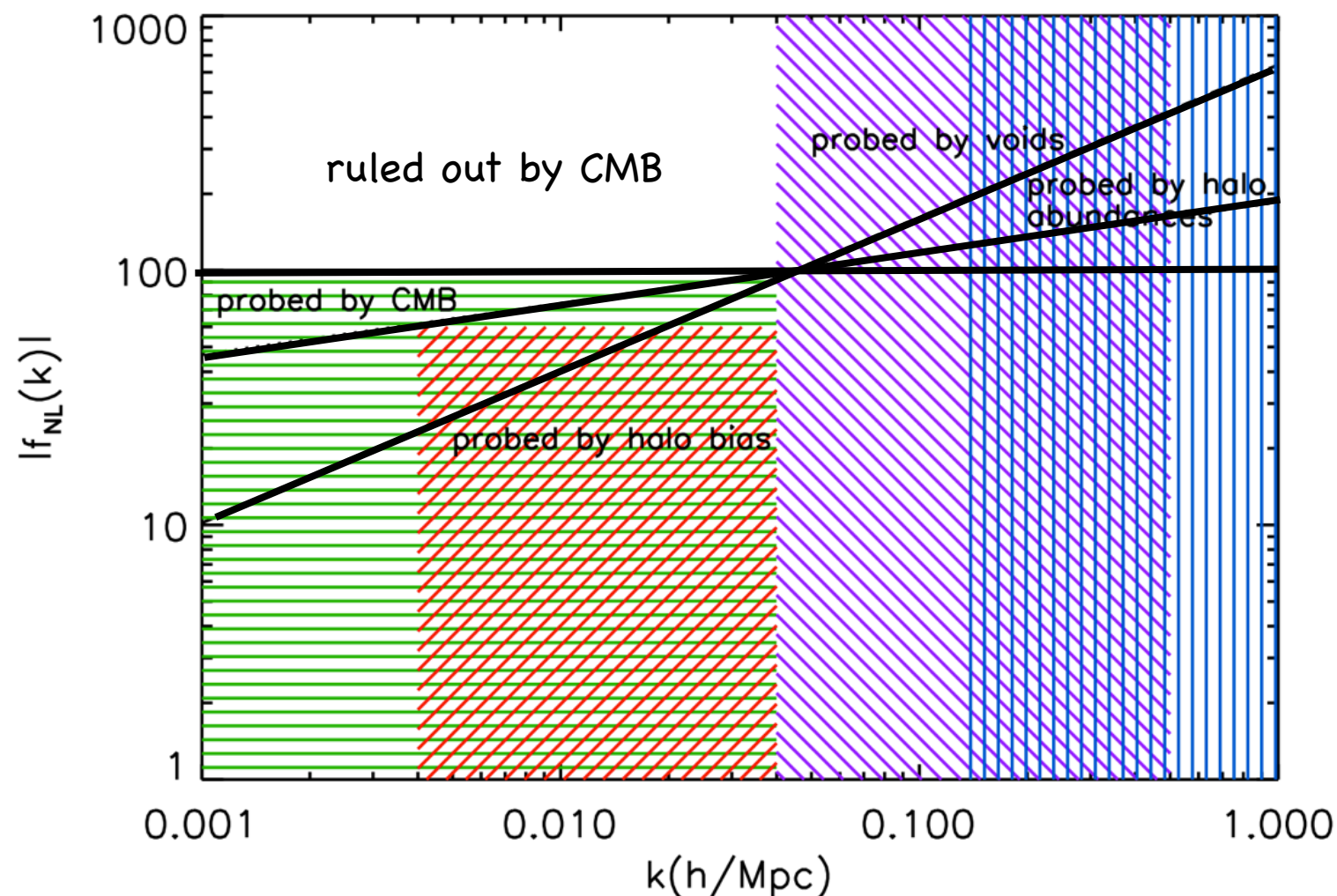


Figure stolen from, Emiliano Sefusatti's talk

$$\begin{aligned} n_{NG} &= 0 \\ n_{NG} &= 0.2 \\ n_{NG} &= 0.6 \end{aligned}$$

LoVerde, Miller, Shandera, Verde (2008)
Sefusatti, Liguori, Yadav, Jackson, Pajer (2009)

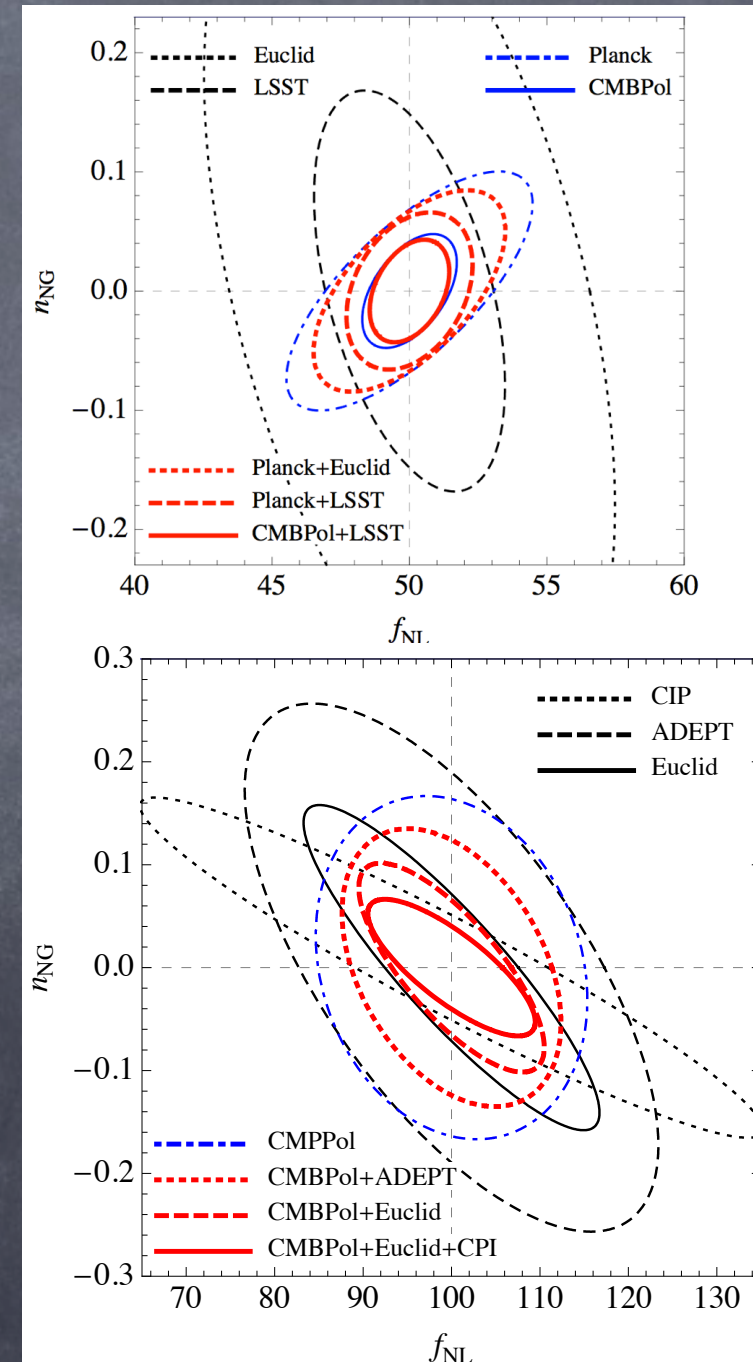
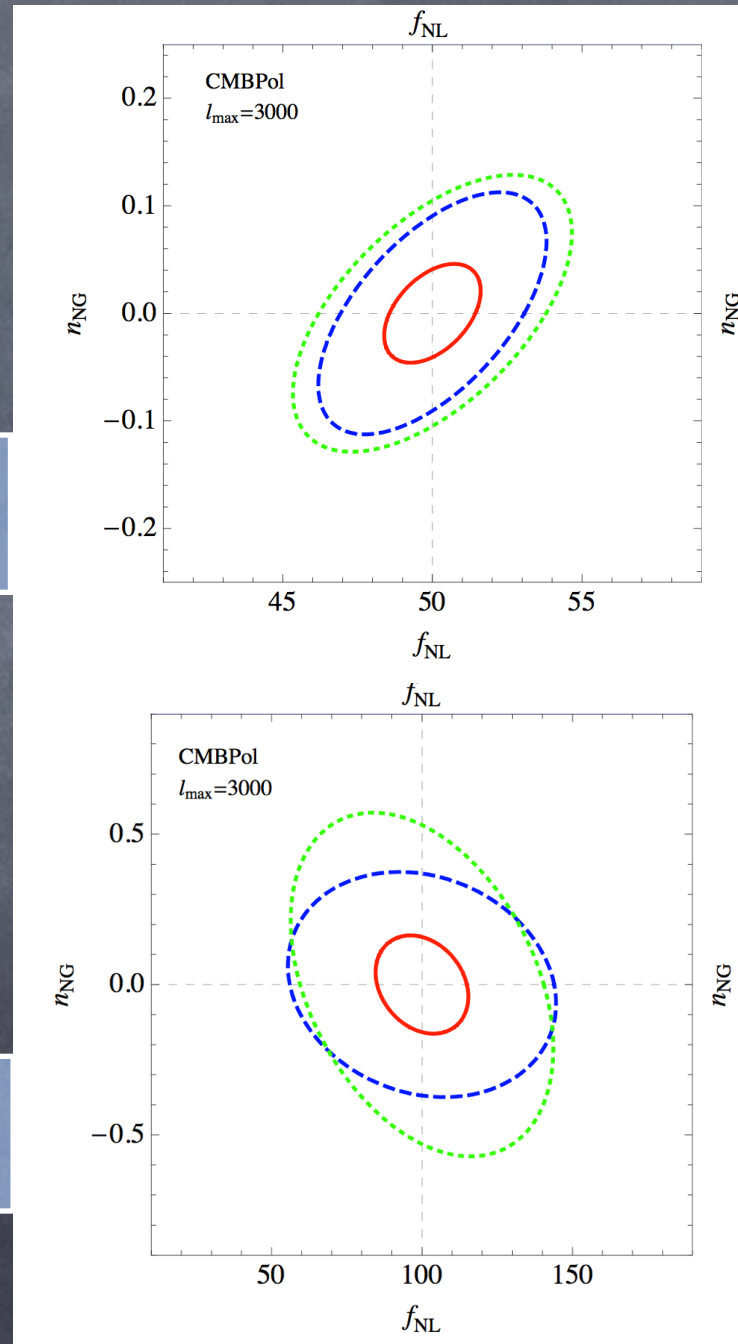
Beyond f_{NL} : the case for running non-Gaussianity

- f_{NL} has been shown to be scale-dependent in several models of with a variable speed of sound, such as DBI inflation.
- Current observations only constrain the magnitude of f_{NL} . However if f_{NL} is large enough, it may be also possible in the near future to constrain its possible dependence on scale.

$$f_{\text{NL}} \rightarrow f_{\text{NL}} \left(\frac{K}{k_{\star}} \right)^{n_{\text{NG}}}$$

Local

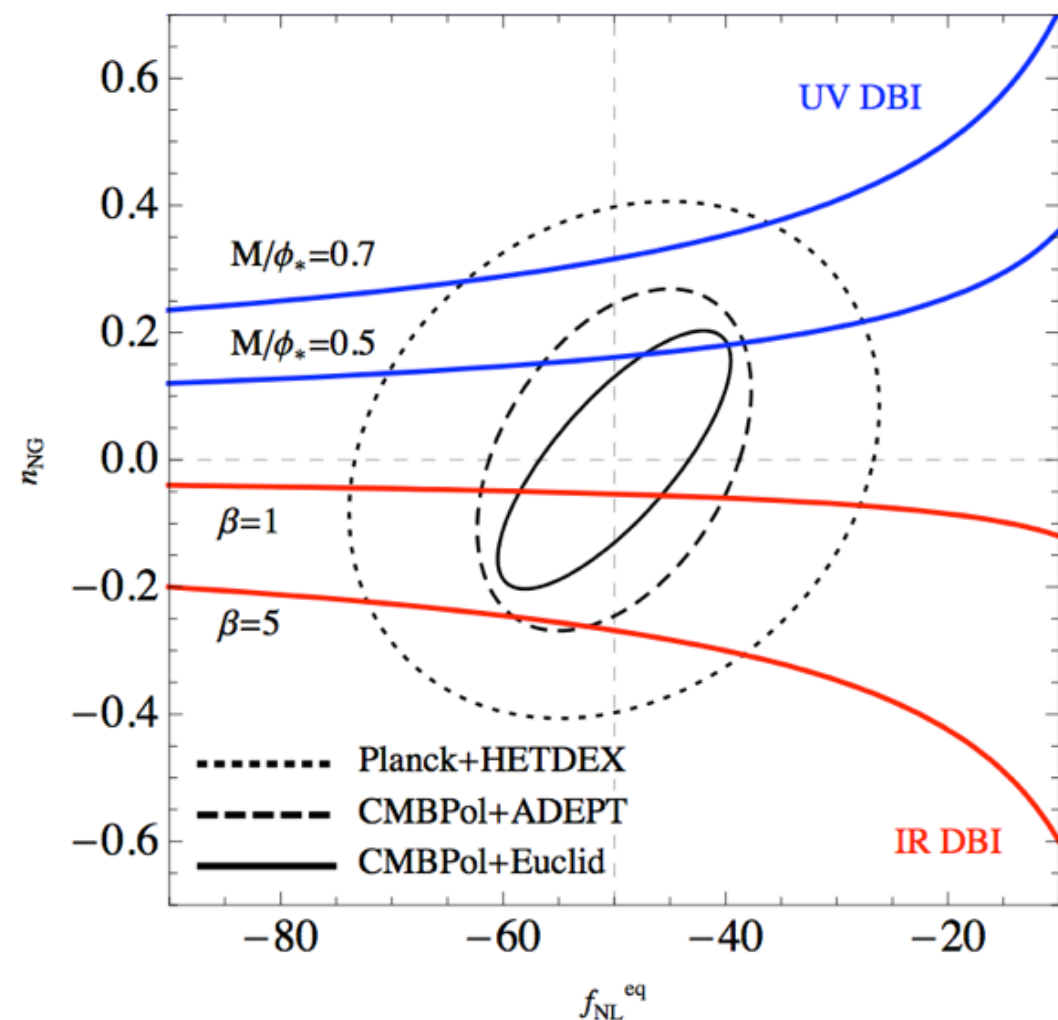
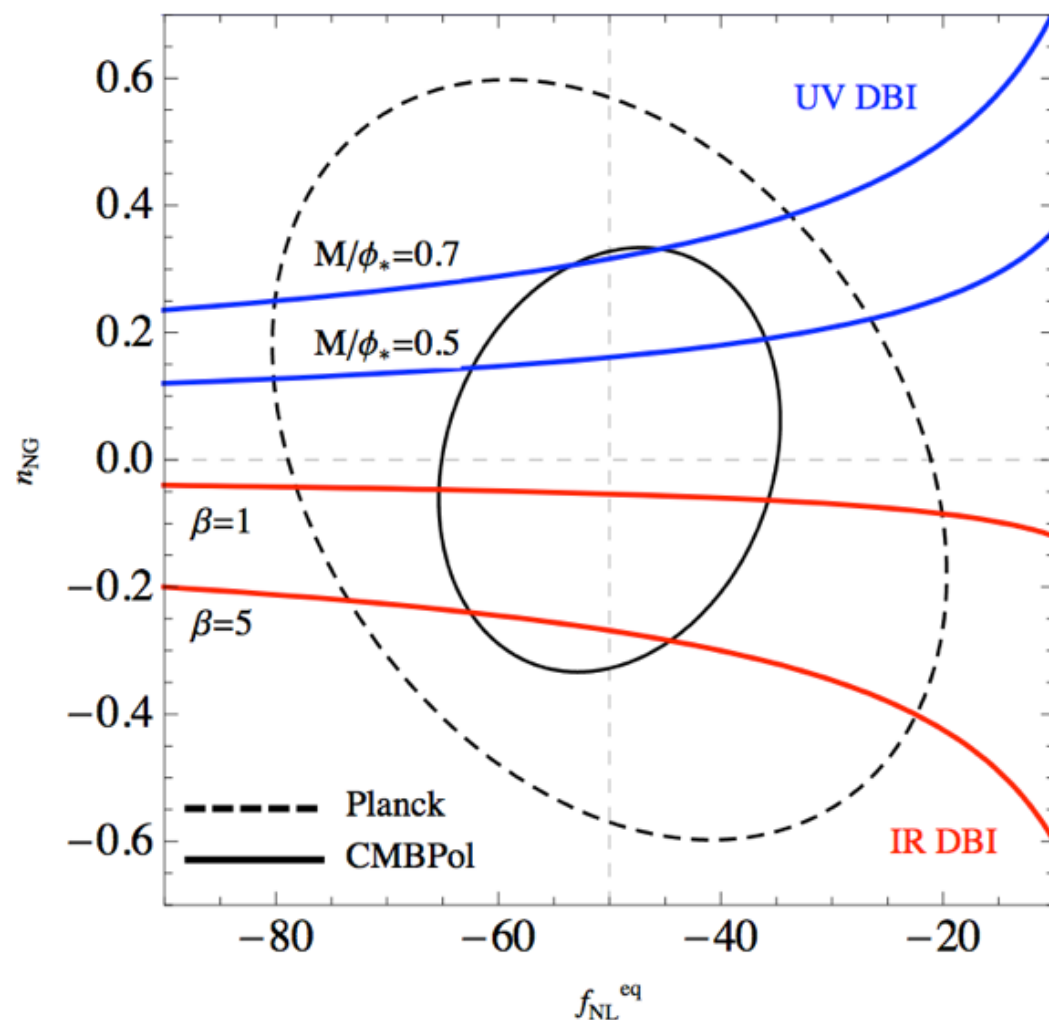
Equilateral



Beyond f_{NL} : the case for running non-Gaussianity

An example: DBI Inflation

$$\mathcal{L}(\phi) = -\frac{\phi^4}{\lambda} \sqrt{1 - \lambda \frac{\dot{\phi}^2}{\phi^4}} + \frac{\phi^4}{\lambda} - V(\phi), \quad V(\phi) = \begin{cases} \frac{1}{2} m \phi^2 & \text{UV} \\ V_0 - \frac{1}{2} m \phi^2 & \text{IR} \end{cases} \quad (m^2 \equiv \beta H^2)$$



Limits on n_{NG} can provide stronger constraints on the parameters of the models

Instrumental effects on CMB Bispectrum

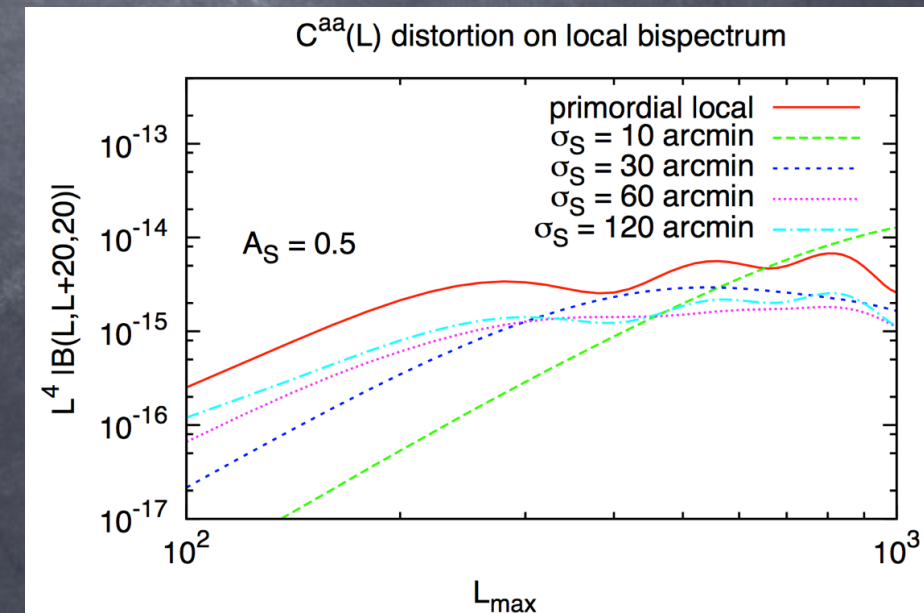
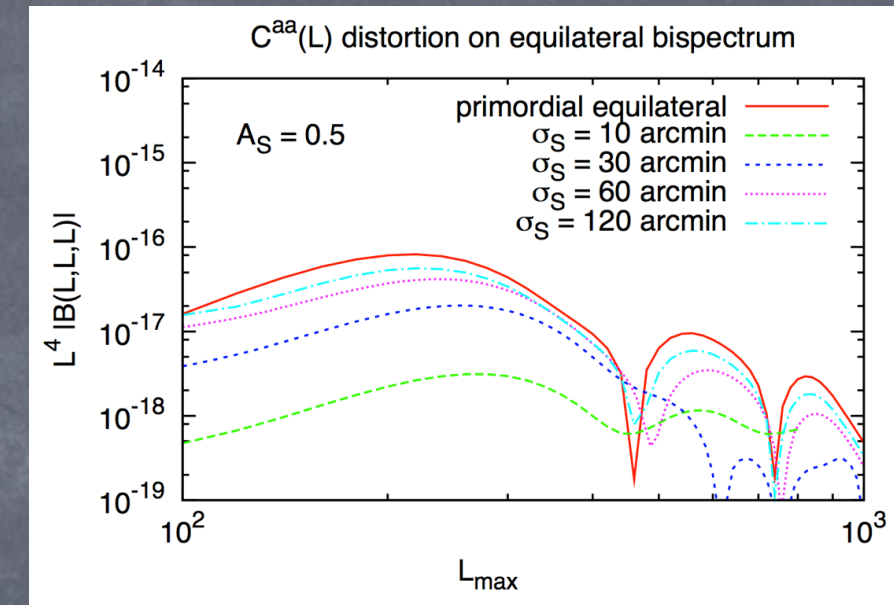
Su, Yadav et al. (2010)

Detector gain: $\delta T^{out}(n) = (1 + a(n))\delta T^{in}(n)$

$$\delta B_{(l_1, l_2, l_3)}^{TTT} = \int \frac{d^2 l'}{(2\pi)^2} C_{l'}^{aa} \left\{ B_{(l_1, l_2 - l', l_3 + l')}^{TTT} + B_{(l_1 - l', l_2 + l', l_3)}^{TTT} + B_{(l_1 + l', l_2, l_3 - l')}^{TTT} \right\}$$

$$C_\ell^{aa} \propto \exp(-\ell(\ell + 1)\sigma_S^2)$$

(1) Linear systematics can only distort primordial bispectrum



Instrumental effects on CMB Bispectrum

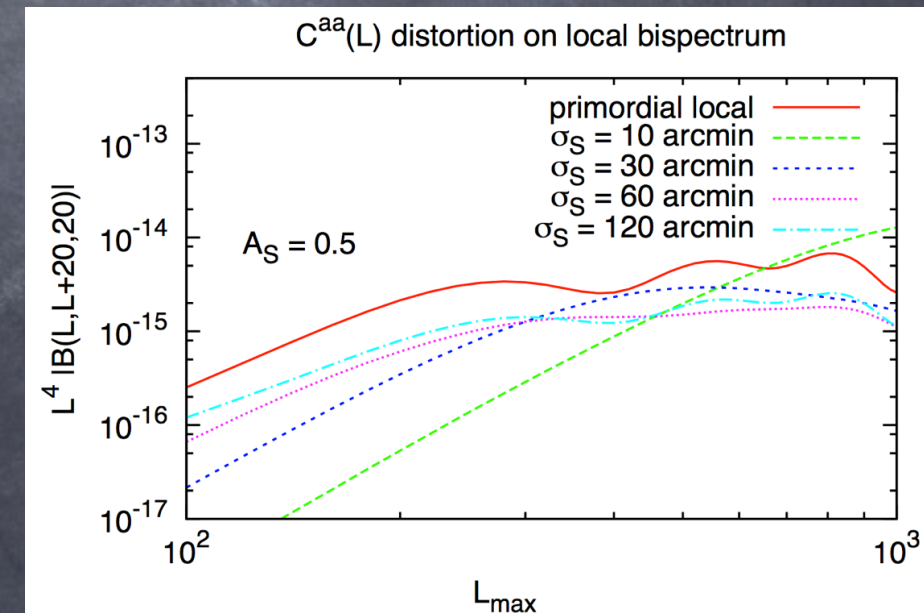
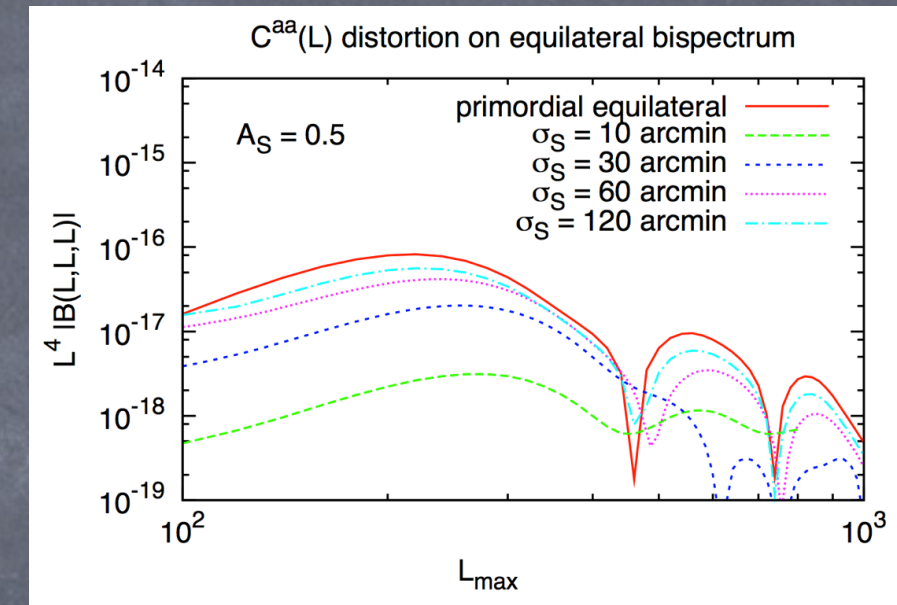
Su, Yadav et al. (2010)

Detector gain: $\delta T^{out}(n) = (1 + a(n))\delta T^{in}(n) + b(\delta T^{in})^2$

$$\delta B_{(l_1, l_2, l_3)}^{TTT} = \int \frac{d^2 l'}{(2\pi)^2} C_{l'}^{aa} \left\{ B_{(l_1, l_2 - l', l_3 + l')}^{TTT} + B_{(l_1 - l', l_2 + l', l_3)}^{TTT} + B_{(l_1 + l', l_2, l_3 - l')}^{TTT} \right\}$$

$$C_\ell^{aa} \propto \exp(-\ell(\ell + 1)\sigma_S^2)$$

(1) Linear systematics can only distort primordial bispectrum



Instrumental effects on CMB Bispectrum

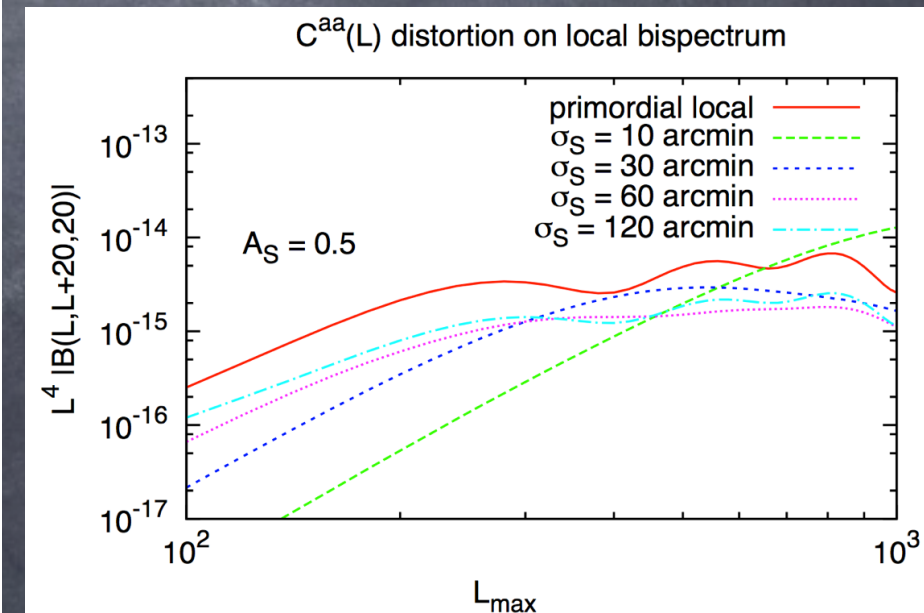
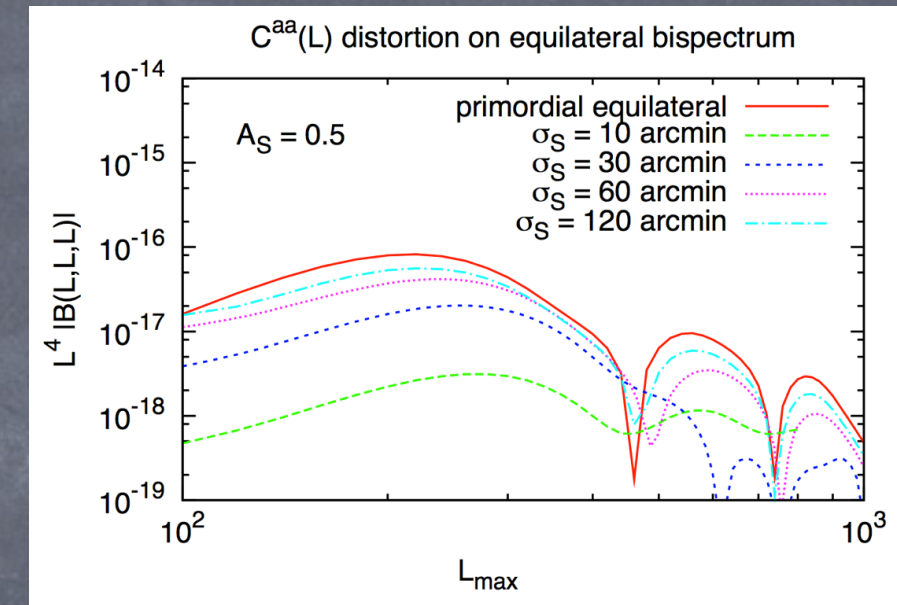
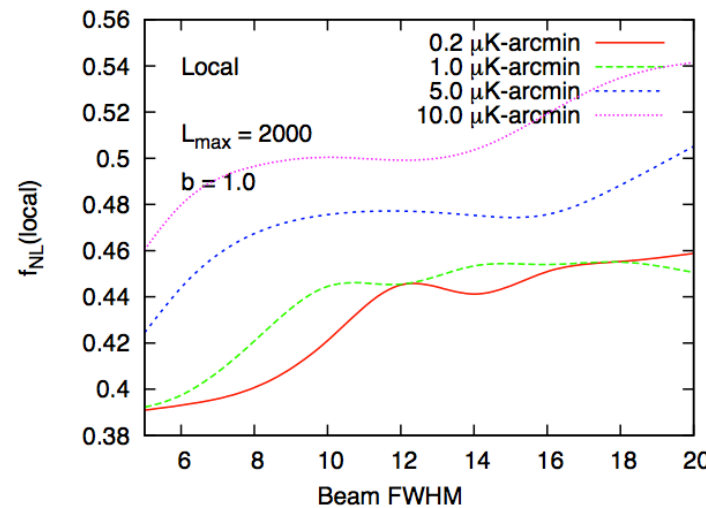
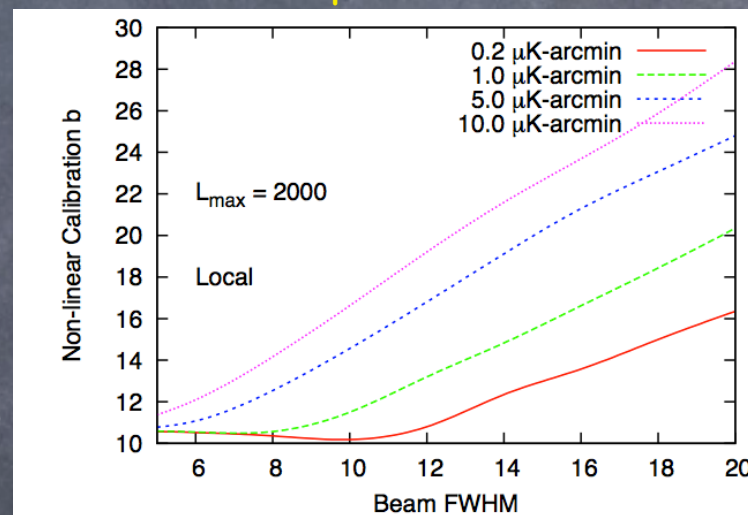
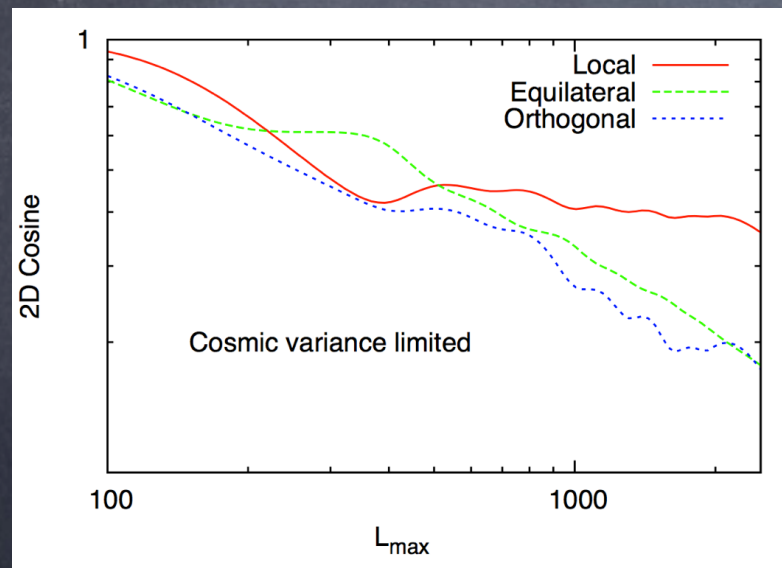
Su, Yadav et al. (2010)

Detector gain: $\delta T^{out}(n) = (1 + a(n))\delta T^{in}(n) + b (\delta T^{in})^2$

$$\delta B_{(l_1, l_2, l_3)}^{TTT} = \int \frac{d^2 l'}{(2\pi)^2} C_{l'}^{aa} \left\{ B_{(l_1, l_2 - l', l_3 + l')}^{TTT} + B_{(l_1 - l', l_2 + l', l_3)}^{TTT} + B_{(l_1 + l', l_2, l_3 - l')}^{TTT} \right\}$$

$$C_\ell^{aa} \propto \exp(-\ell(\ell + 1)\sigma_S^2)$$

- (1) Linear systematics can only distort primordial bispectrum
- (2) Instrumental nonlinearities can generate spurious bispectrum and if these non-linearities are not controlled by dedicated calibration, they can produce $f_{NL} \sim O(10)$ before their effect is visible in CMB dipole.



Instrumental effects on CMB polarization Bispectrum

Su, Yadav et al. (2010)

Differential Gain $\delta[Q + iU](n) = [\gamma_1 + i\gamma_2](n)T(n)$

Differential Pointing $\delta[Q + iU](n) = \sigma[d_1 + id_2](n)[\partial_1 + \partial_2]T(n)$

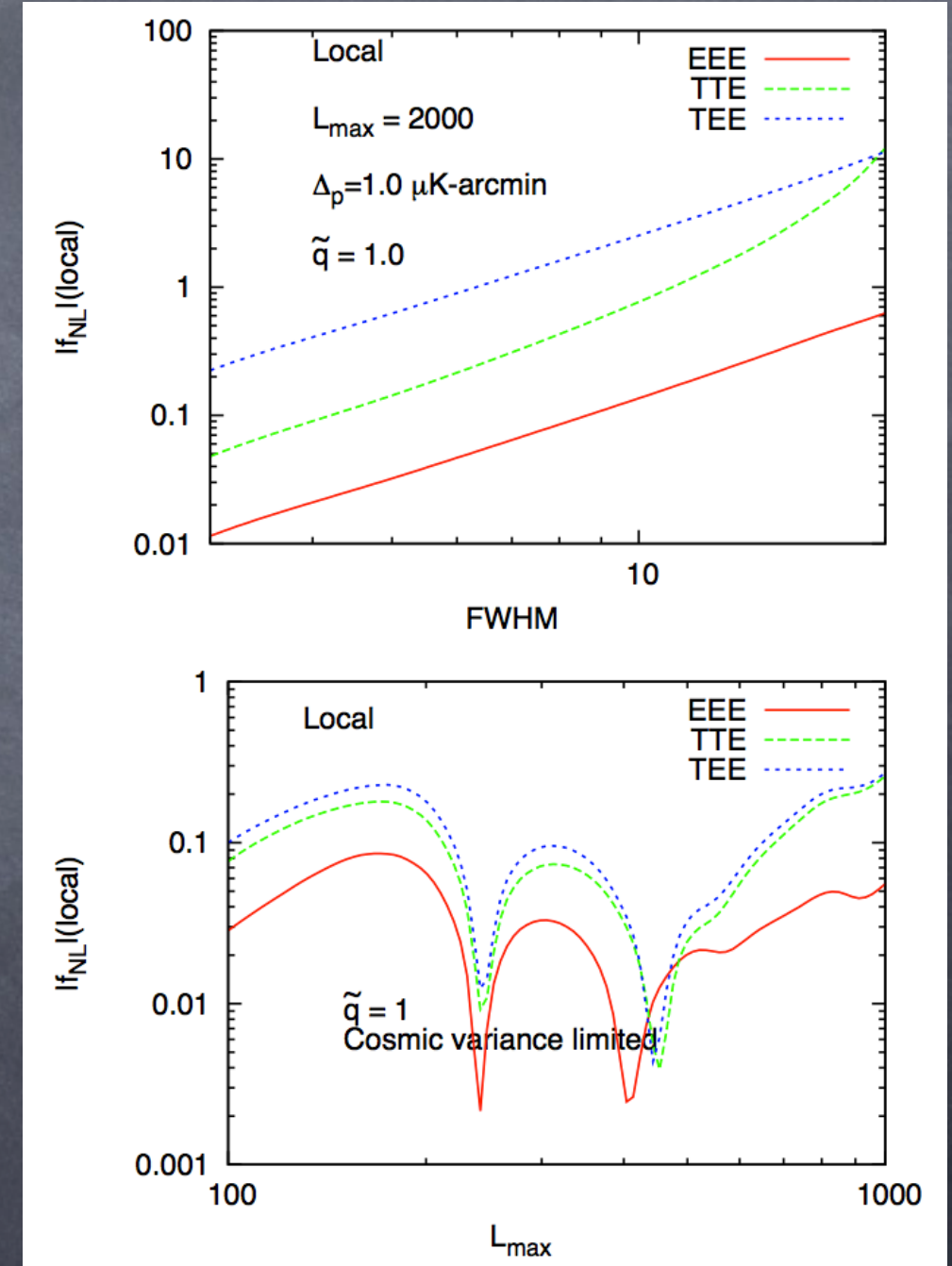
Differential Ellipticity $\delta[Q + iU](n) = \sigma^2 q(n)[\partial_1 + \partial_2]^2 T(n)$

$$B_{(\mathbf{l}_1, \mathbf{l}_2, \mathbf{l}_3)}^{EEE, \tilde{S}, bias} = \tilde{S}(0) C_{\ell_1}^{TE} C_{\ell_2}^{TE} W_E^{\tilde{S}}(\mathbf{l}_3, -\mathbf{l}_1, -\mathbf{l}_2) + \text{perm.}$$

$$B_{(\mathbf{l}_1, \mathbf{l}_2, \mathbf{l}_3)}^{TTE, \tilde{S}, bias} = \tilde{S}(0) C_{\ell_1}^{TT} C_{\ell_2}^{TT} W_E^{\tilde{S}}(\mathbf{l}_3, -\mathbf{l}_1, -\mathbf{l}_2) + \text{perm.}$$

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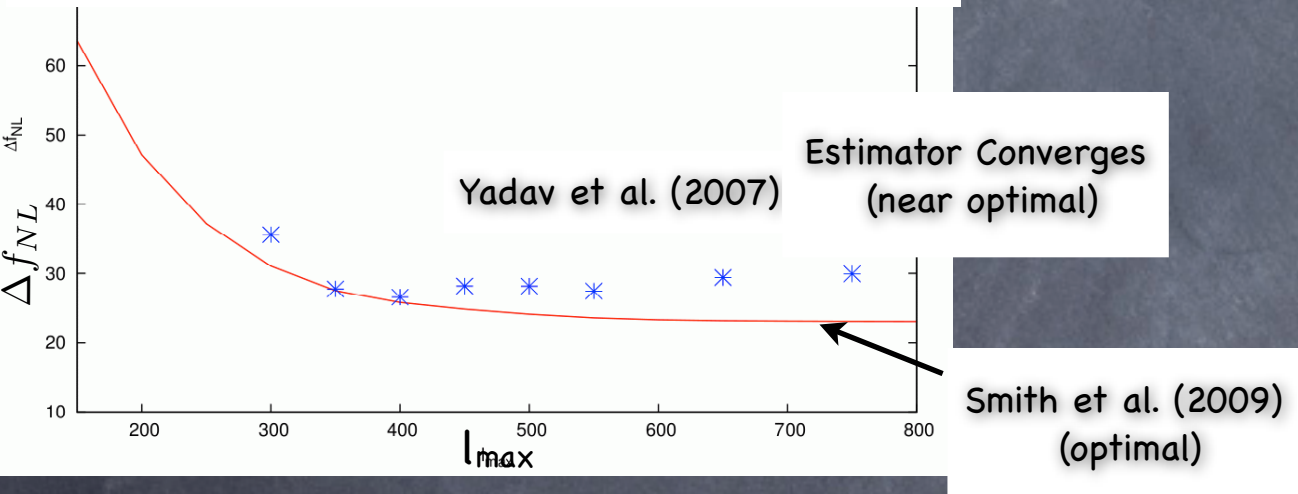
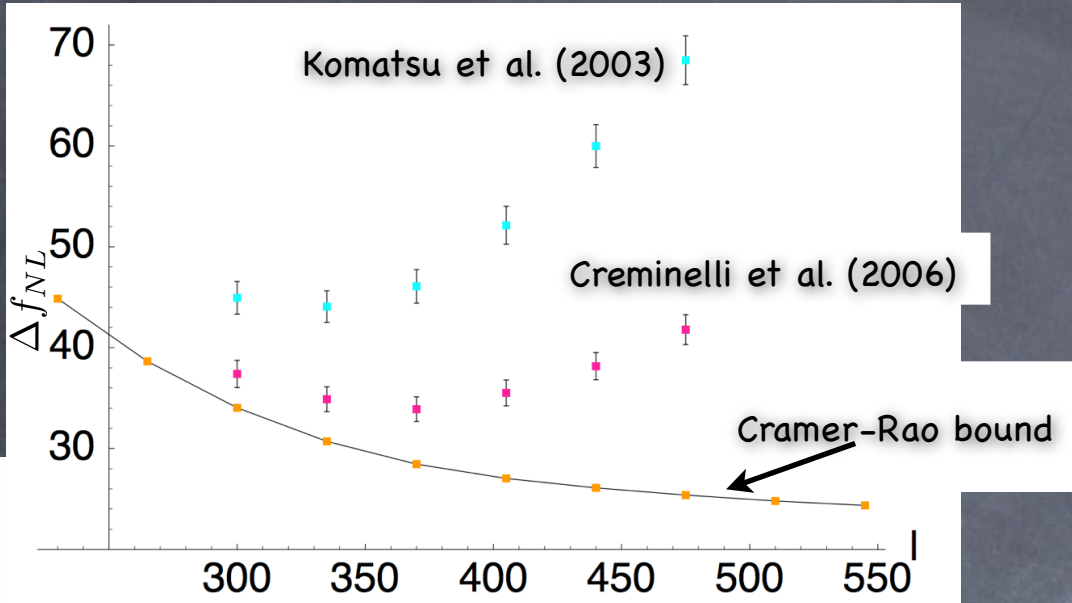
Non-linear Systematics $\tilde{S}$	$W_E^{\tilde{S}}(\mathbf{l}, \mathbf{l}', \mathbf{l}'')$
Monopole leakage $\tilde{\gamma}_a$	$2 \cos[2(\varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}''})]$
Monopole leakage $\tilde{\gamma}_b$	$-2 \sin[2(\varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}''})]$
Dipole leakage $\tilde{d}_a$	$\sigma[\mathbf{l}' \sin(\varphi_{\mathbf{l}'} + \varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}''}) + \mathbf{l}'' \sin(\varphi_{\mathbf{l}''} + \varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}'})]$
Dipole leakage $\tilde{d}_b$	$\sigma[\mathbf{l}' \cos(\varphi_{\mathbf{l}'} + \varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}''}) + \mathbf{l}'' \cos(\varphi_{\mathbf{l}''} + \varphi_{\mathbf{l}-\mathbf{l}'}-\varphi_{\mathbf{l}'})]$
Quadrupole leakage $\tilde{q}$	$-2\sigma^2\{\mathbf{l}'^2 \cos[2(\varphi_{\mathbf{l}'} - \varphi_{\mathbf{l}})] + \mathbf{l}''^2 \cos[2(\varphi_{\mathbf{l}''} - \varphi_{\mathbf{l}})]\}$



Executive Summary

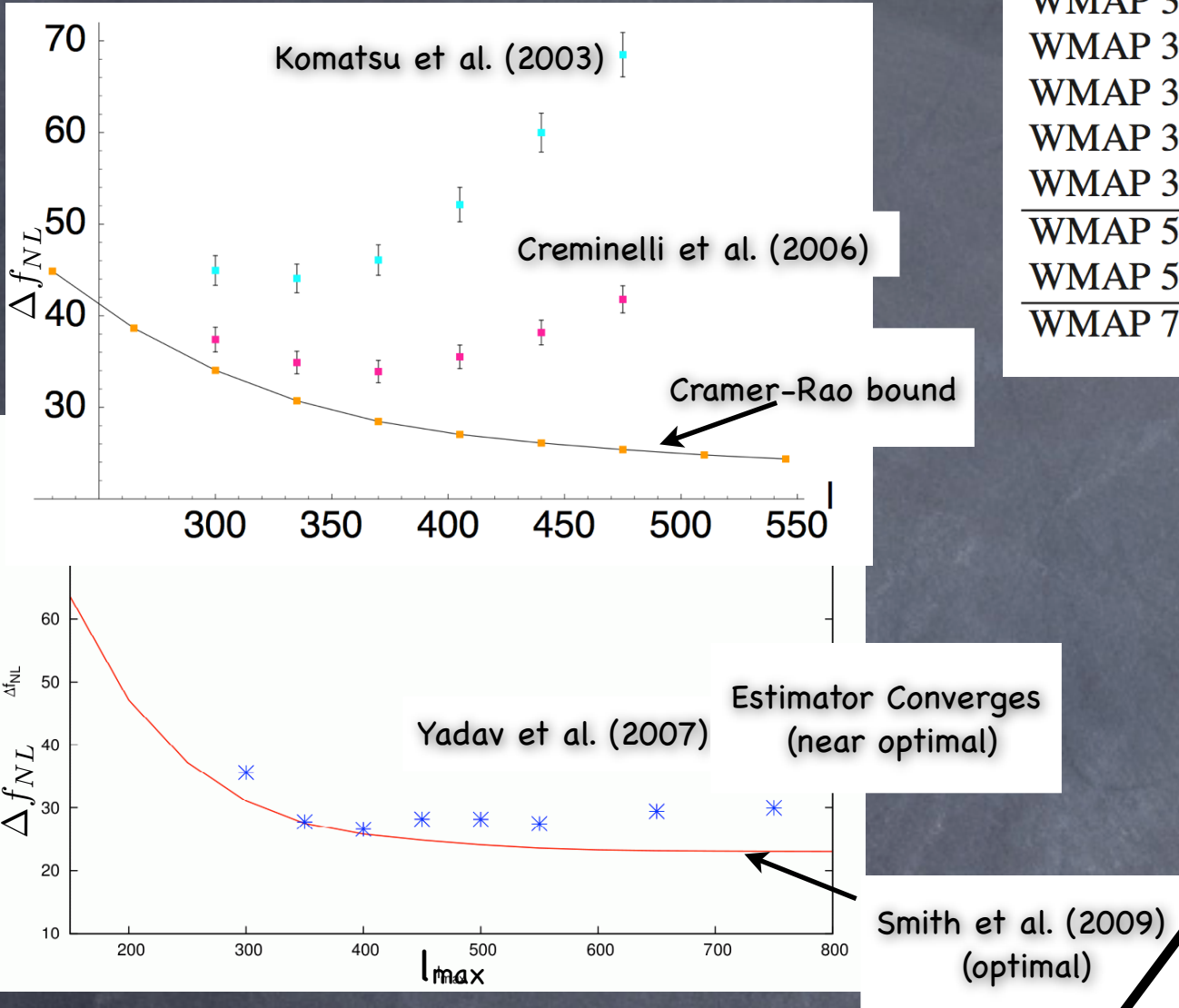
- Characterization of non-Gaussianity is a powerful probe of the the early universe. Non-Gaussianity measures the strength of interaction and field content during inflation.
- Current status with CMB: $\Delta f_{\text{NL}} \sim 20$, $f_{\text{NL}} \sim 40$.
- Using combined T+E data of Planck: $\Delta f_{\text{NL}} \sim 4$ (in few years)
- If nonG is indeed large, running can also be constrained.
- If instrumental non-linearities are not controlled by dedicated calibration, it can generate $f_{\text{NL}} \sim \mathcal{O}(10)$ before showing up in CMB dipole.

Constraints on local f_{NL}



data	Method	$f_{NL}^{local} \pm 2\sigma$ error	
COBE	Bispectrum non-optimal	$ f_{NL} < 1500$	Komatsu et al. (2002)
WMAP 1-year	Bispectrum non-optimal	39.5 ± 97.5	E. Komatsu et al. (2003)
WMAP 1-year	Bispectrum near-optimal-v1	47 ± 74	Creminelli et al. (2005)
WMAP 3-year	Bispectrum non-optimal	30 ± 84	Spergel et al. (2006)
WMAP 3-year	Bispectrum near-optimal-v1	32 ± 68	Creminelli et al. (2006)
WMAP 3-year	Bispectrum near-optimal	87 ± 62	Yadav and Wandelt (2008)
WMAP 3-year	Bispectrum near-optimal	69 ± 60	Smith et al. (2009)
WMAP 3-year	Bispectrum optimal	58 ± 46	Smith et al. (2009)
WMAP 5-year	Bispectrum near-optimal	51 ± 60	Komatsu et al. (2008)
WMAP 5-year	Bispectrum optimal	38 ± 42	Smith et al. (2009)
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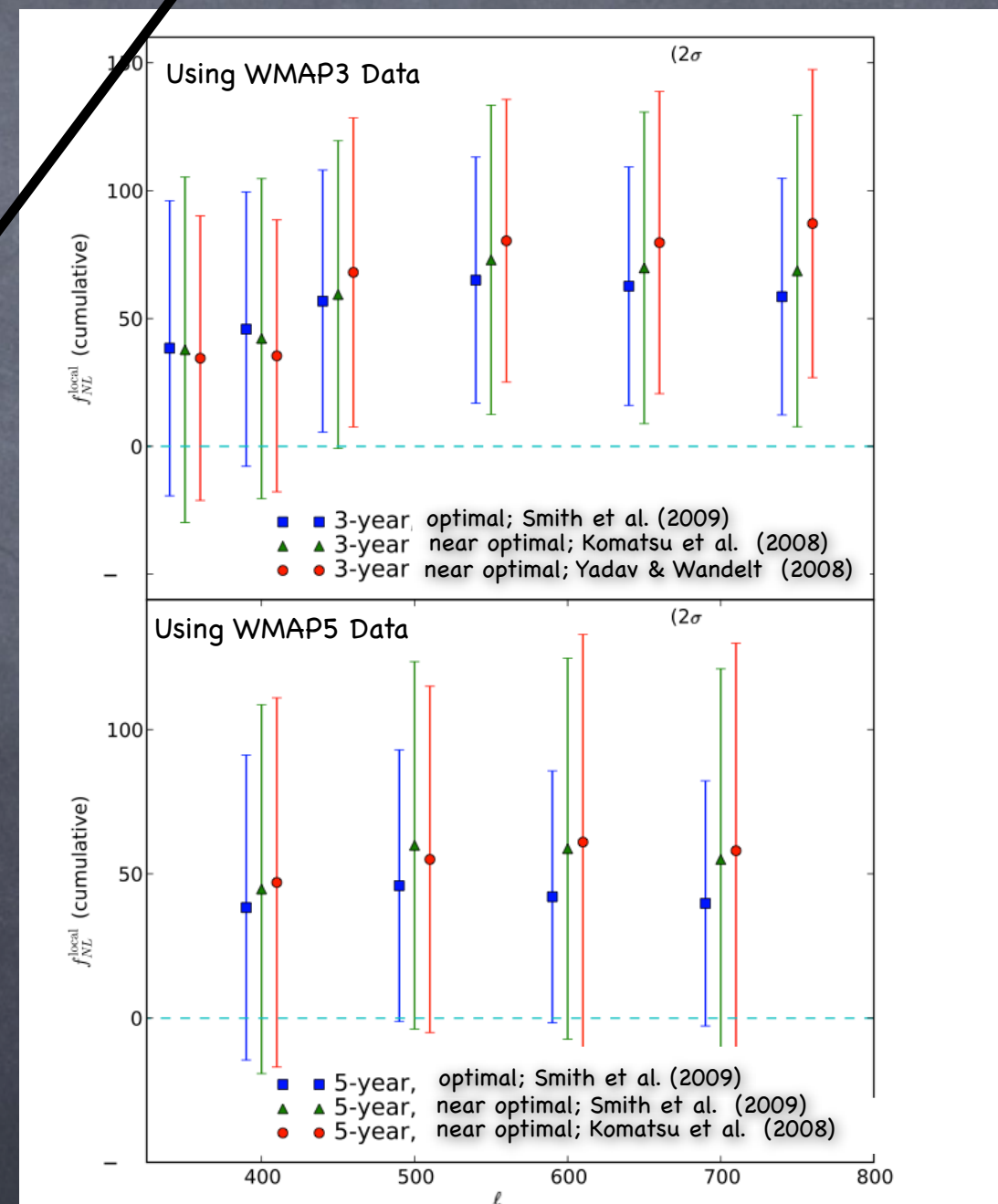
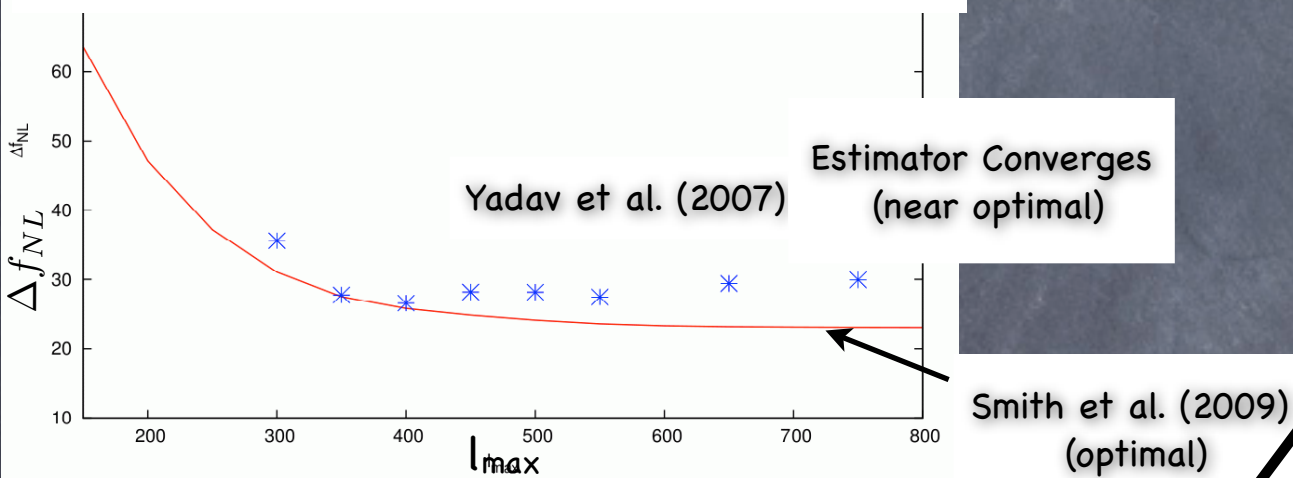
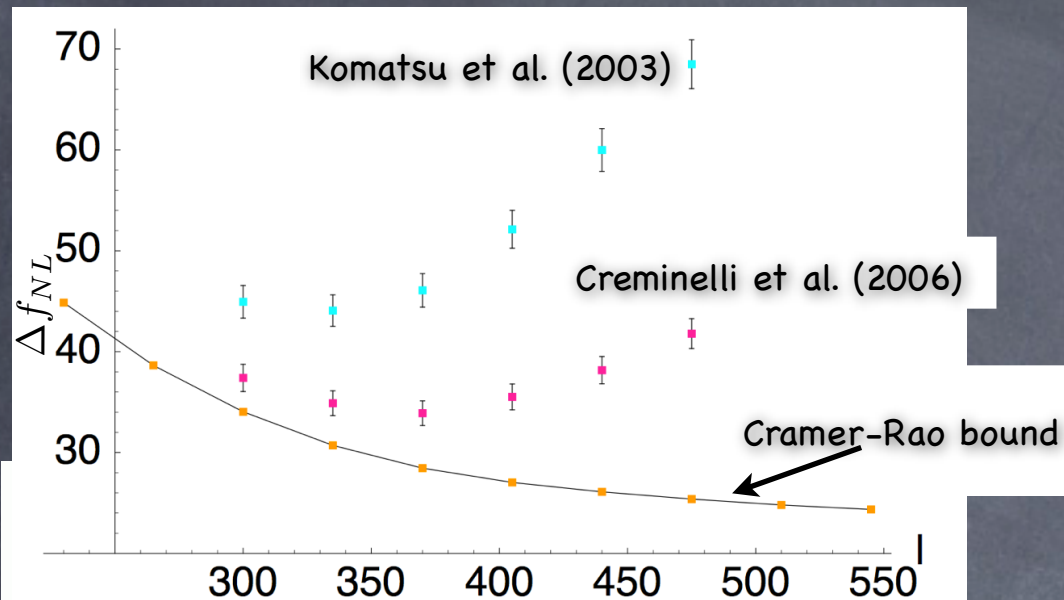
Hint/Evidence of non-Gaussianity

$\sim 2.5\sigma$ deviation from Gaussianity, in wmap3 year data

Yadav & Wandelt (2008); Smith, Senatore and Zaldarriaga 2009)

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